



Experimental study of the aerosol impact on fog microphysics

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Received: 23 August 2018 – Discussion started: 8 October 2018

Revised: 28 February 2019 – Accepted: 4 March 2019 – Published: 4 April 2019

Abstract. Comprehensive field campaigns dedicated to fog life cycle observation were conducted during the winters of 2010–2013 at the Instrumented Site for Atmospheric Remote Sensing Research (SIRTA) observatory in a suburb of Paris. In order to document their properties, in situ microphysical measurements collected during 23 fog events induced by both radiative cooling and stratus lowering are examined here. They reveal large variability in number, concentration and size of both aerosol background before the fog onset and fog droplets according to the different cases. The objective of this paper is to evaluate the impact of aerosol particles on the fog microphysics. To derive an accurate estimation of the actual activated fog droplet number concentration N_{act} , we determine the hygroscopicity parameter κ , the dry and the wet critical diameter and the critical supersaturation for each case by using an iterative procedure based on the κ -Köhler theory that combines cloud condensation nuclei (CCN), dry particle and droplet size distribution measurements.

Our study reveals low values of the derived critical supersaturation occurring in fog with a median of 0.043 %. Consequently, the median dry and wet activation diameters are 0.39 and 3.79 μm , respectively, leading to a minor fraction of the aerosol population activated into droplets. The corresponding N_{act} values are low, with median concentrations of 53.5 and 111 cm^{-3} within the 75th percentile. The activated fraction of aerosols exhibits remarkably low correlation with κ values, which reflects the chemical composition of the aerosols. On the contrary, the activated fraction exhibits a strong correlation with the inferred critical diameter throughout the field campaigns. This suggests that the variability in the activated fraction is mostly driven by particle size, while variations in aerosol composition are of secondary importance. More-

over, our analysis suggests that the supersaturation reached in fog could be lowered by the aerosol number concentration, which could contribute to the sink term of water vapor during the radiative cooling. Although radiative fogs are usually associated with higher aerosol loading than stratus-lowering events, our analysis also reveals that the activated fraction at the beginning of the event is similar for both types of fog. However, the evolution of the droplet concentration during the fog life cycle shows significant differences between both types of fog.

This work demonstrates that an accurate calculation of supersaturation is required to provide a realistic representation of fog microphysical properties in numerical models.

1 Introduction

As they reduce visibility, fog events strongly perturb the aviation, marine and land transportation. Gultepe et al. (2009) pointed out the extremely high level of human losses and financial cost related to fog and low-visibility events. Moreover, present numerical weather prediction models are usually unable to predict the exact location and time evolution of the life cycle of a fog layer (Zhou and Ferrier, 2008; Van der Velde et al., 2010; Bergot, 2013; Boutle et al., 2018).

The occurrence and development of fog result from the nonlinear interaction of competing radiative, thermodynamic, microphysical and dynamical processes (see the review of Gultepe et al., 2007b). Water droplets are formed by heterogeneous nucleation of aerosol particles when the relative humidity (RH) exceeds 100 %. The Köhler theory (Köhler, 1936) forms the basis of our understanding of cloud

droplet formation. The ability of aerosol particles to act as cloud condensation nuclei (CCN) depends largely on their size, composition and phase state. The number size distribution of the droplets activated during the cloud formation phase depends on the supersaturation values reached by the air mass (Pruppacher et al., 1998). Particles having a critical supersaturation below the maximum value are activated and will then further grow by water vapor diffusion as long as the RH remains high enough, while the others particles remain at their equilibrium diameter. At the same time, water droplets are the particles whose diameters are larger than the critical threshold corresponding to that critical supersaturation.

Because of the weak supersaturation inside the fog layers (Hudson, 1980; Pandis et al., 1990; Svenningsson et al., 1992; Hammer et al., 2014; Boutle et al., 2018), the separation between activated fog droplets and non-activated particles, so-called hydrated aerosols, is not explicit (Frank et al., 1998). To date, the microphysics in numerical weather prediction (NWP) models takes into account activated droplets only. Recently, a two-moment microphysical cloud scheme, which relies on the prognostic evolution of an aerosol population to predict cloud droplets, has been developed for low supersaturation clouds (Thouren et al., 2012) and needs to be tested against observations. It is thus essential to precisely evaluate total droplet number concentration values and their corresponding size range. Currently, their definition mainly deals with the instrumentation limitation that generally begins at 2 μm in diameter to measure droplet size distribution. However, Hammer et al. (2014) and Elias et al. (2015) experimentally determined threshold diameter higher than 2 μm to characterize droplets. As the smallest hydrated particles are more numerous, considering a smallest threshold diameter could lead to a significant overestimation of the droplet concentration. Moreover, as shown by Boutle et al. (2018), an accurate representation of droplet concentration for fog could also impact climate projection.

As a general rule in cloud physics, increase in aerosols causes an increase in the droplet number concentration for a given cloud type (Gultepe and Isaac, 1996; Ramanathan et al., 2001). Hudson (1980) and Liu et al. (2017) found systematic differences in the fog microstructure between fogs formed in maritime, continental and urban air masses with a simultaneous increase of fog condensation nuclei and aerosol concentration. Recent numerical simulations exhibit a strong positive correlation between aerosol and droplet concentrations (Zhang et al., 2014; Stolaki et al., 2015; Maalick et al., 2016). They also predict fog events with high values of droplet concentrations, typically a few hundred per cm^{-3} or more. By contrast, numerical study of Bott (1991) shows that an increase in aerosol particles leads to a decrease of the supersaturation that in turn decreases the activated droplet number.

Recently, Boutle et al. (2018) have obtained impressively low droplet concentrations with large eddy simulation (LES) by numerically solving the condensation equation at every

time step and grid point, but they had very few activable aerosol particle.

In this study, we examined in situ microphysical measurements collected during the three winter campaigns to derive typical fog droplet concentrations in semi-urban conditions. From October 2010 to March 2013, a suite of active and passive remote sensing and in situ sensors was deployed at the Instrumented Site for Atmospheric Remote Sensing Research (SIRTA) in a suburb of Paris (Haefelin et al., 2010). Dupont et al. (2016) analyzed 117 fog events to provide typical values of thermodynamical and radiative variables for the fog formation and the mature and dissipation phases. Burnet et al. (2012) have shown that, in terms of droplet number concentration and effective diameter, the microphysical properties of fog events sampled at SIRTA present a large variability. Elias et al. (2015) examined the data collected during November 2011 and assessed the contribution of hydrated aerosols that led to the extinction of the visible radiation in the mist–fog–mist cycle. Hammer et al. (2014) focused on the activation properties of developed fog events observed during the winter of 2012/2013, but they did not provide activated particle concentrations values.

Here, we determine the hygroscopicity parameter, the dry and wet critical diameters and the critical supersaturation by using an iterative procedure based on the κ -Köhler theory that combines CCN measurements, dry particle size distribution and composite wet particle size distribution at ambient humidity. The objectives are to derive accurate estimations of the fog droplet concentrations and to investigate the impact of aerosols on the fog microphysics. This paper provides a comprehensive data set on activation properties and an explanation of the aerosol impact on the activation process.

Data and methods are described in Sects. 2 and 3, respectively. Results are presented in Sect. 4, with first the statistics of fog activation properties of the 23 fog events analyzed in this study. The link between aerosol particles and fog droplets is examined in Sect. 4.2, and the impact of CCN concentration on fog microstructure is discussed in Sect. 5. Conclusions are finally given in Sect. 6.

2 The data set

2.1 Instrumentation

Data presented here were collected at the SIRTA observatory in the framework of the ParisFog field campaigns (Haefelin et al., 2010). During the winters of 2010 to 2013, specific instrumentation were deployed for the PreViBOSS project (Elias et al., 2012) to provide continuous observation of aerosol and fog microphysics. The experimental setup was already presented in Burnet et al. (2012), Hammer et al. (2014), Elias et al. (2015) and Dupont et al. (2016). The instruments used in this study are listed in Table 1.

Table 1. Instruments deployed during the campaigns that are used for this study.

Instruments	Measured parameters	Time resolution
Palas WELAS-2000	Hydrated and activated aerosol particles' number size distribution $D = 0.39\text{--}42\ \mu\text{m}$	5 min
DMT fog monitor (FM-100)	Droplet number size distribution $D = 2\text{--}50\ \mu\text{m}$	1 min
PVM Gerber	Liquid water content $D = 3\text{--}50\ \mu\text{m}$	1 min
Degreanne DF20+	Horizontal visibility	1 min
TSI SMPS	Dry aerosol particles' distribution $D = 10.6\text{--}493\ \text{nm}$	5 min
CCNC DMT	CCN concentration SS = 0.1–0.5 %	20 min for a cycle
CPC (TSI 3025)	Aerosol particles' concentration $D = 2.5\text{--}2500\ \text{nm}$	5 min

Particle size distribution at ambient humidity is derived from a combination of two optical spectrometers: the WELAS-2000 (Palas GmbH, Karlsruhe, Germany) and the FM-100 (Droplet Measurement Technologies Inc., Boulder, CO, USA), which do not cover the same size range. The WELAS-2000 (hereafter referred to as WELAS) provides particle size spectrum between 0.4 and 40 μm in diameter according to the constructor. However, the detection efficiency decreases drastically below $\sim 1\ \mu\text{m}$ (Heim et al., 2008; Elias et al., 2015), resulting in a strong underestimation of the concentration of the submicronic particles. Hammer et al. (2014) choose to consider only data with diameters larger than 1.4 μm . Statistics over the whole data set reveal that the most frequent mode diameter of the WELAS size distribution is 0.96 μm . Thus, we choose to use this value as the lowest threshold and only measurements for bin diameter larger than 0.96 μm will be considered in this study. Note that the activated diameter in fog is expected to be larger than 1 μm ; thus, this instrumental bias will not affect the results presented here, which is confirmed next in this study for SIRTA's fog. The sampling time period was fixed to 5 min as a compromise between time resolution and statistical significance of the measurements. Indeed, this corresponds to a volume of air of 6.40 cm^{-3} per sample.

