

MHD Slip Flow past a Shrinking Sheet

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ABSTRACT

An analysis is made for the steady two-dimensional laminar boundary layer flow of a viscous, incompressible, electrically conducting fluid near a stagnation point past a shrinking sheet with slip in the presence of a magnetic field. The governing boundary layer equations are transformed to ordinary differential equations by taking suitable similarity variables and solved numerically by Shooting method. The effects of the various parameters such as velocity ratio parameter, slip parameter, Prandtl number, Eckert number and magnetic parameter for velocity and temperature distributions have been discussed in detail through graphical representation.

Keywords: MHD; Boundary Layer; Slip Flow; Shrinking Sheet

1. Introduction

The viscous flow and heat transfer in the boundary layer region due to a stretching sheet has wide theoretical and technical applications in manufacturing process and in industries such as extraction of polymer sheet, paper production, wire drawing, glass-fiber production, the cooling and drying of paper and textiles. The study of heat transfer and flow field is necessary for determining the quality of the final products of such process. Crane [1] was first who considered steady boundary layer flow of a viscous incompressible fluid over a linearly stretching plate and gave an exact similarity solution in closed analytical form. The effects of heat and mass transfer and magnetic field under various physical conditions have been investigated by several authors such as Chen and Char [2], Chiam [3], Andersson [4], Ariel *et al.* [5], Jat and Chaudhary [6,7], Wang [8], Fang *et al.* [9], Nadeem *et al.* [10], Bhattacharyya and Layek [11] and recently Bhattacharyya *et al.* [12] studied the slip effects on boundary layer stagnation-point flow and heat transfer towards a shrinking sheet.

Based on the above mentioned investigations and applications, this paper is concerned with a steady, two dimensional stagnation flow of an electrically conducting fluid past a shrinking sheet in the presence of a magnetic field. The results of velocity and temperature distributions for different parameters such as the velocity ratio parameter, the slip parameter, the Prandtl number, the Eckert number and the magnetic parameter were obtained.

2. Formulation of the Problem

Consider the steady two-dimensional laminar flow of a viscous incompressible electrically conducting fluid towards a linearly shrinking sheet such that the sheet is shrunk in its own plane with velocity proportional to the distance from the stagnation point in the presence of an externally applied normal magnetic field of constant strength H_0 . The shrinking surface has a linear velocity $u_w(x)$ and uniform temperature T_w , while the velocity of the flow external to the boundary layer is $u_e(x)$ and temperature T_∞ as shown in **Figure 1**. Therefore, under the usual boundary layer approximations, the governing equations of motion are:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \tag{1}$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \tag{2}$$

$$= u_e \frac{du_e}{dx} + \nu \frac{\partial^2 u}{\partial y^2} - \frac{\sigma_e \mu_e^2 H_0^2 u}{\rho}$$

$$\rho C_p \left(u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) \tag{3}$$

$$= \kappa \frac{\partial^2 T}{\partial y^2} + \mu \left(\frac{\partial u}{\partial y} \right)^2 + \sigma_e \mu_e^2 H_0^2 u^2$$

where ν is the coefficient of kinematic viscosity, σ_e the electrical conductivity, μ_e the magnetic permeability, ρ the density, C_p the specific heat at constant pressure, κ the thermal conductivity and μ the coefficient of viscosity. The other symbols have their usual

