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STUDY OF <sup>132</sup>La OF A FLUENT BEAMLINE OF AN ION THERAPEUTIC MACHINE

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STATE OF  $^{112}\text{Cd}$  BY ELECTRON SCATTERING OF  
50-MeV POLARIZED PROTONS

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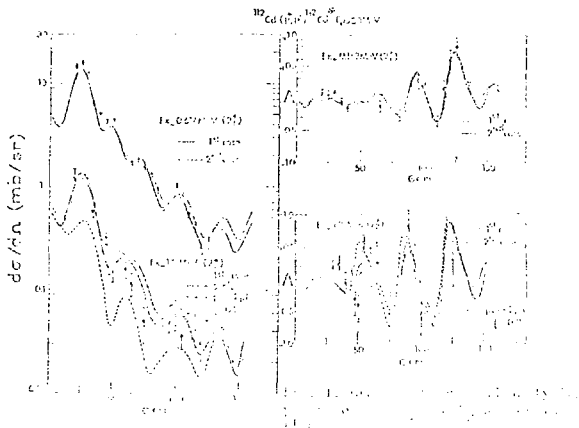
M. H. Susskind, *ibid.* (1974)

to first appear in the *Journal of Nuclear Energy, Part B*, 12, 1974, pp. 1-10.

It is shown that the  $1^+_{ph}$  and  $2^+_{ph}$  states would disappear completely if the proton inelastically scattered from a free nucleon rather than from a nucleon bound in a nucleus. The  $1^+_{ph}$  state fits the data generally but fails to fit the data at the inelasticity state  $Q^2 = 0.25 \text{ GeV}^2$  and at a scattering angle of  $30^\circ$ . The  $2^+_{ph}$  state fits this region and is found to be in a  $1^+_{ph}$  state. However, the  $1^+_{ph}$  has been reported  $^{20}$  that a proton inelastically scattered should be  $1^+_{ph}$  at the inelasticity state of the first  $1^+_{ph}$  ( $Q^2 = 0.042 \text{ GeV}^2$ ) to  $2^+_{ph}$  to describe satisfactorily this state. All the previous results, as have been discussed in detail, are of a proton inelastically scattered from a nucleon bound in a nucleus. In  $^{112}\text{Cd}$ , the data taken at the transfer scattering region are fitted with the present  $1^+_{ph}$   $^{20}$  by replacing the  $1^+_{ph}$  ( $Q^2 = 0.042$ ) and the  $2^+_{ph}$  ( $Q^2 = 0.042$ ) and deformation parameters of table 1. The calculations for the data are presented on Fig. 1.

We determined for each state the mixture of one-phonon and two-phonon components in the vibrational model. In the first case, the  $1^+_{ph}$  state is considered as a pure one-phonon state and the  $2^+_{ph}$  state by

$$|2^+_{ph}\rangle = 0.5|1^+_{ph}\rangle + 0.5|2^+_{ph}\rangle$$



In the second case, the  $2^1_1$  and  $2^3_1$  orthogonals states are represented by

$$\begin{aligned} |2^1_1\rangle &= 0.912 |1\rangle + 0.419 |2\rangle + 0.167 |3\rangle \\ |2^3_1\rangle &= 0.390 |1\rangle - 0.711 |2\rangle + 0.581 |3\rangle \end{aligned}$$

The coupling parameter  $\lambda$  is 0.10 eV in the first case and 0.12 eV in the second case (see Table 1).

For the first case (solid line) we find again that the dynamic parameters as well as the relative composition are similar to those obtained in ref. 11. Several lines are better observed with the data in case 1 than with the data in case 2 but this effect is more visible on the  $2^1_1$  than on the  $2^3_1$  (Fig. 2). It is possible that the case 1  $2^1_1$  obtained for the data in case 2 state remaining  $2^1_1$  coupling with  $1^1_1$  is 0.16.

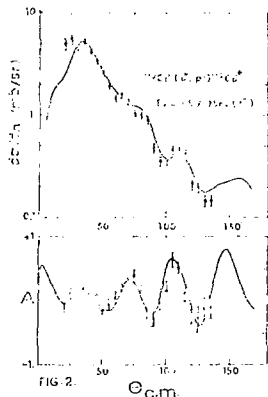


FIG. 2. Grand reaction and analyzing power of the  $Y$  state in  $^{112}\text{Cd}$ .

Fig. 2. Grand reaction and analyzing power of the  $Y$  state in  $^{112}\text{Cd}$ .

Table 1. Optical model and deformation parameters used in the coupled-channel calculations.

	$V_0$ (mev)	$r_0$ (f)	$a_0$ (mev)	$K_V$ (mev)	$W_D$ (mev)	$T_1$ (f)	$a_1$ (f)	$V_{20}$ (f)	$U_0$ (f)	$U_2$ (f)	$\beta_2$	$\beta_{02}$
1st case	55.59	1.14	0.76	3.60	5.63	1.30	0.65	5.16	1.01	0.45	0.20	0.24
2nd case	"	"	"	"	4.50	"	"	"	"	"	"	0.616

#### References

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- See paper by G. G. Skobelev in *Proceedings of the International Conference on Nuclear Physics*, Prague, 1968, *Journal of Nuclear Energy*, **1968**, 103 (1968).
- For a discussion of the spin of the  $2^1_1$  state, see ref. 11.