The FM-100 provides 1 Hz droplet size distribution from 2 to 50 μm in diameter according to the constructor, hence providing a large overlapping range with the WELAS measurements. However, Burnet et al. (2012) and Elias et al. (2015) show that these distributions only overlap each other at a diameter which fluctuates between 5 and 9 μm . Their comparisons reveal a high discrepancy between both these probes with a large underestimation by the FM-100 for particles lower than about 5 μm but conversely a large underestimation by the WELAS for droplets larger than 10 μm (Burnet

et al., 2012; Elias et al., 2015). This is illustrated in Fig. 1, which shows the size distributions measured by the WELAS (black) and FM-100 (cyan) for two fog cases with contrasting properties. Figure A1 in Appendix A shows comparisons of the particle number concentration values of the FM-100 vs. the integrated WELAS measurement over 2–10 μm diameter range that correspond to the four first bins of the FM-100. The default manufacturer's first four bins of the FM-100 are 2–4, 4–6, 6–8 and 8–10 μm , which correspond to 9, 7, 4 and 3 class bins of the WELAS, respectively. Following the WELAS sampling time period, FM-100 and WELAS data points are a 5 min average that represents 2851 samples for the whole data set. Figure A1 confirms that the FM-100 strongly underestimates the particle counts in the first size bin and that the WELAS underestimates the concentration of particles larger than 8 μm . The measurements of the two instruments do not match perfectly in the size range of 4 and 8 μm , which reflects the large fluctuations of the crossing diameter over this range. To derive a composite size distribution, Elias et al. (2015) chose a constant value of 7 μm . Because of this high variability, we compute here the composite size distribution by using WELAS data up to 6 μm , FM-100 data above 8 μm and the average of both from 6 to 8 μm . Results are illustrated in Fig. 1 with red segments corresponding to the junction between the WELAS and FM-100 distribution. This procedure allows us to take advantage of the finer bin resolution of the WELAS for the smaller particles and also to reduce uncertainties of the FM-100 due to Mie ambiguities for smaller diameter, as described by Spiegel et al. (2012), since our FM-100 was using the default manufacturer's bin threshold. Indeed, particularly for smaller diameter, the droplet concentration may be overestimated or underestimated; a few droplets from adjacent classes can be included or some can be counted in an adjacent class (Gonser

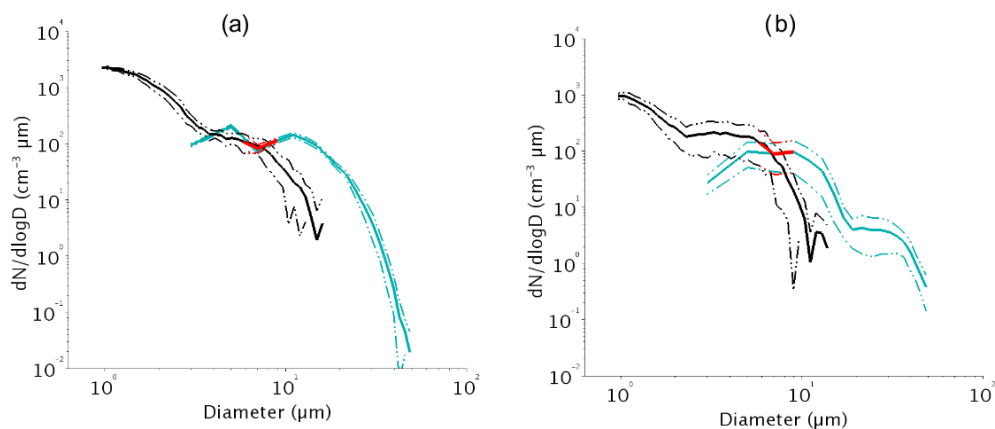


Figure 1. The 1 h average composite number size distributions at ambient humidity derived from WELAS (black) and FM-100 (cyan) at the beginning of the fog event for cases (a) f6 and (b) f20. The solid line represents mean value and the dotted lines indicate ± 1 standard deviation. The red segments correspond to the result of merging both distributions between 6 and 8 μm .

et al., 2012). Nevertheless, because of the large discrepancies between the two sensors on the crossing diameter, the next calculation using this crossing diameter also considers a standard deviation using the minimum and the maximum size distributions over this size range. Both instruments are located on a scaffolding about 2.5 m high, close to a PVM-100 from Gerber Scientific Inc. used as a reference for the liquid water content (hereafter LWC) measurements, which measures on the 2–50 μm diameter. Measurements of visibility and its vertical evolution are given by two Degreanne difusometers (DF20+ and DF320) located at 4 and 18 m above ground, respectively. Confidence is given in FM-100 measurements by comparison of the integrated LWC over its size range with the LWC measured by the PVM-100 and the visibility trend; see Burnet et al. (2012) for winter 2010–2011. A paper will follow on that issue.

Aerosol particle measurements are performed under dry conditions by instruments placed in a shelter. The sampled air mass enters through an aerodynamic size discriminator $\text{PM}_{2.5}$ inlet and a dryer which reduces the relative humidity to less than 50%. This value could be higher than the efflorescence point for some particles (Tang et al., 1995). A scanning mobility particle sizer (SMPS), which provides the dry aerosol particle number size spectrum, consists of a differential mobility analyzer (DMA; TSI 3071) which selects particles from 10.6 to 496 nm and of a condensation particle counter (TSI CPC 3022). Hydrophilic filters in the DMA circuit should lower the relative humidity to less than 30% (Denjean et al., 2014), which should be enough to avoid water uptake issues. Another CPC (TSI 3025) measures the total particle number concentration from 2.5 nm to 2.5 μm . Finally, a continuous-flow streamwise thermal gradient CCN chamber (CCNC) (Roberts and Nenes, 2005) is used to measure the CCN number concentration at five different supersaturations from 0.1 to 0.5%, by steps of 0.1%.

2.2 Selected fog cases and aerosol properties

During the three winter campaigns of 2010–2013, Dupont et al. (2016) report the occurrence of 117 fog events. However, due to instrument failures and technical difficulties in operating the whole set of instruments on a 24/7 mode, 42 events were sampled simultaneously with both the WELAS and the FM-100, and only a subset of 23 cases was also sampled with both the CCNC and the SMPS. They are listed in Table 2 with their classification type (RAD for radiative cooling fog and STL for stratus-lowering ones), as determined by the scheme of Tardif (2007) and their vertical development based on the comparison of both difusometers: a developed fog produces low-visibility conditions simultaneously at 4 and 18 m, while a thin fog produces low-visibility conditions at 4 m only (Elias et al., 2009; Dupont et al., 2016). About the same proportion (40%) of RAD and STL fog events occurred at the SIRTAs site (Haeffelin et al., 2010; Dupont et al., 2016). Among the 23 fog events analyzed here, 13 are radiation fogs and 10 stratus-lowering fogs, 19 are developed and 4 thin. Thermodynamics conditions during the field campaigns are described by Dupont et al. (2016).

To characterize the aerosol background prior to a fog event, statistics of the total number aerosol particle concentration, N_a , as measured by the SMPS between 10 and 496 nm diameter, are computed over the last hour before the beginning of the fog event. This time lapse is chosen to characterize aerosol before a fog event because CCNC, which allows to characterize aerosol chemistry, has a sampling time of 20 min; thus, a 1 h time lapse is supposed to provide reliable measurements. Median, 25th and 75th percentiles are reported in Table 2. N_a values range from ≈ 2000 to $\approx 20\,000\text{ cm}^{-3}$. Consistent with Haeffelin et al. (2010), the smallest values are observed in westerly flow conditions, while highest values are associated mostly with an easterly flow when SIRTAs is exposed to continental conditions. In-

Table 2. List of fog events analyzed here. Type RAD corresponds to radiation fog and STL to stratus lowering, “thick” to fog developed on the vertical and “thin” to fog layer with top altitude lower than 18 m. N_a is the number concentration of aerosol particles derived from SMPS data. The 25th, 50th and 75th percentiles are computed from the distribution of 5 min samples over the last hour before the fog beginning. N_{FM} is the droplet number concentration as measured by the FM-100 over the range 2–50 μm . Statistics of N_{FM} are computed with minute average data of samples with LWC > 0.005 g m^{-3} collected during the whole fog event. The last column provides the minimum of visibility observed during the event.

Fog ID	Start time (dd/mm/yyyy) (UTC)	End time (UTC)	Type		N_a (cm^{-3})			N_{FM} (cm^{-3})			VISI _{min} (m)
			(1)	(2)	25th	50th	75th	25th	50th	75th	
f1	16/11/2010 22:00	17/11/2010 05:40	RAD	thick	18 868	20 175	20 625	52	67	80	53
f2	19/11/2010 05:40	19/11/2010 10:10	STL	thick	5877	6235	6783	30	51	73	108
f3	19/11/2010 15:40	19/11/2010 17:50	STL	thick	3540	4296	5213	5.3	6.3	7.4	389
f4	10/11/2011 18:00	11/11/2011 17:30	RAD	thick	10 723	12 602	13 086	5 4	89	114	122
f5	15/11/2011 02:30	15/11/2011 09:40	RAD	thick	5703	6043	6285	45	60	75	235
f6	16/11/2011 01:10	16/11/2011 13:30	RAD	thick	7391	8188	9087	103	127	145	128
f7	16/11/2011 16:00	17/11/2011 00:10	STL	thick	7213	7648	8310	129	147	160	112
f8	18/11/2011 01:30	18/11/2011 04:10	RAD	thin	5245	5396	5447	71	85	112	167
f9	19/11/2011 22:00	20/11/2011 08:30	RAD	thin	13 094	13 587	14 135	75	110	150	90
f10	21/11/2011 23:50	22/11/2011 08:10	RAD	thin	10 160	10 526	10 737	87	123	156	87
f11	22/11/2011 20:50	22/11/2011 22:30	RAD	thin	15 209	16 290	19 031	105	132	164	116
f12	23/11/2011 03:25	23/11/2011 10:05	RAD	thick	8920	9320	9723	87	115	157	115
f13	24/11/2011 06:20	24/11/2011 14:00	STL	thick	4770	5258	5586	18	33	55	193
f14	24/11/2011 16:10	24/11/2011 18:15	STL	thick	4180	4340	4566	19	33	48	267
f15	25/11/2011 21:40	26/11/2011 10:30	STL	thick	3328	3618	3762	45	68	121	106
f16	28/11/2011 06:30	28/11/2011 10:40	RAD	thick	9308	9618	11 258	13	19	24	195
f17	16/11/2012 20:45	17/11/2012 09:20	STL	thick	9145	10 051	11 959	31	39	49	151
f18	20/11/2012 03:00	20/11/2012 09:10	RAD	thick	3108	3168	3410	16	26	44	177
f19	20/11/2012 20:15	20/11/2012 22:50	STL	thick	3198	3567	4282	39	62	84	139
f20	22/11/2012 03:15	22/11/2012 09:10	RAD	thick	4673	4780	4827	33	46	56	144
f21	30/11/2012 19:00	01/12/2012 02:45	RAD	thick	17 791	18 712	19 151	42	63	87	102
f22	10/01/2013 02:15	10/01/2013 03:36	STL	thick	1975	2098	2273	21	36	64	233
f23	10/01/2013 05:00	10/01/2013 06:40	STL	thick	1904	1955	2009	24	29	33	351

deed, SIRTa is located 25 km southwest of Paris in a semi-urban environment, composed of agricultural fields, wooded areas, housing and industrial developments (Haefelin et al., 2005), and is exposed to air mass charged with pollution originating from regional background according to Crippa et al. (2013).

These various conditions are reflected on the statistics that reveal a large case-to-case variability. This is illustrated in Fig. 2, which shows the scatterplot of the droplet number concentration derived by the FM-100, N_{FM} on 2–50 μm diameter range as a function of N_a , that represent the aerosol loading, for the 23 fog events (data in Table 2). Symbols depend on the fog type: blue and red colors for STL and RAD fogs, respectively, and open and solid symbols for thin and developed fogs, respectively. For STL fogs, N_a ranges from 2000 to 11 000 cm^{-3} with a median value of 4340 cm^{-3} and 25th and 75th percentiles of 2833 and 6942 cm^{-3} , respectively, while for RAD cases, N_a values are spread over a larger range from 3000 to 20 000 cm^{-3} with a median value of 8822 cm^{-3} and 25th and 75th percentiles of 5719 and 13 094 cm^{-3} , respectively. Thus, over this 3-year campaign,

it appears that at SIRTa STL fogs are associated with lower aerosol loading than RAD ones, as already reported by Elias et al. (2015) for the month of November 2011.