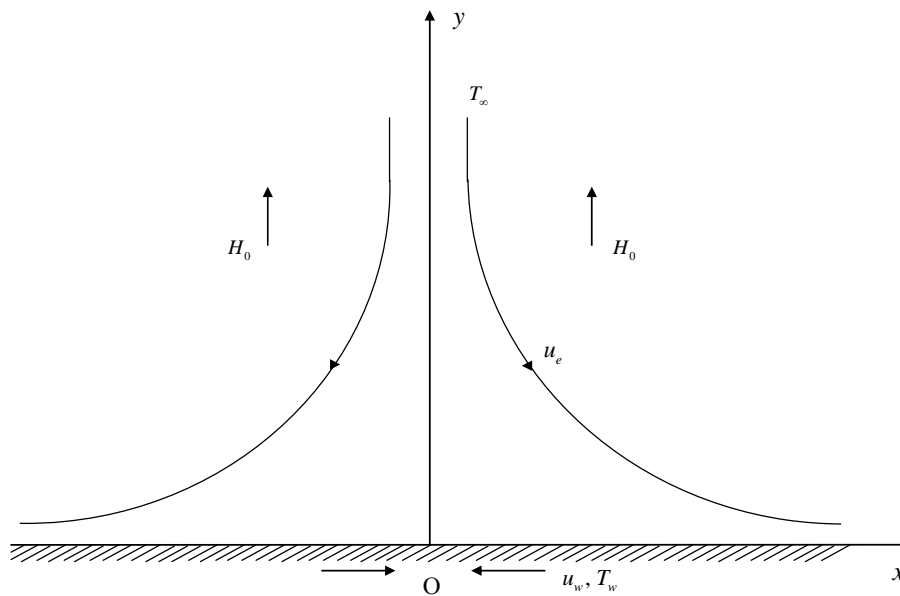


Figure 1. Physical model of two-dimensional stagnation point flow past a shrinking sheet.

meanings.

The boundary conditions are:

$$y = 0 : u = u_w = cx + L \left(\frac{\partial u}{\partial y} \right), v = 0; T = T_w \quad (4)$$

$$y = \infty : u = u_e(x) = ax; T = T_\infty$$

where c is a proportionality constant of the velocity of shrinking sheet, L is a slip length and a is a constant proportional to the free stream velocity for away from the sheet.

3. Analysis of the Velocity and the Thermal Boundary Layers

The continuity Equation (1) is identically satisfied by stream function $\psi(x, y)$, defined as

$$u = \frac{\partial \psi}{\partial y}, \quad v = -\frac{\partial \psi}{\partial x} \quad (5)$$

For the solution of the momentum and the energy Equations (2) and (3), the following dimensionless variables are defined:

$$\psi(x, y) = \sqrt{av}xf(\eta) \quad (6)$$

$$\eta = \sqrt{\frac{a}{\nu}}y \quad (7)$$

$$\theta(\eta) = \frac{T - T_\infty}{T_w - T_\infty} \quad (8)$$

Equations (5) to (8), transform Equations (2) and (3) into

$$f''' + ff'' - f'^2 - \text{Re}_m^2 f' + 1 = 0 \quad (9)$$

$$\theta'' + \text{Pr} f \theta' + \text{Pr} Ec f'^2 + \text{Pr} Ec \text{Re}_m^2 f'^2 = 0 \quad (10)$$

where a prime (') denotes differentiation with respect to

η , $\text{Re}_m = \mu_e H_0 \sqrt{\frac{\sigma_e}{\rho a}}$ is the Magnetic parameter,

$\text{Pr} = \frac{\mu C_p}{\kappa}$ is the Prandtl number and $Ec = \frac{u_e^2}{C_p (T_w - T_\infty)}$

is the Eckert number.

The corresponding boundary conditions are:

$$\eta = 0 : f = 0, f' = \frac{c}{a} + \delta f''; \theta = 1 \quad (11)$$

$$\eta = \infty : f' = 1; \theta = 0$$

where $\frac{c}{a}$ is the velocity ratio parameter and

$\delta = L \left(\frac{a}{\nu} \right)^{\frac{1}{2}}$ is the slip parameter.

