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## *Lawrence Liverrncre Laboratory*

**PHOTODISSOCIATION RATE CALCULATIONS** 

**Frederick M. Luther Donald J. Wuebbles** 

**December 1976** 

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### **PHOTODISSOCIATION RATE CALCULATIONS"**

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December 1976

#### **INTRODUCTION**

**Stratospheric and tropospheric models that contain photochemically active species have almost universally invoked an assumption of pure absorption for computing solar fluxes and photodissociat ion rates. This has been due primarily to nothing more than limitations involving computer running time and capacity. For example, even with the simplest representations of atmospheric kinetics, models are often beset with computational constraints that normally lead to invoking the assumption of a purely absorbing atomic and molecular atmosphere for determining solar fluxes and the corresponding photodissociation rates. These limiting factors have been (and continue to be) sufficiently overriding that the roles of multiple scattering, the earth's surface reflection, clouds, and aerosols have remained a side issue in evaluating photochemical rates in atmospheric photochemical** 

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**trace gas models; yet it is well-known that these factors can be significant in determining stratospheric and tropospheric radiative intensities at photodissociative wavelengths.** 

**We will briefly review the approach used in pure molecular absorption calculations, and we will discuss the effect of including multiple scattering effects in a stratospheric model.** 

#### **RADIATIVE TRANSFER FORMULATIONS IN ATMOSPHERIC PHOTOCHEMICAL MODELS**

**The solution of the one-dimensional purely absorbing source-free**  radiative transfer equation at a particular altitude z<sub>n</sub> (km), solar zenith angle  $\theta_o$ , and atmospheric composition  $\{N_A(z)\}$  is given by

$$
F_{\lambda}(z_p, \theta_o, \{N_A\}, t) = F_{\lambda}(\infty) \exp\left[-\tau_{\lambda}(z_p, \theta_o, \{N_A\}, t)\right]
$$
 (1)

where  $F_{\lambda}$  d<sub> $\lambda$ </sub> is the flux of photons (in number per square centimeter per second) in the wavelength interval  $d\lambda$  about  $\lambda$ .  $F_{\lambda}(\omega)$  represents the solar **flux at 1 AU, and the optical depth**  $\tau_{\lambda}$  **is given by** 

$$
\tau_{\lambda}(z_{p}, \theta_{o}, \{N_{A}\}, t)
$$
\n
$$
\equiv \int_{z_{p}}^{\infty} dz \sum_{A} N_{A}(z, t) \sigma_{T}^{A} [\lambda, T(z)] \sec \theta_{o}(t) \qquad .
$$
\n(2)

**In Eq. (2) the summation on A includes all atmospheric absorbers, each**  having number density  $N_A(z,t)$  cm<sup>-3</sup> and a total absorption cross section **oT <sup>A</sup> [X,T(z)] cm<sup>2</sup> . Most generally, CTT [X,T(z)l is a function of the temperature T(z).** 

**The photodissociation rate for transforming species i to species j Is denoted by** 

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$$
J_{i+j}(z_p, \theta_0, t) N_i(z_p, t)
$$

**where** 

$$
J_{i+j}(z_p, \theta_o, t) = \int_{a11-\lambda} d\lambda \sigma_0^{-i} [j, \lambda, T(z_p)] F_{\lambda}(z_p, \theta_o, (N_A), t) \quad . \quad (3)
$$

The microscopic photodissociation cross section  $\sigma_{D}^{-1}\left[j,x,T(z_{n})\right]$  cm<sup>2</sup> is often written in terms of the so-called quantum yield  $Q_2(i+j)$  as

$$
\sigma_0^{-i} \left[ j, \lambda, T(z_p) \right] = \sigma_1^{-i} \left[ \lambda, T(z_p) \right] Q_{\lambda}(i+j) \tag{4}
$$

Given microscopic cross-section data,  $F_{\lambda}(\infty)$ ,  $(N_{A}(z))$ , and  $\infty$ , **it is a straightforward matter to compute photodissociation coefficients.** 

**It should be noted that uncertainties stiil remain in the data used to calculate the photodissociation coefficients in the models. For example, major uncertainties remain in the branching of 0, photolysis near 310 nm**  to either  $O({}^3P)$  or  $O({}^1D)$ , in the branching and quantum yield for NO<sub>3</sub> photo**lysis, and in the methodology for calculating the photolysis of species**  having banded or line absorption cross-section structures, such as 0<sub>2</sub> or NO.