There is also a significant difference between a STL and RAD fog with respect to N_{FM} . Despite the maximum value being reached for a STL fog (147 cm^{-3} for the case of f7), N_{FM} for a STL fog is generally lower, with median values between 38 and 83 cm^{-3} and 25th–75th percentiles of 31–63 and 53–118 cm^{-3} , for STL and RAD, respectively. Moreover, for a RAD fog, thin cases exhibit higher values of N_{FM} with a median and 25th–75th percentiles of 97 and 73–131 cm^{-3} compared to 64 and 49–81 cm^{-3} for a developed fog. Therefore, it appears that a radiative fog is also associated with higher droplet concentrations than stratus-lowering fog over the range 2–50 μm as measured by the FM-100. One can note that these values are rather low for continental conditions (Wendisch et al., 1998; García-García et al., 2002; Gultepe and Milbrandt, 2007a; Niu et al., 2012; Price, 2011; Liu et al., 2011; Lu et al., 2013; Zhao et al., 2013; Boutle et al., 2018).

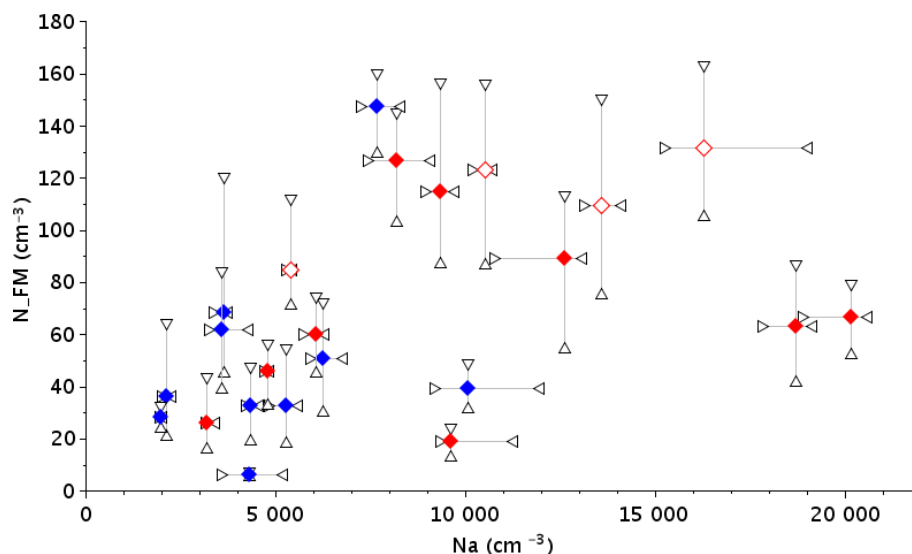


Figure 2. Droplet number concentration as derived from the FM-100 collected during the whole fog event as a function of the aerosol number concentration as derived from the SMPS measurements collected over the last hour before the fog beginning. Median values (diamonds) are indicated for each case of Table 2, with the 25th and 75th percentiles (triangles) to represent the variability. Symbols depend on the fog types, STL (blue) and RAD (red), with open and solid symbols for thin and developed fogs, respectively.

As shown in Fig. 2, the general trend points towards a slight increase in the number of droplets with the aerosol loading but the scatter is very large. In cloud physics, the connection between an increase in aerosol particles and an increase in the cloud droplet number concentration (CDNC) has been supported by many in situ observations (Twohy et al., 2005; Lu et al., 2007; Levin and Brenguier, 2009) and even if the discrepancy is large in the compilation of the diverse results (Ramanathan et al., 2001), the general trend is much more pronounced to what is presently observed in Fig. 2. However, CDNC is derived here from FM-100 measurements only between 2 and 50 μm diameter. In warm clouds, the supersaturation at the cloud base is high enough to activate droplets whose size will increase by water vapor condensation during the ascent of the air parcel. The spectrum then gets narrower because the growth rate of a droplet is inversely proportional to its size, while the remaining interstitial particles keep their own equilibrium diameter at 100 % relative humidity. Higher up in the cloud, the spectrum is widened by different processes such as turbulent mixing and collision–coalescence and new CCN can also be activated. Droplet spectra vary considerably in space and in time but the general trend is that the activated droplet population is clearly separated from the interstitial non-activated particles. See, for example, measurements into stratocumulus cloud as reported in Martin et al. (1994), Brenguier et al. (2011) and Ditas et al. (2012). Measurements by optical counters such as FSSP (Forward Scattering Spectrometer Probe) or CDP (Cloud Droplet Probe) over the range (about 2–50 μm) are considered to provide an accurate estimation of the droplet size distribution.

In contrast, supersaturation in fog is much lower and, as already pointed out by Hudson (1980), in continental air the droplet number size distribution does not exhibit a clear separation. This is clearly illustrated in Fig. 1. To derive an accurate estimate of the CDNC, it is thus essential to determine carefully the wet activation diameter in order to integrate the composite size distribution derived from both WELAS and FM-100 measurements. The method used to discriminate the droplets from the hydrated aerosol particles using κ -Köhler theory is described in the next section.

3 Methods

There are different methods to separate fog droplets from non-activated aerosol particles. A fixed value can be used as a rough estimate, as described in previous studies by Noone et al. (1992), Hoag et al. (1999) and Elias et al. (2009), among others, with thresholds ranging from 2.5 to 5 μm . To take into account case-to-case variability, Elias et al. (2015) used WELAS measurements. The volume size distribution is fitted with two log-normal distributions and the transition diameter is defined as the intersection between them. Hammer et al. (2014) investigated the activation properties by measuring the total and interstitial dry particle number size distribution behind two different inlet systems, and with WELAS and CCN measurements. They compared two methods: one by fitting surface distributions of SMPS+WELAS measurements similar to Elias et al. (2014) and one by retrieving the dry activation diameter from the difference between interstitial and total particle size distribution measurements.

In this study, we use a different approach based on an iterative procedure that combines dry particle distribution from SMPS and composite wet particle size distribution derived from both WELAS and FM-100 measurements. The CCN measurements are also used to derive the hygroscopicity parameter (κ) needed to link the dry and wet activation diameters (D_d and D_w) with the critical supersaturation peak (SS_{peak}) using the κ -Köhler theory.

3.1 κ -Köhler theory

The Köhler theory (Köhler, 1936) expresses the equilibrium saturation vapor pressure over a solution droplet considering the solute effect (Raoult) and water surface tension effect (Kelvin). Accurate information is needed on the particle dry diameter and about the chemical composition to determine whether or not it can act as a CCN. Recently, Petters and Kreidenweis (2007) developed a method, named κ -Köhler, to describe the relationship between particle dry diameter and CCN using a single hygroscopic parameter, κ . This method allows to study the activation process without considering aerosols' complex chemical composition (McFiggans et al., 2006). The formulation of the method is expressed in Eq. (1):

$$S(D) = \frac{D^3 - D_d^3}{D^3 - D_d^3(1 - \kappa)} \exp\left(\frac{4\sigma_{s/a}M_w}{RT\rho_w D}\right), \quad (1)$$

where S is the saturation ratio over a solution droplet, D is the droplet diameter, D_d is the dry droplet diameter, ρ_w is the density of water, M_w is the molar mass of water, $\sigma_{s/a}$ is the surface tension of the solution air/interface (of pure water here), R is the universal gas constant, T is the temperature, and κ is the hygroscopicity parameter. κ represents a quantitative measure of aerosol particles water uptake characteristics and CCN activity. The critical supersaturation SS_c and critical wet activation diameter D_w correspond to values at the maximum supersaturation. They are linked to a couple (D_d , κ) through the following relation:

$$\kappa = \frac{4A^3}{27D_d^3 \ln^2 SS_c} \quad A = \frac{4\sigma_{s/a}M_w}{RT\rho_w}. \quad (2)$$

SS_{peak} represents the maximum supersaturation that the air mass experienced for a sufficiently long time (Hammer et al., 2014); all particles whose SS_c is less than SS_{peak} are activated and will further grow by water vapor condensation as long as RH remains high enough, while other particles remain stable at their equilibrium diameter at the actual RH value.

3.2 Method to determine fog activation properties

The iterative procedure used to derive fog activation properties is illustrated in Fig. 3. To derive D_d and D_w simultaneously, measured CCN concentration (N_{ccn}), dry and wet aerosols' size distributions were used. From that, it follows

that all particles become activated to droplets when their dry or wet size is similar to or larger than D_d and D_w . Indeed, there can only exist one trio (N_{ccn} , D_w and D_d) that is linked to κ and SS_{peak} through Eq. (2). For each fog event, the activation properties are determined as follows: N_{ccn} is set corresponding to a value of D_d . D_w and SS_{peak} are further calculated by using Eq. (2) for SS_{peak} and by numerically searching for the maximum of Eq. (1) for D_w for a given value of κ . Then the integral of the droplet distribution from D_w provides the droplet concentration N_d . Iterations are made on N_{ccn} until N_{ccn} equals N_d : this is the activated aerosol number concentration for this case, hereafter referred to as N_{act} . N_{act} corresponds to the concentration of activated aerosols at SS_{peak} . Two sets of iterations are made with the two κ values and the average particle size distribution. To take into account the size distribution variability within the selected time period, two other sets of iterations are made by using the average size distribution ± 1 standard deviation corresponding to the extrema of the WELAS and FM-100 crossing diameter, respectively, and with the κ value maximizing the scatter.

Our method to determine fog activation properties has to assume that the aerosol particles are internally mixed. However, atmospheric aerosols are frequently externally mixed with particles of different sizes exhibiting different chemical composition. Discussion on this hypothesis is presented in Sect. 3. Moreover, surface tension properties (surface tension) of atmospheric aerosols could be modified by surfactant. This effect has largely been attributed to organics when the organic concentration in the aerosol population is sufficiently large (Facchini et al., 2000; Noziere et al., 2014). The extent to which this will occur is strongly dependent upon the specific molecular properties of the organic molecules. Unfortunately, exploring these aspects would need to combine CCN measurements and a detailed analysis of the molecular species contained in organic particles. This is out of the scope of this paper. In addition, number concentration of aerosol particles larger than 496 nm was assumed to be negligible in our calculation.

The air mass is sampled through a PM_{2.5} head and is dried to RH < 50 % before entering the SMPS and the CCN chamber. Hence, during a fog event, activated aerosols larger than 2.5 μm are missed. Aerosol properties of the air mass must thus be characterized before the occurrence of a fog event. In addition, the aerosol particle distribution changes continuously according to the boundary layer evolution and aerosol sources (wood burning, road traffic, etc.). This is illustrated in Fig. A2, showing a large variability in the aerosol particle distribution along the diurnal cycle for the fog event f6. There is a large variability along the diurnal cycle. To characterize the air mass just before the fog event, we average data recorded during the last hour before the fog onset. This corresponding time interval is delimited by red segments in Fig. A2.