For numerical solution of the Equations (9) and (10), we apply the following power series in a small magnetic parameter Re_m^2 as:

$$f(\eta) = \sum_{i=0}^{\infty} (\text{Re}_m^2)^i f_i(\eta) \quad (12)$$

$$\theta(\eta) = \sum_{j=0}^{\infty} (\text{Re}_m^2)^j \theta_j(\eta) \quad (13)$$

Substituting Equations (12) and (13) and its derivatives in Equations (9) and (10) and then equating the coefficients of like powers of Re_m^2 , we get the following set of equations:

$$f_0''' + f_0 f_0'' - f_0'^2 + 1 = 0 \quad (14)$$

$$\theta_0'' + \text{Pr } f_0 \theta_0' = -\text{Pr } Ec f_0''^2 \tag{15}$$

$$f_1''' + f_0 f_1'' - 2f_0' f_1' + f_0'' f_1 = f_0' \tag{16}$$

$$\begin{aligned} \theta_1'' + \text{Pr } f_0 \theta_1' \\ = -\text{Pr } f_1 \theta_0' - \text{Pr } Ec (2f_0'' f_1'' + f_0' f_1'^2) \end{aligned} \tag{17}$$

$$\begin{aligned} f_2''' + f_0 f_2'' - 2f_0' f_2' + f_0'' f_2 \\ = -f_1 f_1'' + (f_1' + 1) f_1' \end{aligned} \tag{18}$$

$$\begin{aligned} \theta_2'' + \text{Pr } f_0 \theta_2' = -\text{Pr} (f_1 \theta_1' + f_2 \theta_0') \\ - \text{Pr } Ec (2f_0'' f_2'' + f_1''^2 + 2f_0' f_1') \end{aligned} \tag{19}$$

with the boundary conditions:

$$\begin{aligned} \eta = 0: f_i = 0, f_0' = \frac{c}{a} + \delta f_0'', f_j' = \delta f_j''; \theta_0 = 1, \theta_j = 0 \\ \eta = \infty: f_0' = 1, f_j' = 0; \theta_i = 0 \quad i \geq 0, j > 0 \end{aligned} \tag{20}$$

4. Results and Discussion

The Equation (14) is that obtained by Bhattacharyya *et al.* [12] for the non-magnetic case and the remaining equations are ordinary linear differential equations and have been solved numerically by Newton's shooting method with fourth-order Runge-Kutta integration scheme for various values of the parameters. The velocity and temperature distributions for various values of parameters are shown in **Figures 2-6** respectively.