## **RADIATIVE TRANSFER CALCULATIONS INCLUDING MOLECULAR MULTIPLE SCATTERING AND SURFACE ALBEDO**

**When molecular multiple scattering and surface albedo are included in the radiative transfer calculation, (I) is no longer the solution of the radiative transfer equation. However, the flux in the direct solar beam** 

 $F_1^S$  is given by an equation similar to  $(1)$ :

$$
F_{\lambda}^{S}(z_{p}, \theta_{o}, \{N_{A}\}, t) = F_{\lambda}(\infty) \exp\{-\tau_{\lambda}^{S}(z_{p}, \theta_{o}, \{N_{A}\}, t)\}
$$
 (5)

where the optical depth 
$$
r_{\lambda}
$$
 along the slant path is given by  
\n
$$
\tau_{\lambda}^{s}(z_{p}, \theta_{o}, (N_{A}), t)
$$
\n
$$
\equiv \int_{P}^{\infty} dz \sum_{A} N_{A}(z, t) \sigma_{T}^{A}(i, T(z)) \sec \theta_{o}(t)
$$
\n
$$
+ \int_{P}^{\infty} dz \sum_{i} N_{i}(z, t) \sigma_{R}^{i}(i) \sec \theta_{o}(t) \qquad (6)
$$

In (6) the summation on i includes  $\epsilon$ il atmospheric species, and  $\frac{1}{R}$  is the **Rayleigh scattering cross section for species i.**  $F_{\lambda}^{S}(z_{0}, e_{0}, (N_{A}^{+}, t))$  **differs** from  $F_{\lambda}$  ( $z_{\mathsf{p}}, \theta_{\mathsf{p}}, \mathsf{in}_{\mathsf{A}}$ ), t) as defined by (1) in that attenuation due to both **absorption and scattering is included in (5), whereas only absorption is included in (I).** 

**The photodissociation coefficient also depends upon the scattered (diffuse) radiation given by** 

$$
J_{i+j}(z_{p}, \theta_{o}, t) = \int_{a11-\lambda} d\lambda \cdot c_{0}^{-1} [j, \lambda, \tau(z_{p})]
$$
  
 
$$
\cdot \left[ F_{\lambda}^{-s}(z_{p}, \theta_{o}, (t_{A}), t) + \int_{\frac{1}{4} - \tau} \lambda(z_{p}, \omega) d\omega \right]
$$
 (7)

where  $\zeta_{\lambda}$  is the specific intensity of the diffuse radiation and  $\omega$  is a **solid angle. Including the effect of molecular multiple scattering and surface albedo in the calculation is simply expressed by changing the value**  of  $F_{\lambda}$  appearing in (3). For clarity we define

$$
F_{\lambda}^{\text{MS}}(\text{multiple scattering}) = F_{\lambda}^{5}(z_{p}, b_{p}, \{u_{A}\}, t) + \int_{\frac{1}{4} - \pi} I_{\lambda}(z_{p}, \omega) d\omega
$$
 (8)

**Henceforth, F PA (pure absorption) will refer to the flux defined by (1). Aside**  from substituting  $F_3$ <sup>MS</sup> in place of  $F_3$ <sup>PA</sup> in (3), all other aspects of the **photodissociation rate calculation are the same as for the purely absorbing**   $m$ olecular atmosphere. Since  $F_\lambda^S(z_p^{}, \theta_o^{}, \{N_A\}), t)$  is less than  $F_\lambda^{'PA}$ ,  $F_\lambda^{'MS}$ may be greater than or less than  $F^{\text{PA}}_{\lambda}$ , depending upon the intensity of the **diffuse radiation.** 

#### **MODEL DESCRIPTION**

**The effect of molecular scattering and surface albedo on ambient species concentrations and on model sensitivity has been assessed using a one-dimensional transport-kinetics model. The model has been described previously [Chang, 1975; Muebbles and Chang, 1975), so only the main features of the model will be summarized.** 

**The governing equation regarding the temporal variation in the number density of the i constituent c. is given by the continuity equation,** 

$$
\frac{\partial c_i}{\partial t} = P(c) - L(c) c_i - \frac{\partial}{\partial z} \left[ K_z \rho \frac{\partial}{\partial z} \left( \frac{c_i}{\rho} \right) \right]
$$
 (9)

where P(c) is the production of c<sub>1</sub> due to photochemical interactions of the other c<sub>;</sub> species; L(c) c<sub>;</sub> is the loss of c<sub>;</sub> due to chemical interaction of c<sub>i</sub> with the other c<sub>i</sub> species; K<sub>2</sub> is the vertical transport coefficient;  $\rho = \rho(z)$  is the ambient air density; and t and z are time and altitude, **respectively.** 

**The model atmosphere extends from the ground to 55 km and is divided into W layers. A total of 92 chemical (and photochemical) reactions are**  included in the model; 45 reactions describe  $0<sub>v</sub>$ , NO<sub>v</sub>, and HO<sub>v</sub> chemistry,

**and** *k7* **reactions describe CIO chemistry. These reactions are used to**  dynamically describe the stratospheric vertical distributions of 20 minor **dynamically describe the stratospheric vertical distributions of 20 minor**   $c10NO_2$ ,  $c10$ ,  $c10_2$ ,  $0c10$ ,  $c1NO$ ,  $c1NO_2$ ,  $Hc1$ ,  $cc1_4$ ,  $cr_2Cl_2$  and  $crc1_3$ ). Three **Species** (H, N, O(<sup>1</sup>D)) are assumed to be in instantaneous equilibrium, and the vertical distributions of  $N_2$ ,  $0_2$ ,  $H_2$ 0,  $CH_{\underline{h}}$  and  $H_2$  are assumed constant **the vertical distributions of N** *0* **HO , CH. and H, are assumed constant**  throughout the calculations.