CCN chamber measurements are used to derive κ by using Eq. (2) knowing SS_c and D_d . The chamber supplies the con-

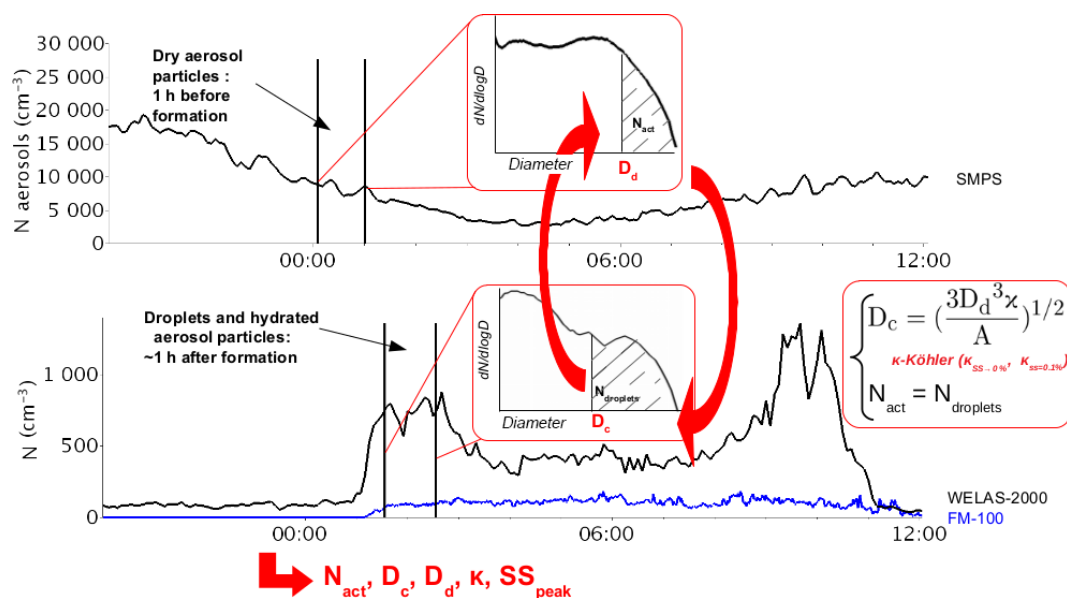


Figure 3. Schematic of the iterative method used to retrieve fog activation properties from dry and ambient humidity particle size distributions.

centration of activated aerosol particles N_{ccn} at five supersaturations from 0.1 % to 0.5 %. Scanning takes about 20 min. Therefore, three scans can be made during the 1 h interval and thus provide meaningful estimation of the aerosol hygroscopicity. However, supersaturation values in the chamber are greater than 0.1 % and it is well known that SS_{peak} in fog is generally lower than 0.1 % (Hammer et al., 2014; Ming and Russell, 2004; Svenningsson et al., 1992; Hudson, 1980). Therefore, assumptions must be made to extrapolate κ at lower supersaturation. This is illustrated in Fig. A3, which shows statistics of κ as a function of supersaturation for the selected time period of Fig. A2. Data points correspond to calculations with CCN data at each supersaturation and the SMPS aerosol size distribution averaged over the 1 h time interval. Solid and dashed lines superimposed on the data points correspond to mean and mean ± 1 standard deviation values, respectively. Two extrema of κ are defined. κ_{inf} , which is computed as the mean value of κ at 0.1 % minus the standard deviation, corresponds to the lowest expected value. κ_{sup} is calculated as the linear extrapolation of the mean κ at 0 % plus the standard deviation and corresponds to the highest expected value. Both values are indicated in Fig. A3 by blue and red dots, respectively. These two extreme values of κ will then provide extreme values of activation properties for a given dry activation diameter.

Once droplets have been activated, their size distribution evolves during the fog life cycle. They can grow by water vapor diffusion or by coalescence with other droplets due to gravitational and turbulent motions and they can also evaporate in the case of mixing with clear air or changes in temperature. Following Noone et al. (1992), the mixing with clear air has similar consequences to an external mixture for

aerosols, which means that some particles can deactivate and their diameter can be under the critical wet diameter because of evaporation. In convective clouds, the activation occurs mainly within the first tenth of meters above the cloud base, and it is thus possible to directly sample the resulting droplet spectra with instrumented aircraft. By contrast, in the case of fog, the activation first occurs at the fog onset; the ensuing vertical development of the fog layer depends on many processes, among which the radiative cooling at the fog top plays a key role. The fog onset is generally defined by a drop of visibility below the 1 km threshold, but this is still subject to debate (see discussion in Elias et al., 2015). Time series show indeed various case-to-case CDNC evolution during the first fog hour of the fog event, with an instantaneous formation, or conversely, a slower increase towards a stable value as illustrated in Fig. A2. To estimate the most representative fog droplet size distribution of the activated distribution, we average the composite wet particle size distribution derived from both WELAS and FM-100 measurements over a time interval from the beginning of the fog event during which the CDNC reaches a stable value for a sufficiently long time. On average, this time period is selected from 30 min to 1 h after the fog beginning. It is delimited by blue segments in Fig. A2. We therefore also assume that local measurements at 2 m height are representative of the fog layer.

3.3 Method validation

Figure 4 illustrates the impact of κ variability for two fog cases of Fig. 1. The average aerosol dry size distribution measured by the SMPS recalculated at 100 % RH with the κ_{inf} and κ_{sup} values (blue and purple lines, respectively) is

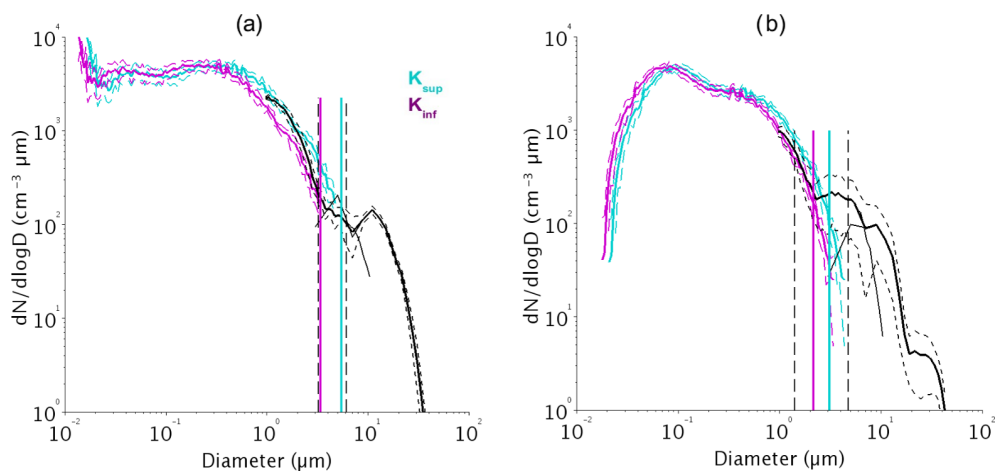


Figure 4. The 1 h wet number size distribution before the fog event resulting from the hygroscopic growth at $RH = 100\%$ of the dry distribution measured by the SMPS calculated with κ_{inf} (purple) and κ_{sup} (cyan) for cases (a) f6 and (b) f20. The in-fog composite size distribution of Fig. 1 is superimposed (black). The solid lines represent mean values and the dashed lines show ± 1 standard deviation. The vertical segments indicate the mean values of D_w corresponding to each κ (solid color lines) and the extreme values (dashed lines).

superimposed on the composite fog droplet number size distribution. Vertical segments indicate the corresponding D_w diameters. These values are 3.37 and 5.45 μm , corresponding to N_{act} of 91 and 61 cm^{-3} , respectively, for the f6 case (Fig. 4a), and 2.15 and 3.08 μm , which correspond to N_{act} of 116 and 86 cm^{-3} , respectively, for the f20 case (Fig. 4b). The dashed lines correspond to the two calculations which produce extreme values and thus delimit the range of possible values. If the spectrum remains constant as in Fig. 4a, these values are rather identical. But if the shape of the spectrum evolves during the 1 h time period, the range of possible values increases as in Fig. 4b, with extreme values of 1.4 and 4.75 μm corresponding to N_{act} of 266 and 21 cm^{-3} , respectively, leading to an uncertainty of a factor of 10 for N_{act} . There is good agreement between both distributions on the overlap area corresponding to hydrated (non-activated) particles. This indicates that the hygroscopic growth of all dry particles measured before the fog event with the derived κ values is consistent with the ambient measurements at the beginning of the fog. This also allows to neglect the mixture issue. Indeed, the dry (growth) and wet aerosol particles' shape could not fit with our assumption for a significant external mix, as there is no reason for an external mix to be homogeneously distributed on the dry aerosol number size distribution. Nevertheless, a significant external mix would have implied less aerosol particles activated per diameter with our method that would lead to an overestimation of D_d .

Our method provides a satisfactory estimation of κ , and the selected time periods are adequate to estimate the activation properties.

4 Results

4.1 Fog activation properties

Table 3 presents results for the 23 fog events with mean values and uncertainty intervals for each parameter (κ , D_d , D_w , SS_{peak} and N_{act}). The mean value is the average between the two values obtained from the above procedure with the average wet composite size distribution and the two κ values. The uncertainty interval is determined by the extreme values resulting from the average wet composite size distribution minus (plus) 1 standard deviation and the lowest (highest) κ value, respectively.

Values of κ_{inf} and κ_{sup} are relatively close. Indeed, the relative standard deviation to the average value κ_{mean} ranges between 9% and 25%, except for the f14 case, which reaches 35%. It follows that κ_{mean} which corresponds roughly to κ value at $SS \approx 0.05\%$, is representative of the κ value at the actual supersaturation. As a result, κ_{mean} values range from 0.09 to 0.3 with a median of 0.17. They are in agreement with those determined by Hammer et al. (2014), who found κ between 0.06 and 0.27 centered at 0.14 ($SS \leq 0.11\%$), and Jurányi et al. (2013), who reported κ between 0.08 and 0.24 ($SS = 0.1\text{--}1\%$) at the same site. They are slightly lower than those observed for continental aerosols; Andreae and Rosenfeld (2008), for example, suggested to use $\kappa = 0.3 \pm 0.1$. As already pointed out by Hammer et al. (2014), local emissions of road traffic and residential wood burning are likely responsible for these low values of κ . Figure 5 shows the scatterplot of the hygroscopic parameter as a function of the critical supersaturation for each of the 23 fog events with their uncertainty interval. Overall, the variability of κ remains very low, since 80% of the κ values range between 0.13 and 0.27. This suggests that the possible variations in aerosol composition

Table 3. Activation properties of the 23 fog events: hygroscopicity parameters κ_{inf} and κ_{sup} , mean values and uncertainty intervals of the dry diameter D_{d} , the wet diameter D_{w} , the supersaturation SS_{peak} and the number concentration of activated particles N_{act} . The 5th, 25th, 50th, 75th and 95th percentiles of the distribution of the 23 cases are indicated on the last lines for each parameter.