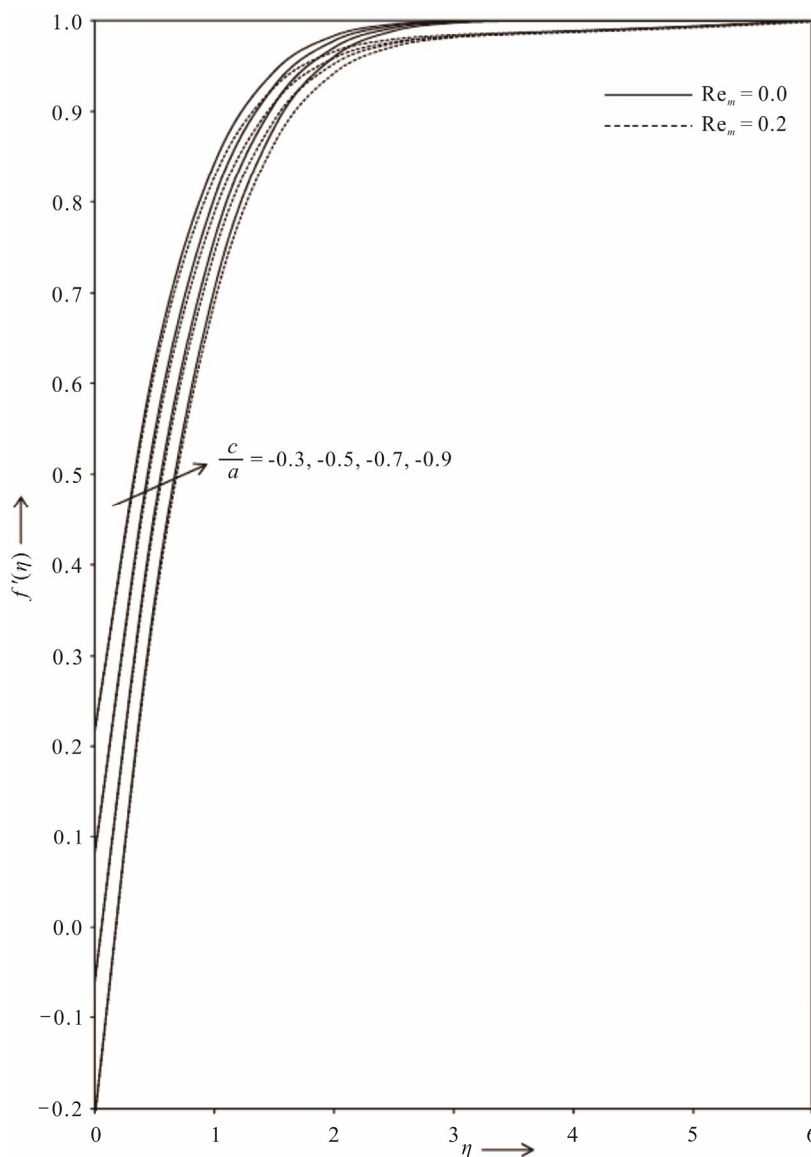


Figure 2. Velocity distribution against η for various values of $\frac{c}{a}$ and Re_m with $\delta = 0.5$.

