

**EFFECT OF LOCAL EQUILIBRATION IN DISSIPATIVE  
HEAVY-ION COLLISION**

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The main objective of this communication is to point out and emphasize the influence of the local equilibration process on the observable quantities in the collision of two heavy ions. No attempt has been made to provide a best fit to the experimental data. The calculations have been performed within the framework of the Fokker-Planck equation. The collective variables of the reaction are the relative position  $r$  and momentum  $p$ , the angle  $\theta$  and angular momentum  $l$ , the mass asymmetry  $\alpha$ , and the deformation  $\varepsilon$  related to the neck size. The Fokker-Planck equation for the probability distribution function  $f(r, p, \theta, l, \alpha, \varepsilon, t)$  is written as

$$\begin{aligned} \partial f / \partial t = & - p / \mu \partial f / \partial r + \partial U_{\text{dyn}} / \partial r \partial f / \partial p - \partial (v_p f) / \partial p + \partial^2 (D_{pp} f) / \partial p^2 \\ & - l / J_{\text{rel}} \partial f / \partial \theta - \partial (v_l f) / \partial l + \partial^2 (D_{ll} f) / \partial l^2 \\ & - \partial (v_\alpha f) / \partial \alpha + \partial^2 (D_{\alpha\alpha} f) / \partial \alpha^2 \\ & - \partial (v_\varepsilon f) / \partial \varepsilon + \partial^2 (D_{\varepsilon\varepsilon} f) / \partial \varepsilon^2 \end{aligned} \quad (1)$$

where  $\mu$  and  $J_{\text{rel}}$  denote, respectively, the reduced mass and the relative moment of inertia. The transport coefficients  $v_i$  and  $D_{ij}$  are obtained from the microscopic expressions of refs.<sup>3,4</sup>) with a strength factor  $f(r)$ . The form factor  $f(r)$  is given by the square root of the density overlap normalized at the contact radius. The effect of local equilibration is treated<sup>2</sup>) through a time dependent dynamical potential

$$\begin{aligned} U_{\text{dyn}} = & U_{\text{ad}} [1 - \chi(t)] + U_{\text{diab}}(r) \chi(t) = U_{\text{ad}} + (\Delta U)_{\text{diab}} \chi(t) \\ \chi(t) = & \exp[-(1/\tau_{\text{loc}}) \int_{t_0}^t f(r(\tau')) d\tau'] \end{aligned} \quad (2)$$

where the decay factor  $\chi(t)$  describes a smooth transition to equilibrium within the time interval  $\tau_{\text{loc}}$ . The quantities  $r(t)$  and  $t_0$  denote, respectively, the trajectory and a time well before the collision. The additional repulsive po-

tential  $(\Delta U)_{\text{diab}}$  is that of ref<sup>2</sup>). The adiabatic potential  $U_{\text{ad}}$  is described in terms of the three shape parameters  $r$ ,  $\alpha$ , and  $\tau$ , and is adapted<sup>5</sup>) from the model of Möller and Nix<sup>6</sup>). In order to account for a consistent treatment of the diabaticity at the initial stage of the collision, the transport coefficients are multiplied by the factor  $[1-\chi(\tau)]$ . The diffusion coefficient for the mass asymmetry  $D_{\text{aa}}$ , however, is not modified to allow for the fact that the diabaticity does not inhibit the exchange of particles due to the residual interaction.

For the numerical treatment of the Fokker-Planck equation we use the moment expansion up to second order. For small impact parameters, however, it is found that the method of moment expansion breaks down. Such a difficulty in employing the moment expansion has recently been confirmed by Nix<sup>7</sup>). A critical limit for the validity of the moment expansion can be expressed through the variance for the distance between the surfaces of the two nuclei,  $\sigma_{\text{SS}}^2 < 1 \text{ fm}^2$ . After a trajectory attains  $\sigma_{\text{SS}}^2 > 0.3 \text{ fm}^2$  we, therefore, perform the calculations within an approximation<sup>8</sup>) justified for creeping motions. This approximation, however, does not provide any information regarding the distribution in the  $s$  variable. In figs. 1 and 2 we show, as representative examples, the results for the