Fog ID	κ_{inf}	κ_{sup}	D_{d} μm	D_{d_i} μm	D_{d_s} μm	D_{w} μm	D_{w_i} μm	D_{w_s} μm	SS %	SS_i %	SS_s %	N_{act} cm^{-3}	N_{act_i} cm^{-3}	N_{act_s} cm^{-3}
f1	0.11	0.16	0.46	0.35	0.50	4.24	2.56	5.14	0.04	0.03	0.06	33.5	1	206
f2	0.10	0.16	0.32	0.25	0.39	2.42	1.49	3.57	0.06	0.04	0.10	131	51	266
f3	0.06	0.12	0.45	0.44	0.49	3.31	2.66	4.25	0.05	0.04	0.06	13.5	6	16
f4	0.12	0.24	0.38	0.26	0.44	3.65	1.73	5.15	0.04	0.03	0.09	104	46	341
f5	0.20	0.35	0.4	0.32	0.45	4.83	2.89	6.56	0.03	0.02	0.05	41	16	111
f6	0.20	0.41	0.37	0.34	0.41	4.41	3.19	6.08	0.04	0.02	0.05	76	51	106
f7	0.14	0.28	0.35	0.31	0.41	3.46	2.27	5	0.05	0.03	0.07	111	61	186
f8	0.17	0.29	0.36	0.24	0.50	3.79	1.73	6.9	0.04	0.02	0.09	53.5	1	226
f9	0.14	0.25	0.47	0.39	0.50	5.22	3.43	6.37	0.03	0.02	0.04	18.5	6	66
f10	0.12	0.19	0.35	0.27	0.47	3.03	1.83	5.15	0.05	0.03	0.08	146	31	371
f11	0.13	0.20	0.43	0.42	0.49	4.67	3.59	5.63	0.03	0.03	0.04	53.5	21	106
f12	0.14	0.21	0.47	0.42	0.50	4.9	3.76	5.84	0.03	0.03	0.04	28.5	11	61
f13	0.17	0.22	0.34	0.24	0.38	3.25	1.73	4.04	0.05	0.04	0.09	106	71	281
f14	0.13	0.38	0.41	0.34	0.47	4.91	2.66	7.25	0.03	0.02	0.06	31	11	81
f15	0.19	0.31	0.43	0.37	0.47	5.15	3.54	6.6	0.03	0.02	0.04	21	11	46
f16	0.10	0.16	0.43	0.39	0.47	3.69	2.87	4.65	0.04	0.03	0.05	23.5	16	36
f17	0.12	0.18	0.45	0.41	0.47	4.15	3.23	4.97	0.04	0.03	0.05	36	26	66
f18	0.12	0.18	0.30	0.20	0.39	2.37	1.13	3.86	0.07	0.04	0.13	141	41	401
f19	0.13	0.22	0.39	0.34	0.44	3.82	2.66	5.02	0.04	0.03	0.06	56	26	101
f20	0.21	0.34	0.26	0.19	0.37	2.62	1.4	4.75	0.06	0.03	0.11	101	21	266
f21	0.10	0.16	0.37	0.29	0.47	3.03	1.85	4.67	0.05	0.03	0.08	161	31	366
f22	0.14	0.21	0.21	0.17	0.29	1.44	0.98	2.65	0.11	0.06	0.15	264	116	371
f23	0.14	0.21	0.43	0.37	0.47	4.27	3.03	5.33	0.04	0.03	0.05	18.5	11	41
5th	0.12	0.18	0.35	0.25	0.41	3.03	1.73	4.65	0.03	0.02	0.05	18.5	11	41
25th	0.13	0.21	0.39	0.34	0.47	3.79	2.66	5.14	0.04	0.03	0.05	28.5	11	66
50th	0.14	0.21	0.43	0.37	0.47	4.27	3.03	5.33	0.04	0.03	0.06	53.5	21	111
75th	0.17	0.29	0.43	0.39	0.49	4.67	3.19	6.08	0.05	0.03	0.09	111	46	281
95th	0.21	0.41	0.47	0.44	0.50	5.22	3.76	7.25	0.11	0.06	0.15	264	116	401

occurring during the campaign due to variations in air mass type or source contributions have limited effect on the CCN ability of aerosol in fog conditions. Based on this study, a single κ value of 0.17 could be used for further modeling studies in this area.

The mean dry activation diameter values, 0.39 μm for the median and 0.35–0.43 μm for the 25th–75th percentiles, are rather high, which indicates that only largest aerosol particles are activated. This is in agreement with previous studies of Noone et al. (1992), Ming and Russell (2004) and Hammer et al. (2014) of the supersaturation occurring in fog. The wet activation diameter D_{w} is 3.79 μm in median and range from 3.03 and 4.67 μm for the 25th and 75th percentiles. These values are slightly higher than the median value of 2.6 μm reported by Hammer et al. (2014) but are consistent with the mean value of 4 μm by Elias et al. (2015). The latter determined $D_{\text{w}_{50\text{th}}}$ from a November 2011 data set as the intersection between the two log-normal distributions fitting the particle volume distribution measured by the WELAS. That is also highly consistent with statistics on a no-fog event at

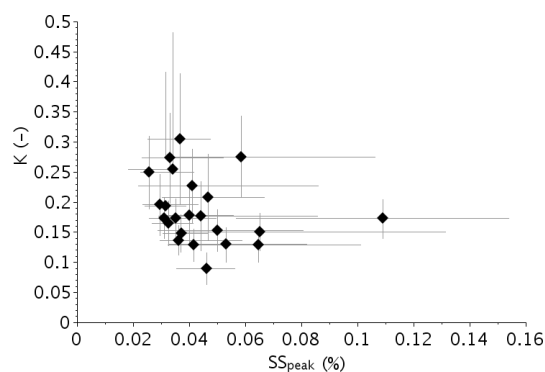


Figure 5. Hygroscopicity parameter κ as a function of the critical supersaturation for the 23 fog events. Diamonds correspond to mean values and error bars to uncertainty intervals.

SIRTA, whereby particle diameters as measured by WELAS do not exceed 2 μm . Except for the f22 event, mean values of D_{w} are larger than 2.37 μm throughout the field campaign. The differences between hydrated and activated particles can

be seen in Fig. 4a and b, and correspond fairly well to the average wet activation diameters derived for such cases that are equal to 4.41 and 2.62 μm , respectively. Moreover, it appears that the extreme values indicated by the dashed vertical segments largely maximize the uncertainty of the retrieval.

Critical supersaturation occurring in fog is very low with a median value of 0.043 % and 25th–75th percentiles of 0.035 %–0.051 %. This result was expected for fog conditions, as already shown in experimental and modeling studies (Hudson, 1980; Svenningsson et al., 1992; Ming and Russell, 2004; Hammer et al., 2014; Shen et al., 2018; Boutle et al., 2018).

Finally, the corresponding concentration of activated particles (N_{act}) is 53.5 cm^{-3} for the median and 28.5–111 cm^{-3} for the 25th–75th percentiles. These values are slightly lower than the N_{FM} ones derived from FM-100 data (reported in Fig. 2) which give a median and 25th–75th percentiles of 61 and 34–103 cm^{-3} , respectively. Moreover, the median of the average concentration values of particles with diameter in the range of 0.96–50 μm at the beginning of the fog events reaches 389 cm^{-3} with 25th–75th percentiles of 260–660 cm^{-3} . It follows that only a small fraction of the fog hydrated particles can be activated to droplets due to low critical supersaturations encountered in fog. Obviously, because of the uncertainty interval, N_{act} is very broad. For instance, this uncertainty is twice the N_{act} value itself on average but could be as high as a factor of 6 for the f1 event. However, as mentioned before, N_{act_i} and N_{act_f} are the extreme possible values derived by cumulating uncertainties on κ and on the variability of measurements during the 1 h time period.

4.2 Impact of aerosol particles on fog droplet concentration

Mean values of N_{act} are shown in Fig. 6 as a function of the other activation parameters (SS_{peak} , D_{d} and κ) for the 23 fog events. N_{act} values increase with SS_{peak} (Fig. 6a) and consequently decreased with D_{d} (Fig. 6b). In contrast to Fig. 6b, Fig. 6c reveals no trend between N_{act} and κ , suggesting that size matters more than chemical composition for a particle to be activated. Indeed, the largest aerosol particles are activated at low SS independently of their κ values. There is some scatter in Fig. 6b for values of $D_{\text{d}} \leq 0.4 \mu\text{m}$, which suggests that some variability of the dry aerosol size number distribution between the different cases occurs mainly below this threshold.

Our results extend to the lower supersaturation conclusion of Fitzgerald (1973) and Andreae and Rosenfeld (2008) on clouds that particle size plays a much greater role than composition in regulating cloud droplet nucleation. The fundamental reason is that soluble mass changes with the third power of particle diameter but only linearly with soluble fraction (Andreae and Rosenfeld, 2008). The color code in Fig. 7a shows that, at a given supersaturation, the highest values of the CCN ratio are associated with the highest values of

the hygroscopic parameter (i.e., chemical composition), and that at a given CCN ratio, decrease as the supersaturation increases. However, more samples are needed before a robust interpretation of that feature is to be established.

To remove the influence of the aerosol number concentration, we normalize N_{act} by the number concentration of activable aerosols (N^*), defined as the number concentration of aerosol particles with diameter $\geq 200 \text{ nm}$, which corresponds to the smallest dry diameter of Table 3. Figure 7b shows a good correlation between N_{act} normalized by N^* and the inferred critical diameter. We obtained similar correlation when N^* was calculated for the range of critical diameter inferred during the campaign (150 to 400 nm). This suggests that N^* using a single critical diameter can be considered as a good proxy of the number of activable particles. This can be understood since the supersaturation occurring in fog has a narrow range of value. CCN at a supersaturation between 0.02 % and 0.11 % are mostly composed of accumulation-mode particles.

Figure 8 reveals that almost no relationship exists between N_{act} and N^* . This surprising behavior demonstrates that the concentration of fog droplets was roughly independent of the aerosol number concentration. This is most likely due to the low SS occurring in fog leading to very selective criteria on particle diameter for droplet activation. This is also consistent with the very large derived dry diameter. Nevertheless, this figure also shows that, for high concentration of activable particles (N^*), lower N_{act} values are obtained. It follows that the number concentration of fog droplets is mainly controlled by the peak supersaturation with low influence from the aerosol loading. This is illustrated in Fig. 9, where N_{act} values are plotted as a function of SS_{peak} in log scale (black diamonds, same data as Fig. 6a) superimposed on the statistics of CCN measurements (grey diamonds). Compared to the compilation of CCN spectra reported in Andreae and Rosenfeld (2008) (their Fig. 2), our CCN data are spread between SCMS (Small Cumulus Microphysics Study) data (Hudson and Yum, 2001) and continental cases. However, we observe a strong decrease of N_{act} for $\text{SS}_{\text{peak}} \leq 0.1 \%$, similar to the ASTEX (Atlantic Stratocumulus Transition Experiment) data collected in maritime stratocumuli and to Hudson (1980) fog activation spectra measurements. This indicates that only a very small fraction of aerosols is activated for such low values of supersaturation.

These results have been fitted using three different equations commonly used in modeling. The typical CCN parameterization is proposed by Twomey (1959): $N_{\text{CCN}} = CS^k$, where C represents the CCN number concentration at $\text{SS} = 1 \%$ and the parameter k that varies significantly (Martins et al., 2009). The next formula is suggested by Ji et al. (1998): $N_{\text{CCN}} = N(1 - \exp(-BS^k))$, where N is the total number concentration of CCN and B and k are empirical coefficients to be determined. Finally, the more general description of the activation spectra is proposed by Cohard et al. (1998): $N_{\text{CCN}} = CS_{v,w}^k F(\mu, \frac{k}{2}, \frac{k}{2} + 1; -\beta S_{v,w}^2)$, where C is propor-

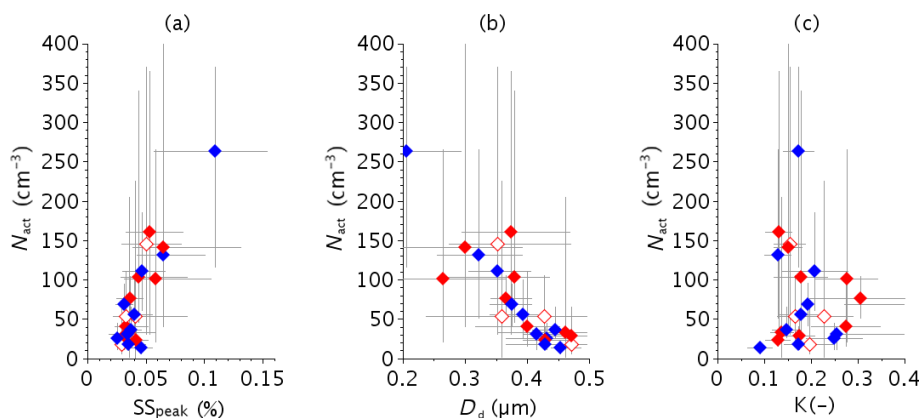


Figure 6. Number concentration of activated particles N_{act} as a function of (a) the critical supersaturation, (b) the dry critical diameter and (c) κ for the 23 fog cases. Diamonds correspond to mean values and error bars to uncertainty intervals. The same symbols as in Fig. 2 are used.