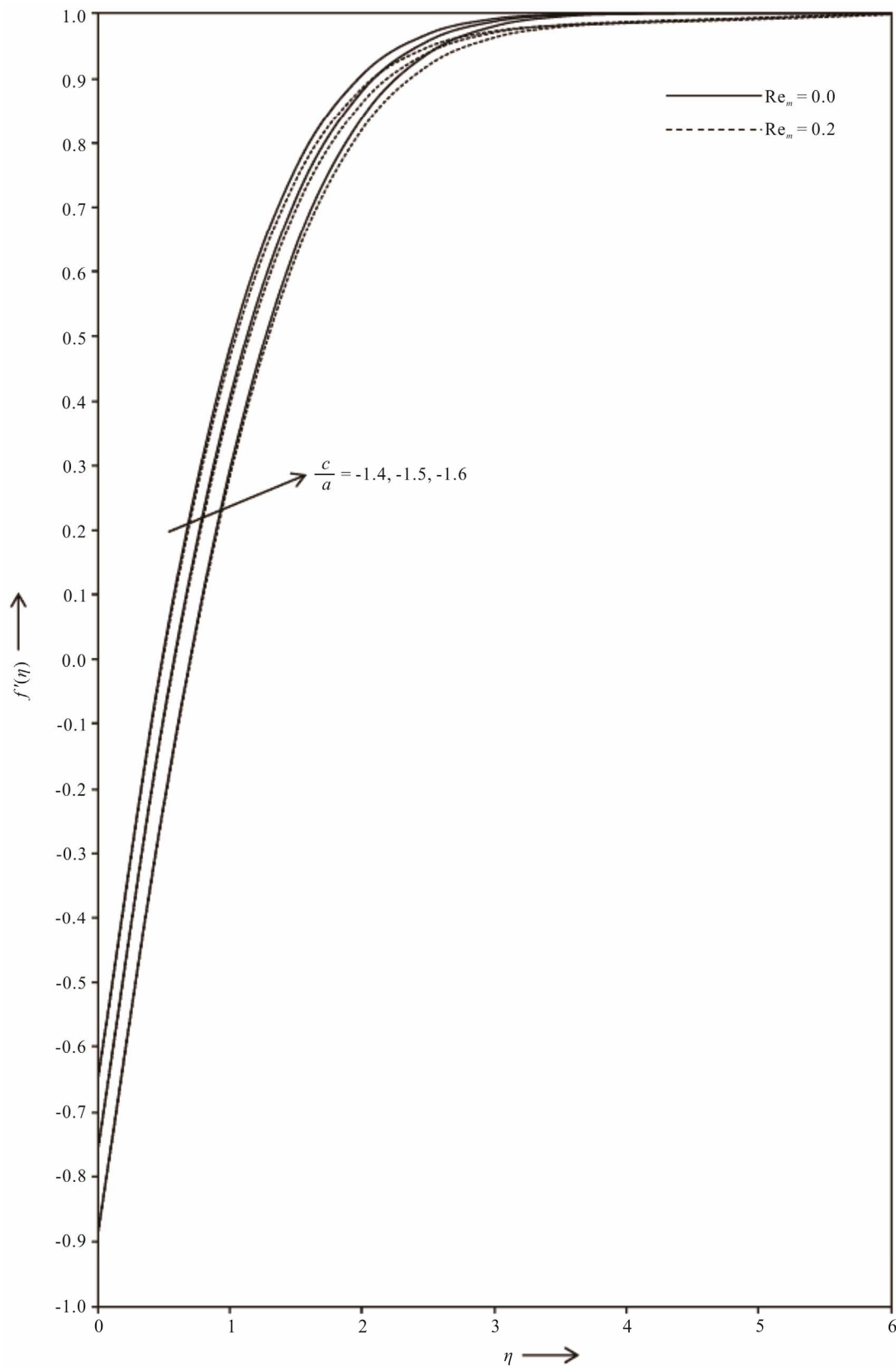


Figure 3. Velocity distribution against η for various values of $\frac{c}{a}$ and Re_m with $\delta = 0.5$.

