

A PHONON CRANKING MODEL[†]

Torunn Fogel, P. Haapakoski and P.O. Lipas
University of Helsinki, Finland

According to the cranking description [1], $H_\omega = H - \omega J_x$ is the Hamiltonian in the nuclear rest frame. Its eigenstates ϕ_ω give the total laboratory energy as

$$E = \frac{(\phi_\omega, H\phi_\omega)}{(\phi_\omega, \phi_\omega)}$$

with ω determined from $(\phi_\omega, J_x \phi_\omega) = \hbar \sqrt{J(J+1)}$.

We use this formalism to calculate the ground band rotation energies from the phonon Hamiltonian

$$\hbar\omega_2 \left[\sum_m b_m^\dagger b_m + \frac{5}{2} - d(b_0^\dagger + b_0) + d^2 \right],$$

which describes harmonic quadrupole vibrations in an axially deformed body ($d \propto \beta_0$). We have used it previously as the intrinsic Hamiltonian in projection calculations [2].

To solve the H_ω problem we express J_x in terms of the phonon operators and seek a ground state solution of the form $\exp \sum_m W_m b_m^\dagger |0\rangle$ as an extension of the $\omega=0$ case [2]. The solution is exact, and the rotational energies are given by

$$E_{\text{rot}} = \frac{1}{2}\omega^2 \frac{6d^2\hbar\omega(4\omega^2 + \omega_2^2)}{(4\omega^2 - \omega_2^2)^2}, \quad \frac{6d^2\omega_2^3\omega}{(4\omega^2 - \omega_2^2)^2} = \sqrt{J(J+1)}$$

in terms of two parameters, ω_2 and d .

The moment of inertia in the limit $\omega \rightarrow 0$, $6d^2\hbar/\omega_2$, agrees with the result given by the Inglis formula. Further, taken to order ω^4 , our E_{rot} agrees with Harris' (and Mariscotti's) two-parameter rotational description. Fits to data are similar. Our upper and lower limits on E_4/E_2 are $10/3 = 3.33$ and $\sqrt{10/3} = 1.83$, in agreement with the Goldhabers [3].

In contrast to most rotational phenomenology, ours also provides wave functions upon projection from $\exp \sum_m W_m b_m^\dagger |0\rangle$. The resultant $B(E2)$ values and quadrupole moments resemble those from the projection model [2].

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[2] P. Haapakoski, T. Honkaranta and P.O. Lipas, Phys. Lett. 31B (1970) 493.

[3] G. Scharff-Goldhaber and A.S. Goldhaber, Phys. Rev. Lett. 24 (1970) 1349.