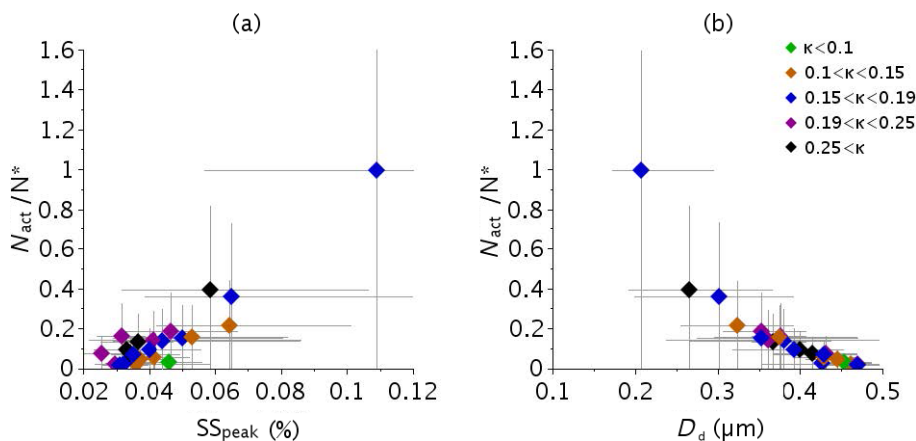


Figure 7. Ratio of activated particles to concentration of aerosol particles with diameter > 200 nm (activated fraction) for the 23 fog cases as a function of (a) the critical supersaturation and (b) the dry critical diameter. Diamonds correspond to mean values and error bars to uncertainty intervals. Color as a function of κ is indicated in the legend.

tional to N_{CCN} that would be activated when supersaturation tends to infinity and parameters k , μ and β are adjustable parameters depending on the aerosol properties. Figure 9 shows that the parameterization of Ji et al. (1998) (green line), with parameters as indicated by the legend, better reproduces the decrease of N_{act} for $\text{SS} < 0.1\%$ compared to the Twomey expression (blue line) but the drop is not as sharp as in the data. In contrast, the parameterization of Cohard et al. (1998) (red line) provides the best fit of the data for lower values of SS.

Figure 10 shows the CCN concentration that would be activated at $\text{SS} = 0.1\%$ as a function of the aerosol activable concentration for the 23 fog events. The expected strong correlation between activated particles and aerosol concentration then clearly appears. This confirms that the apparent independence observed in Fig. 8 is simply due to the low supersaturation values experienced in fogs. The symbol color depends on κ as the preceding figures. Data are aligned according to their hygroscopic parameter. For a given aerosol par-

ticle's activable concentration, the number of activated particles at $\text{SS} = 0.1\%$ increases with κ , which is consistent with Fig. 7a.

One can note in Fig. 8 that the highest values of N_{act} decrease as N^* increases. Indeed, while aerosol particles' activable concentration increases (N^*), variability of SS_{peak} decreases and tends toward low values as reported in Fig. 11; thus, less particles can be activated. This suggests that the supersaturation reached in fog could be limited by the activable aerosol concentration; if numerous aerosol particles are available, they are efficient enough to uptake the water vapor excess and therefore limit the supersaturation. Hudson (1980) found that aerosol particle concentration values do not have a great effect on fog supersaturation. Our result however supports the conclusions drawn by the numerical study performed by Bott (1991) that the higher the particle concentration is, the lower the supersaturation is. Note that no trend appears with κ in Fig. 8, which means that the sensitivity of

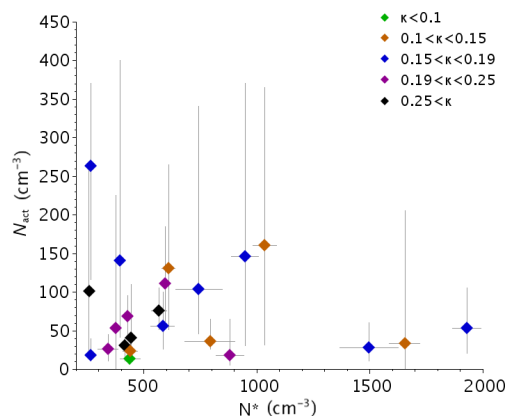


Figure 8. Number concentration of activated particles (N_{act}) as a function of N^* , the concentration of aerosol particles with diameter > 200 nm for the 23 fog events. For N^* , diamonds correspond to median values and error bars to 25th–75th percentile intervals. Color as a function of κ is indicated in the legend.

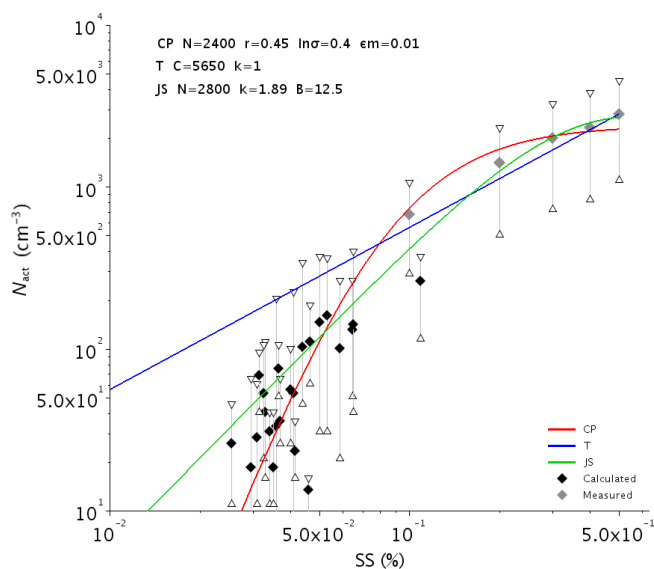


Figure 9. Number concentration of activated particles N_{act} as a function of the critical supersaturation for the 23 fog events (black) superimposed for the statistics of CCN chamber measurements (grey). For CCN measurements, diamonds correspond to median values and error bars to 25th–75th percentile intervals. Color lines correspond to fitted results of $N_{\text{CCN}} = CS_{v,w}^k F(\mu, \frac{k}{2}, \frac{k}{2} + 1; -\beta S_{v,w}^2)$ (red), $N_{\text{CCN}} = CS^k$ (blue) and $N_{\text{CCN}} = N(1 - \exp(-BS^k))$ (green), with parameters as indicated in the legend.

this process to the hygroscopicity is very weak. As suggested by Bott et al. (1990), this could be due to a lowering of vapor supersaturation with the hydration of the numerous aerosol particles.

Figure 6 compared to Fig. 2 shows that there are no obvious differences between STL and RAD cases. As a result, the

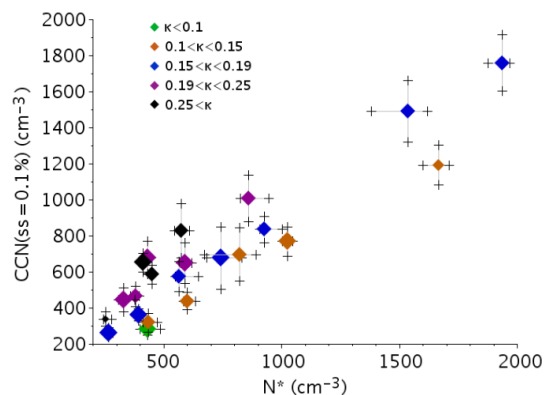


Figure 10. Median values for the 23 fog events of the number concentration measured by the CCN chamber at $SS = 0.1\%$ as a function of the concentration of aerosol particles with diameter > 200 nm. Color as a function of κ is indicated as in Fig. 8. Symbol size is proportional to the mean diameter.

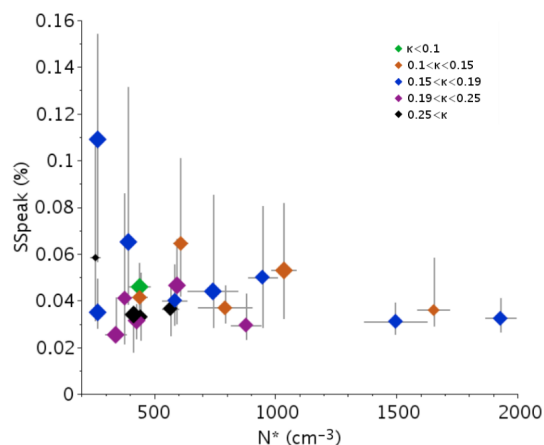


Figure 11. Critical supersaturation as a function of the concentration of aerosol particles with diameter > 200 nm for the 23 fog cases. Diamonds correspond to mean values and error bars to uncertainty intervals. Color and size of symbols are as in Fig. 10.

median values of N_{act} are 40 and 58 cm^{-3} , and the 25th–75th percentiles are 25 – 115 and 33 – 109 cm^{-3} , for STL and RAD cases, respectively. Radiative cases have thus higher values of activated particles but the percentile intervals are rather similar and the difference is less pronounced than the factor of 2 obtained from N_{FM} derived from the FM-100 measurements only.

4.3 Impact of CCN concentration on fog microstructure

The mean fog droplet concentration ($N_{\text{d}_{1h}}$) averaged over the 1 h time interval at the beginning of the fog event is plotted in Fig. 12a as a function of N_{act} . Obviously, they are almost identical, apart from small deviations that come from averaging the N_{act} values derived for both κ values. The color

code in Fig. 12a shows no differences between radiative (red diamonds) and stratus-lowering (blue diamonds) fogs during the formation phase.

Once droplets have been activated, they compete for the available water vapor; when their number increases, the size they can reach by water vapor diffusion growth decreases. This inverse relationship between the number and the size of droplets is clearly depicted in Fig. 12b, which displays the concentration $N_{d_{1h}}$ vs. the corresponding mean diameter D_m . Indeed, the highest $N_{d_{1h}}$ of 255 cm^{-3} corresponds to the lowest D_m of $4\text{ }\mu\text{m}$, and for fog events with $N_{d_{1h}} > 120\text{ cm}^{-3}$ the mean diameter cannot exceed $7\text{ }\mu\text{m}$, while it reaches twice this value for concentration as low as 25 cm^{-3} . However, some scatter appears for samples with $N_{d_{1h}} < 50\text{ cm}^{-3}$. The symbol color corresponds to the mean LWC values, as indicated by the legend. Lowest values of D_m are associated with lowest values of LWC, which suggests that a lack of available liquid water could also limit the droplet growth, even when there are few of them.