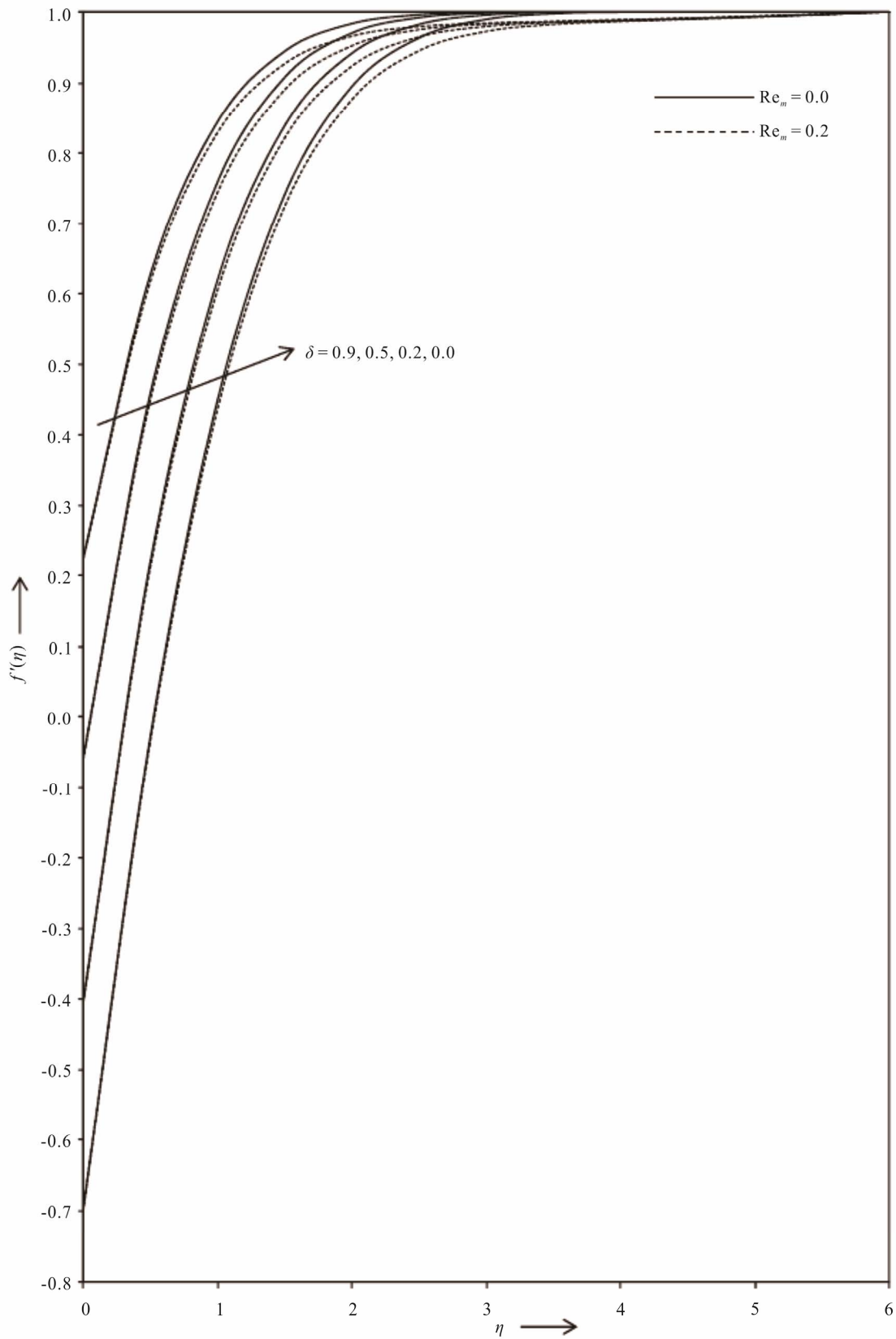


Figure 4. Velocity distribution against η for various values of δ and Re_m with $\frac{c}{a} = -0.7$.