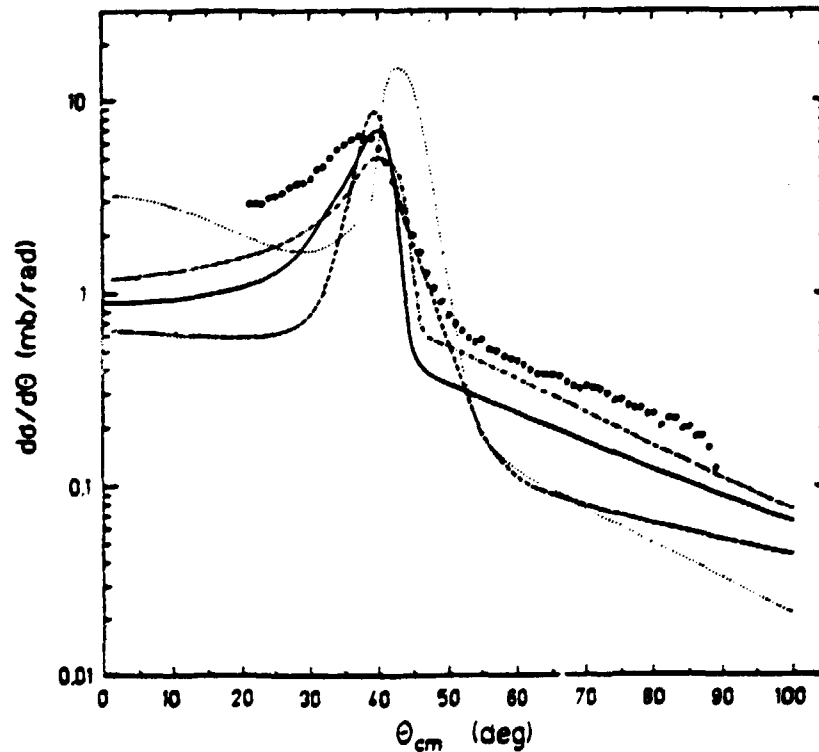
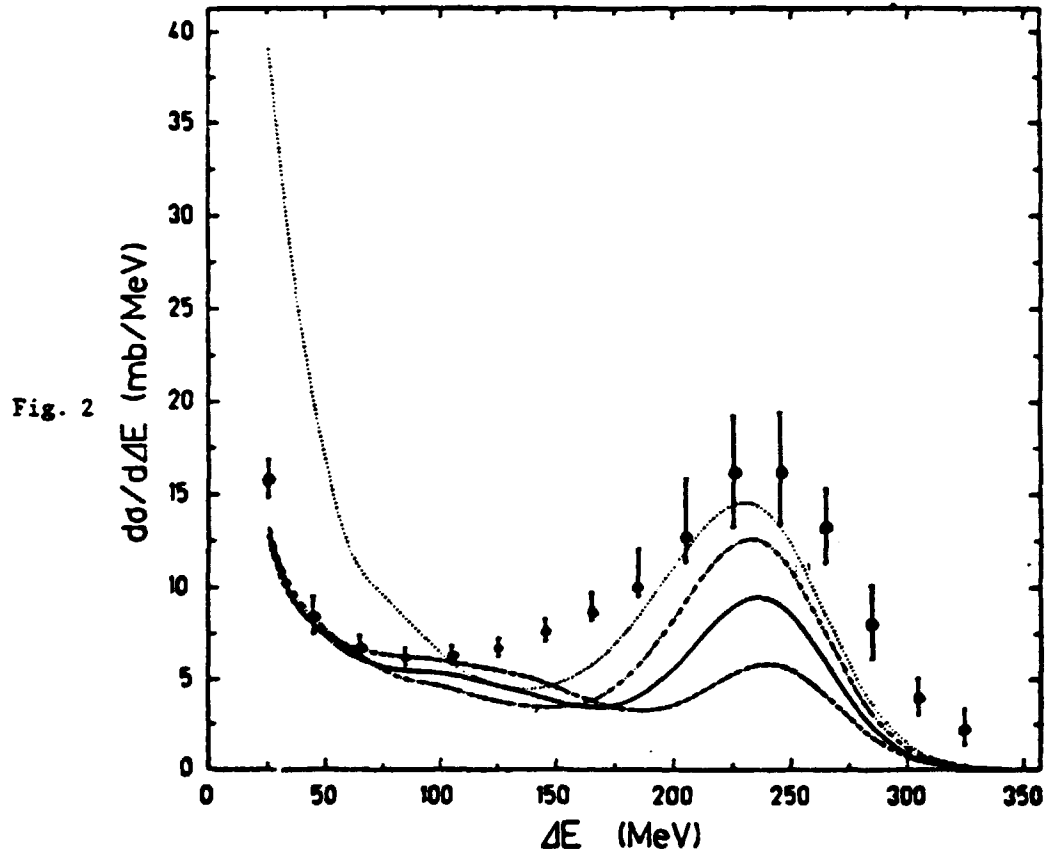


Fig. 1

angular and energy distributions for the reaction  $^{86}\text{Kr}(8.18 \text{ MeV/u}) + ^{166}\text{Er}$  calculated with different values of the local equilibration time ( $\tau_{\text{loc}} \cdot 10^{22} = 0.001(\dots)$ ,  $5(-\cdot-\cdot-)$ ,  $10(----)$ , and  $20(- - -)$ ; experimental data<sup>1,2</sup> are shown by thick dots). It is evident that the results show a remarkable dependence on the size of  $\tau_{\text{loc}}$ . This conclusion holds good for other cross sections as well.



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