During a fog life cycle, several processes contribute to the droplet size distribution. Droplets can grow by water vapor diffusion or by collision–coalescence with other droplets due to gravitational and turbulent motions. Conversely, they can evaporate if the supersaturation decreases due to heating of the air mass, for example, or in the case of mixing with clear air. Figure 13a and b show the same plots except that N_d , D_m and LWC correspond to mean values over the complete fog life cycle. The mean droplet concentration values are significantly lower, especially for cases with high N_{act} values. For instance, for cases with $N_{act} > 50\text{ cm}^{-3}$, the ratio $N_{d_{cycle}}/N_{act}$ ranges from 0.25 to 0.92, with an average value of 0.58. Values reported in Brenguier et al. (2011) exhibit a clear separation between cumulus and stratocumulus clouds, with values from 0.32 to 0.56 (0.46 on average) and from 0.72 to 0.96 (0.87 on average), respectively, that were attributed to differences in entrainment–mixing processes in both cloud types. Here, it appears that values are spanned over a large range with an intermediate average value. It is noteworthy that Fig. 13b reveals that $N_{d_{cycle}}$ for a radiative fog cannot exceed 70 cm^{-3} , which suggests that this reduction is more pronounced in the case of a radiative fog, while STL cases seem less affected. Indeed, for cases with $N_{act} > 100\text{ cm}^{-3}$, the average $N_{d_{cycle}}/N_{act}$ is 0.80 for a STL fog, a value close to stratocumulus clouds, while it drops down to 0.45 for RAD cases. This suggests that the processes involved in the reduction of the CDNC in a radiative fog are as efficient as the entrainment–mixing occurring in cumulus clouds. Note that many studies (Pilié et al., 1975; Choulaton et al., 1981; Gerber, 1991; Bergot, 2013; Maronga and Bosveld, 2017), among others, pointed out the key mechanism of the turbulent mixing in fogs. Recent work of Bergot (2016) has shown, with a LES study, evidence of the key role of dry downdraught at the top of the fog layer on dissipation. These downdraughts would even then reach the surface, allowing

the Sun to warm it. Wærsted et al. (2017), using experimental data, confirmed the importance of top processes for the fog evolution. The dilution ratio being as high as in cumulus clouds in radiative fog brings a new element on that point, meaning that turbulent mixing must be of importance.

Figure 13b shows that mean diameter values averaged over the fog life cycle are similar to the previous ones. The resulting cluster of data points does not reveal any general trend and the anticorrelation between size and number of droplets is no longer noticeable. Indeed, the evolution of microphysical properties during the fog life cycle is complex and highly varies from case to case and depends on many parameters. As already noticed by Burnet et al. (2012) in their Fig. 3, various values of D_{eff} and N_{FM} can be observed for a given LWC_{FM} . A comprehensive analysis on this topic has been performed and will appear in a forthcoming paper.

5 Discussion

In this experimental study, we derived accurate estimations of the activated droplet concentration at the beginning of a fog event. This was done through a careful estimation of the critical supersaturation and of the wet critical diameter, which allowed us to integrate the composite size distribution derived from WELAS and FM-100 measurements. We have shown that the derived parameters are consistent with the parameterization of Cohard et al. (1998). They are also consistent with measurements reported in Hudson (1980), who developed the isothermal haze chamber to extend the CCN measurements to supersaturation range below 0.1 % and derived fog condensation nuclei (FCN). Indeed, the author found effective supersaturation between 0.06 % and 0.1 % as well as sharp changes in FCN spectra below 0.1 % with activated concentration at $SS = 0.4\%$ ranging from 0.8 cm^{-3} over the sea to 250 cm^{-3} for polluted cases.

To our knowledge, no other study has experimentally retrieved concentration of activated particles (and thus droplets) in fog. Usually, measurements of droplet spectra reported in the literature were performed by single particle counters such as FSSP, CDP or FM-100 over the range 2–50 μm . Figure 1 shows typical droplet size distributions observed during the campaign with measures from 2 up to 50 μm .

The corresponding droplet concentrations are in the range of a tenth to several hundred cm^{-3} (Wendisch et al., 1998; García-García et al., 2002; Gultepe and Milbrandt, 2007a; Niu et al., 2012; Price, 2011; Liu et al., 2011; Burnet et al., 2012; Lu et al., 2013; Zhao et al., 2013; Boutle et al., 2018). In our study, droplet concentrations are rather low for semi-urban conditions, with a maximum value of 150 cm^{-3} . For instance, Liu et al. (2017) reported values in China as high as 1000 cm^{-3} . Moreover, our values of critical diameter are large, with a median value of $3.8\text{ }\mu\text{m}$. Consequently, the use of FM-100 data over the range 2–50 μm will lead to an over-

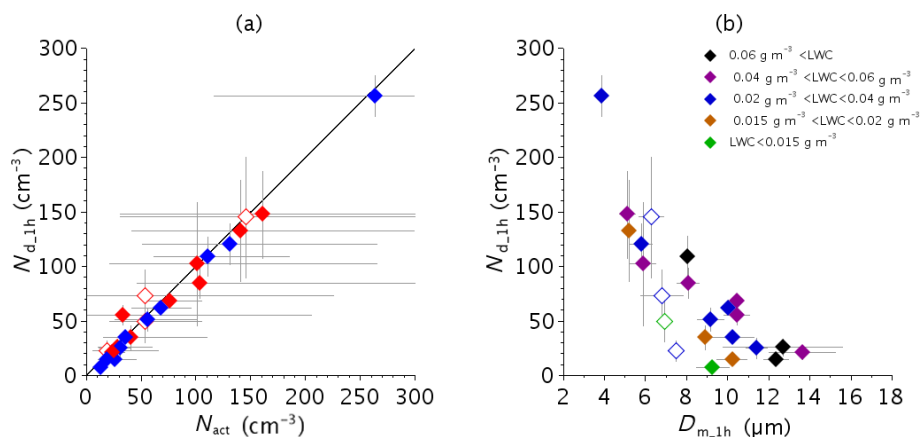


Figure 12. Droplet number concentration at the beginning of the fog for the 23 cases as a function of (a) the activated particle concentration (N_{act}) and (b) the mean fog diameter. Diamonds correspond to median values and error bars to 25th–75th percentile intervals. Symbol colors are as a function of (a) fog type as in Fig. 2 and (b) LWC values as indicated in the legend.

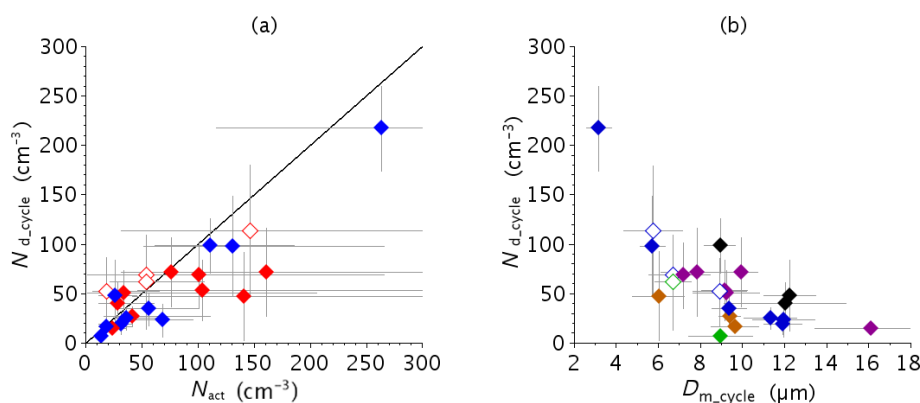


Figure 13. Same as Fig. 12 except that statistics of N_d , D_m and LWC are computed over the complete fog life cycle.

estimation of the droplet concentration, because it includes deliquescent aerosol particles that are not activated. A noticeable exception comes from the f22 case with a maximum value of 264 cm^{-3} . Surprisingly, the median N_{FM} value for this case is only 36 cm^{-3} , which is rather low. This is explained by the fact that this case also corresponds to the lowest critical diameter ($D_w = 1.44 \mu\text{m}$); therefore, the contribution from WELAS data is the most important and emphasizes the underestimation of the FM-100 in the first classes (see Fig. A1). Thus, for cases with low D_w , WELAS data must be taken into account to avoid an underestimation of the concentration. In a modeling purpose, distinguishing hydrated aerosol particles from droplets allows to calculate an accurate repartition of the vapor deposition on the droplet size distribution. Indeed, behavior of hydrated particles and droplets is quite different: the second one grows as long as there is vapor in excess, while the first one stays at an equilibrium diameter. The second one may grow enough to produce drizzle and have strong interaction with radiation. A clear distinction should improve the representation of processes in numerical

weather prediction models and thus on the visibilities' forecast.

A surprising result is the observed decrease of droplet concentrations with increasing aerosol loading. This is in contrast with recent numerical simulations that exhibit a strong positive correlation between aerosol and droplet concentrations in fog (Zhang et al., 2014; Stolaki et al., 2015; Maalick et al., 2016). This could be due to a limitation of the supersaturation according to Bott (1991). The sensitivity study of Stolaki et al. (2015) indicates that the fog droplet concentration during the mature stage is 2.6 times greater (2.9 times lower) when doubling (halving) the CCN number concentration of the accumulation mode. Indeed, at the fog onset, the droplet concentration at the surface reaches 467 cm^{-3} in their simulation with a double CCN concentration. This value is much larger than that of our estimations, which is $\leq 150 \text{ cm}^{-3}$, except in one case where it reaches 264 cm^{-3} . For their reference run, also this value is as high as 304 cm^{-3} . The CCN activation scheme used by the model follows the expression of Cohard et al. (1998). Given the CCN activa-

tion spectrum derived from their values of CCN size distribution parameters (geometric mean radius and standard deviation of 0.1525 μm and 2.33, respectively) and concentration (540 cm^{-3}) and solubility of aerosols (0.4), this corresponds to an estimation of the maximum supersaturation of 0.15 %. Even for activated concentration of 250 cm^{-3} , the corresponding SS is still $\simeq 0.1\%$. We have shown that CCN measurement at $\text{SS} = 0.1\%$ exhibits a linear increase of the droplet concentration with the activable aerosol concentration. Consequently, with such high supersaturation values, it is consistent that their simulations produce a strong positive correlation between aerosol and droplet concentration.

More generally, one shall keep in mind that the use of the CCN activation spectra provides a satisfactory estimation of the activated concentration, providing that the maximum supersaturation is correctly diagnosed. For instance, a current limitation of schemes with adjustment to saturation with parameterized peak supersaturation is that the formulae do not take into account pre-existing liquid water in the model grid box (Thouren et al., 2012). This leads to a significant overestimation of the supersaturation peak value that in turn will overestimate the activated concentration. Another limitation is the use of minimum values for vertical velocities in typical supersaturation parameterization as pointed out by Boutle et al. (2018). They are not necessarily consistent with fog dynamics and can lead to an overestimation of the supersaturation peak value. The same overestimation could occur while using the saturation adjustment assumption instead of a diffusion growth/evaporation equation according to Khain et al. (2015). Moreover, in a sensitivity study in which the CCN concentration is increased (or decreased) while keeping other parameters constant, it intrinsically introduces a strong dependency between droplet and aerosol particles. Indeed, the pre-existing hygroscopic aerosol particles contribute to the sink term of water vapor during the radiative cooling that will also reduce the supersaturation peak value at the fog onset and then limits the number of activated particles. Thus, changing the aerosol properties in models using such a scheme provides a useful way to modify the droplet concentration and study the impact of this latter on the fog life cycle. However, it should not be used to assess the impact of aerosol properties themselves on the fog life cycle.