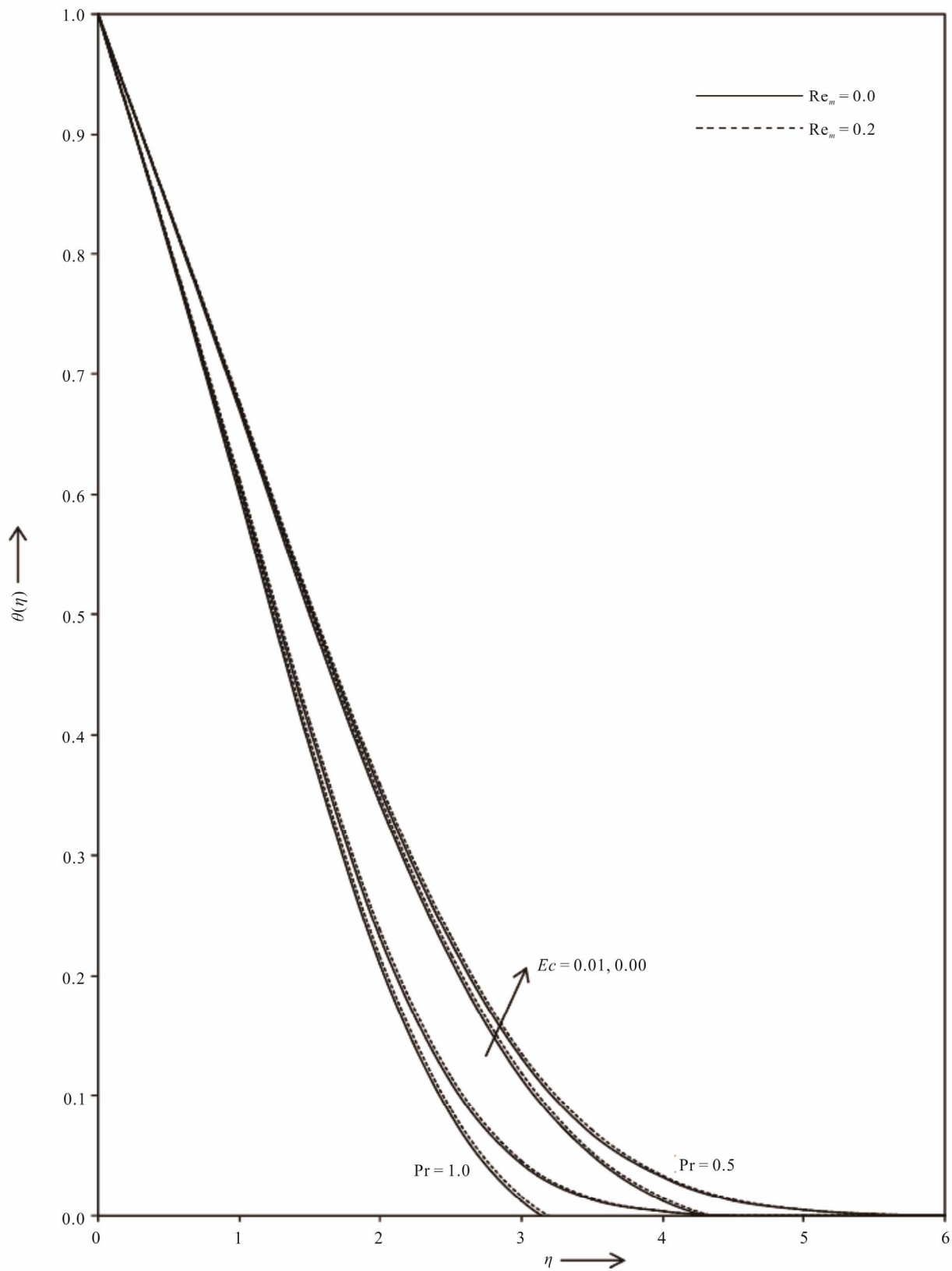


Figure 5. Temperature distribution against η for various values of Pr , Ec and Re_m with $\frac{c}{a} = -1.40$ and $\delta = 0.5$.

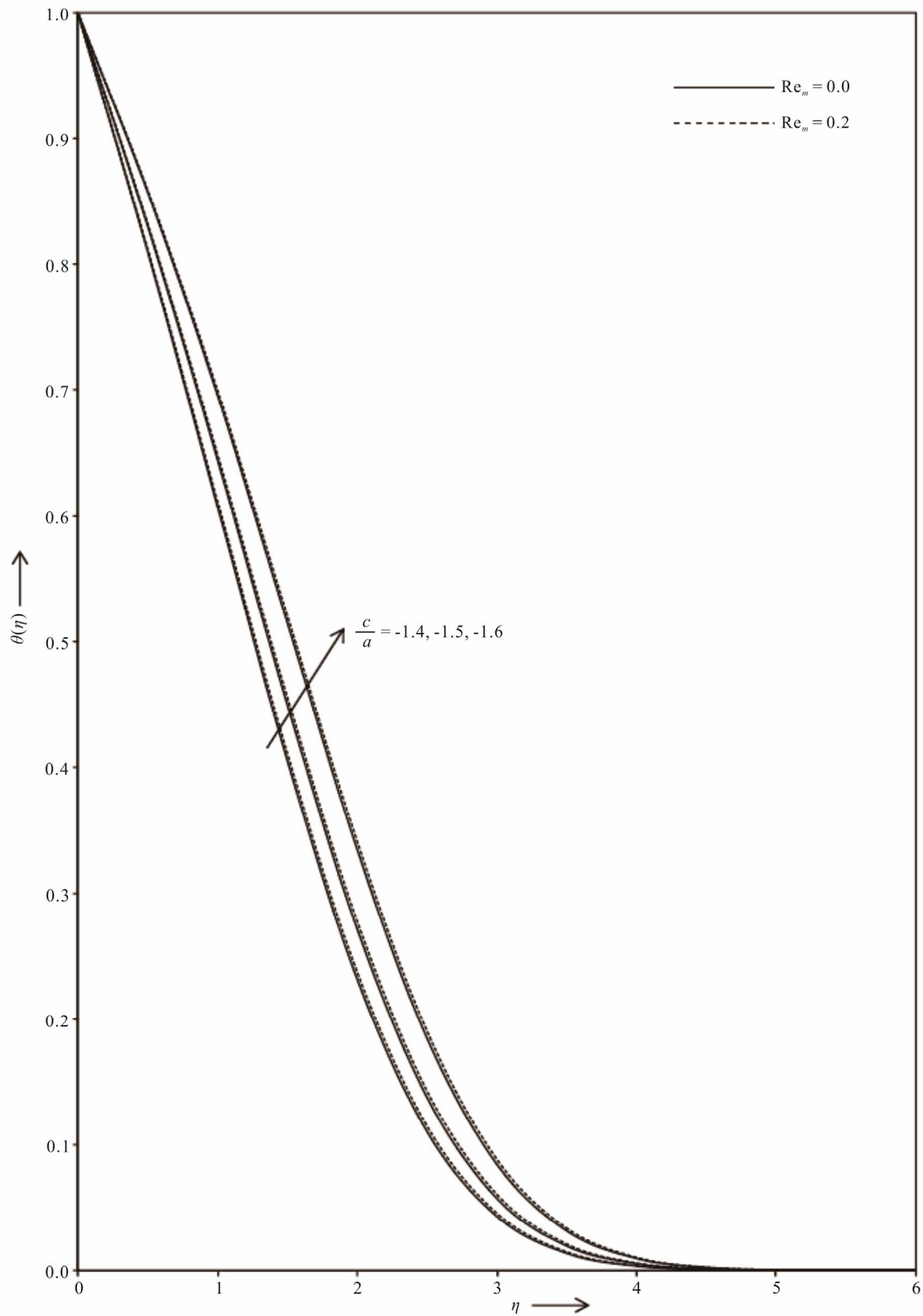


Figure 6. Temperature distribution against η for various values of $\frac{c}{a}$ and Re_m with $\delta = 0.5$, $Pr = 1$ and $Ec = 0.0$.

The velocity profiles $f'(\eta)$ for different values of the velocity ratio parameter $\frac{c}{a}$, the slip parameter δ and the magnetic parameter Re_m are shown in **Figures 2-4**. It is observed that the velocity boundary layer thickness increases with the increasing values of the velocity ratio parameter $\frac{c}{a}$ and the slip parameter δ , whereas it decreases as the magnetic parameter Re_m increases for a fixed η .

The temperature profiles $\theta(\eta)$ for different values of the velocity ratio parameter $\frac{c}{a}$, the slip parameter δ , the Prandtl number Pr , the Eckert number Ec and the magnetic parameter Re_m are plotted in **Figures 5 and 6**. It is observed that for the slip parameter $\delta=0.5$ the thermal boundary layer thickness decreases with the increasing values of the velocity ratio parameter $\frac{c}{a}$, the Prandtl number Pr and the Eckert number Ec and the reverse phenomenon is observed for the magnetic parameter Re_m .

5. Conclusions

In this paper, the stagnation flow for two-dimensional electrically conducting fluid past a shrinking sheet with slip boundaries in the presence of a magnetic field is studied. The similarity equations are derived and solved numerically. It is found that the velocity boundary layer thickness increases with the increasing values of the velocity ratio parameter, the slip parameter. Further we observed that the magnetic parameter decreases with the increasing value of the velocity boundary layer thickness but the reverse phenomenon occurs for thermal boundary layer thickness. Also it observed that the thermal boundary layer thickness decreases with increasing values of the velocity ratio parameter, the Prandtl number and the Eckert number.

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