6 Conclusions

In situ microphysical measurements collected during 23 wintertime fog events sampled from October 2010 to March 2013 at the semi-urban SIRTa site near Paris have been examined to document their properties. They reveal a large variability of their characteristic values in terms of number concentration and size of fog droplets depending on the different cases, as well as various evolution of these properties during the fog life cycle. The aerosol background exhibits also a highly variable number concentration of parti-

cles $> 10\text{ nm}$ before the fog onset. The objective of this paper was to evaluate the impact of aerosol particles on the fog microphysics. As a first step, we focused on the relationship between aerosol and fog droplet number concentrations, as we expected that they follow the same trend with a greater number of fog droplets with an increase of the aerosol loading.

To derive accurate estimations of the actual activated fog droplet number concentration (N_{act}), we determined the hygroscopicity parameter, the dry and wet critical diameters, and the critical supersaturation for the 23 events by using an iterative procedure based on the κ -Köhler theory which combined CCN measurements, dry particle distribution from an SMPS and composite wet particle size distribution at ambient humidity derived from WELAS and FM-100 measurements. These data are averaged over a 1 h time period before and during the fog onset to characterize the air mass and the fog properties, respectively.

Values of the hygroscopicity parameter κ extrapolated at $\text{SS} \simeq 0.05\%$ were found to range from 0.09 to 0.3 which are characteristics for urban aerosol. They are rather similar from case to case; thus, we recommended to use $\kappa = 0.17 \pm 0.05$ for fog modeling studies in this urban area. Our study reveals low values of the derived critical supersaturations with a median of 0.043 % and 25th–75th percentiles of 0.035 %–0.051 %. Consequently, wet and dry activation diameters are high and the corresponding N_{act} values are low for continental conditions, with a median concentration of 53.5 cm^{-3} and 25th–75th percentiles of 28.5–111 cm^{-3} . Our results depict a sharp drop of N_{act} as critical supersaturations decrease which is best fitted by the parameterization of Cohard et al. (1998).

Concentration of activated particles using a single critical diameter can be considered as a good proxy of the number of activable particles. No detectable trend between concentration of aerosol particles with diameter $> 200\text{ nm}$ and N_{act} was observed. But N_{act} could be modulated for high concentration of aerosol particles. Hydration of numerous aerosol particles could indeed lower the supersaturation. Our analysis corroborates the modeling study by Bott (1991) suggesting that an increase of the aerosol concentration limits the SS values. In contrast, the CCN data at 0.1 % supersaturation exhibit a strong correlation with these aerosol concentrations. We therefore conclude that the droplet number concentration is roughly independent of the aerosol one because the actual supersaturations reached in these fog events are too low. Given the very diverse origin of atmospheric particles, the complexity of their composition has long been seen as a major obstacle to modeling aerosol effects on fog properties. Understanding that particles' ability to act as CCN at the supersaturation reached in fog is largely controlled by aerosol size rather than composition greatly facilitates the treatment of aerosol impacts on fog microphysics in NWP models. Nevertheless, some reservation must be issued on that point. Aerosol size distribution impacts fog microphysics more than chemical composition on the measured range on the

hygroscopic parameter during the three winters of observations. Higher or lower values of this parameter could bias this conclusion. Other measurements of aerosol and fog microphysics in other environmental conditions would be needed to pursue this issue.

Despite a stratus-lowering fog appearing to be associated with lower aerosol loading than a radiation fog event, no significant differences were observed with respect to the droplet number concentration at the fog formation when calculations were performed by taking into account the wet critical diameter. In contrast, statistics over the complete life cycle indicate that a radiation fog is subject to a pronounced decrease in the droplet concentration, while stratus-lowering cases seem unaffected. In a radiation fog, this decrease, which occurs mainly for events with $N_{\text{act}} > 50 \text{ cm}^{-3}$, is significant with a ratio of average droplet concentration to N_{act} ranging from 0.25 to 0.67. This reduction is similar to the dilution which results from entrainment–mixing in cumulus clouds, meaning that mixing with clear air may be of importance for fog microphysical evolution. The expected inverse relationship between the number and the size of the droplets at the formation phase is clearly depicted, except for some cases with low liquid water content values. However, this trend is less pronounced on average over the fog life cycle.

Nenes and Seinfeld (2003) have shown that kinetic limitations in droplet activation may retard the growth of CCN sufficiently to limit the number of activated droplets formed. The extent to which potential kinetic limitations to droplet formation exist in fog conditions will be explored in a subsequent publication focusing on the temporal evolution of droplet spectra and visibility during fog life cycle.

Data availability. Advanced data are available in Tables 2 and 3. Data are available upon request to the authors.

Appendix A

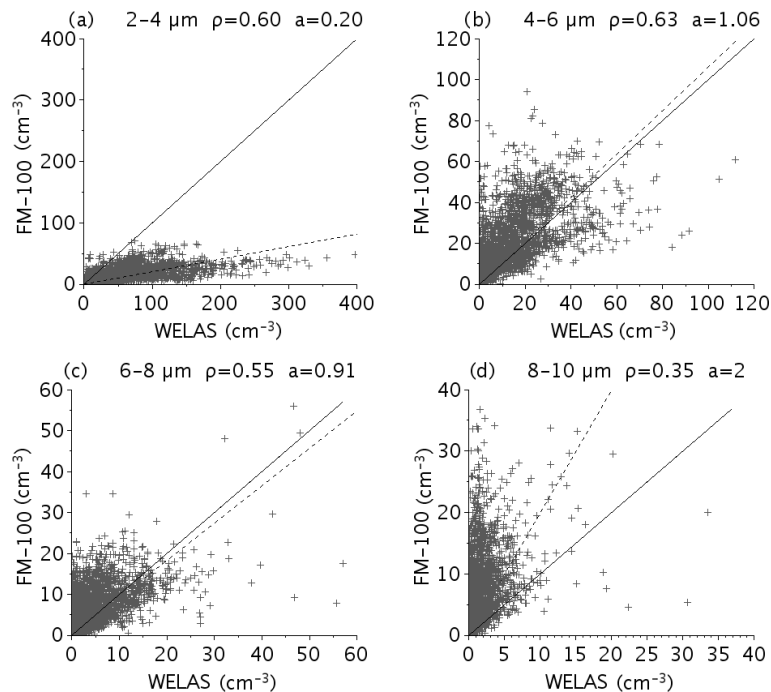


Figure A1. Scatterplot of the 5 min average particle number concentration values as measured by the FM-100 for the four first bins vs. the integrated WELAS measurement over the corresponding diameter range (a) 2–4 μm , (b) 4–6 μm , (c) 6–8 μm and (d) 8–10 μm . The solid line corresponds to 1 : 1 line and the dashed line corresponds to best-fit line with correlation coefficient ρ and slope a , as indicated in the legend.

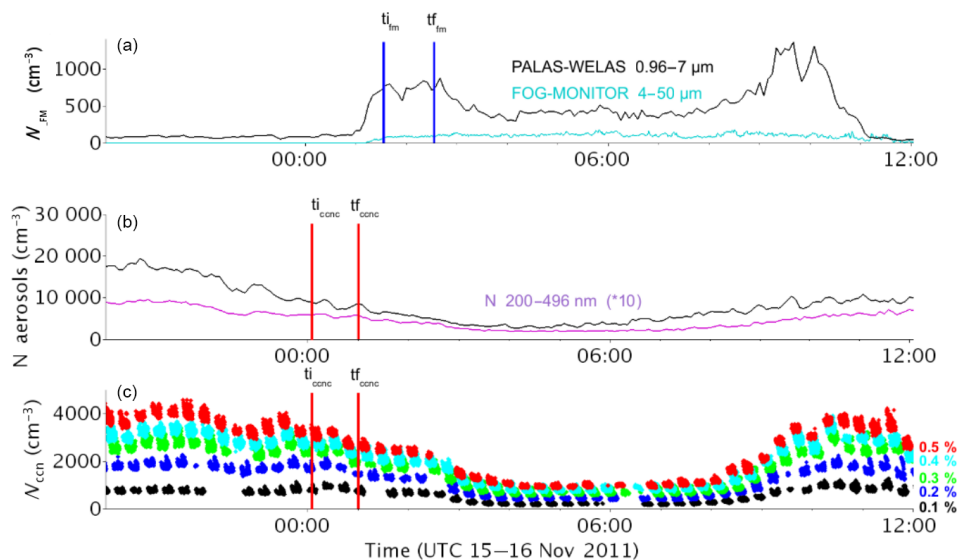


Figure A2. Time series of measurements for the f6 case: (a) particle number concentration at ambient humidity from FM-100 (cyan) and WELAS (black); (b) dry particle number concentration from SMPS: total (black) and particle with diameter > 200 nm (purple); and (c) N_{ccn} from the CCN chamber at each supersaturation with color as indicated on the label. Vertical segments indicate the selected time periods before (red) and during (blue) the beginning of the fog event.

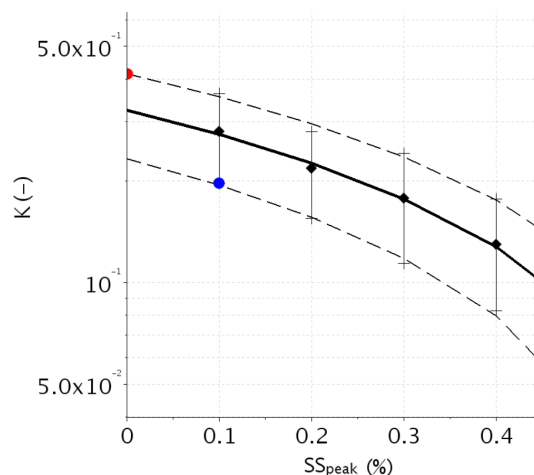


Figure A3. Derived κ values as a function of CCN chamber supersaturations from measurements over the 1 h time period before the fog event for the f6 case. Mean (diamond, solid line) and mean ± 1 standard deviation (dashed lines). Extrapolated κ_{inf} (blue dot) and κ_{sup} (red dot) are also indicated.

Appendix B: Abbreviations

CCN	Cloud condensation nuclei
N_d	Concentration of particles above the wet critical diameter (derived from WELAS+FM-100)
CN	Condensation nuclei
N_{FM}	Concentration of particles measured by the FM-100 between 2 and 50 μm diameter
N_a	Concentration of dry particles measured by the SMPS between 10 and 495 nm diameter
N_{act}	Concentration of activated particles
N^*	Concentration of dry particles between 200 and 496 nm diameter (derived from SMPS)
N_{ccn}	Concentration of activated aerosol measured by the CCN at different supersaturation
D_d	Dry critical diameter
D_w	Wet critical diameter
D_m	Median diameter
SS_{peak}	Critical supersaturation peak
SS	Supersaturation

Author contributions. This work is a part of the PhD of MM. MM and FB analyzed the data and wrote the paper. FB, MH, JCD and TE organized the field campaign. TE coordinated the field campaign. GCR helped to process the CCNC data. MM and CD performed the review process.

Competing interests. The authors declare that they have no conflict of interest.

Acknowledgements. The authors are very grateful to all SIRTAs operators and database managers. This campaign was held in the framework of the PreViBOSS project, supported by DGA/DGIS. This research was partially funded by the European Commission's Seventh Framework Programme (FP7/2007-2013) under the SESAR WP 11.2.2 project, under grant agreement 11-120809-C.

Review statement. This paper was edited by Barbara Ervens and reviewed by four anonymous referees.

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