called "eddy" mixing coefficient K<sub>r</sub> (see Wuebbles and Chang, 1975; Chang, **called "eddy" mixing coefficient K (see Wuebbles and Chang, 1975; Chang, 1975). The numerical technique, which is a variable order, multistep,**  et al. (1974). The boundary conditions are species dependent and are assumed to be either fixed or time varying source dependent concentrations at the surface and flux conditions at 55 km, the upper boundary.

The chemical reaction rates used in the model are shown in Table 1. Other reaction rates in the model are based on the review by Hampson and Garvin (1975). Detailed spectral data used in the calculations are as described by Gelinas (1974).

The effect of multiple scattering was incorporated into the photodissociation rate calculation by applying correction factors to the **flux F used in the pure absorption calculation (Eq. (1)). These MS PA correction factors, which are given by the ratio F /F , were computed for the unperturbed species profiles using a highly decailed solar radiation model (Luther and Gelinas, 1976). A separate factor was**  computed for each of the 44 levels and for each of the 148 wavelength **intervals between 133 and 735 nm. Different sets of correction factors** 

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were computed for each assumed value of surface albedo A<sub>c</sub> using a solar **zenith angle of ^5°• The correction factors were assumed to be constant in the calculations.** 

**Examples of correction factors for a solar zenith angle of 60° are shown in Figure 1 as functions of: wavelength for selected altitudes and various values of A . Since no si-lar radiation reaches the earth's surface at wavelengths less than 290 nm, there is no dependence on surface albedo in this spectral region. The results for a surface albedo of zero demonstrate the effect of molecular multiple scattering alone, which is shown to have its maximum effect near 330 nm. The correction factors are nearly constant with height above 20 km, but they may vary significantly with height in the region below 20 km, which is where most scettering events occur. Correction factors for a solar zenith angle of ^5° are somewhat larger than those shown for 60°, but they are qualitatively similar [Luther and Gelinas, 1976].** 

#### **RESULTS**

**The effects of including multiple scattering in the photodissociation rate calculation may be of three types: (1) the effect on photodissociation coefficients, (2) the effect on ambient species concentration profiles, and (3) the effect on model sensitivity to perturbations.** 

#### **1. Photodissociation Coefficients**

**As shown by Figure 1, the effect of multiple scattering should be largest for those species having strong absorption cross-sections at wavelengths greater than 300 nm. Table 2 compares photodissociation** 

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coefficients for pure absorption and for multiple scattering computed for the ambient model atmosphere. Only those photodissociation reactions significantly affected by multiple scattering are shown in Table 2. In an attempt to account for the diurnal variation of photodissociation in the model, the solar flux was halved in these calculations.

The importance of the significant changes in reaction rate coefficients is reflected in the species concentration profiles and in the model sensitivity.

2. Species Concentration Profiles

The concentration profiles for selected  $0_\chi,$  HO $_\chi,$  NO $_\chi,$  and ClX species are shown in Figures 2-5 for the ambient atmosphere prior to inclusion of multiple scattering and prior to the release of CFM's. These figures are included in order to define the reference conditions for assessing the fractional change in concentration caused by multiple scattering.

The changes in concentration of chemical species due to multiple scattering relative to the pure absorption calculation were computed for cases with and without ClONO<sub>2</sub>. Inclusion of ClONO<sub>2</sub> significantly affected the results for the chlorine containing species, but it had a small effect on the other species  $\langle\leq 5\%$  change in NO<sub>2</sub> and much less for other species). The following figures correspond to the case including C10NO<sub>2</sub> and with  $A_e = 0.25$  unless otherwise specified.

The effect of multiple scattering on  $0<sub>y</sub>$  species is shown in Figure 6. The large percentage increase in  $0(^{\dagger}$ D) near 10 km occurs where the ambient concentration is very small; nevertheless, it is significant in terms of stratospheric chemistry. The increases in  $0(^3P)$  and  $0(^{\text{1}}$   $0)$  in the region 20-30 km are due to increased photolysis of  $v_{3}$ . Because of differences in ambient concentrations, a small percentage decrease in 0, causes large percentage increases in the other species. The increase in  $O({}^3P)$  near 40 km is due

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primarily to increased photolysis of NO<sub>2</sub>. The increase in O<sub>3</sub> at this height occurs because the chemical production of  $0<sub>3</sub>$  by  $0 + 0<sub>2</sub>$  is increased more than the photolysis of  $0<sub>2</sub>$ .

Figure 7 shows the effect of multiple scattering on HO<sub>2</sub> species. The increase in OH results from HNO<sub>2</sub> "+ OH + NO<sub>2</sub> and H<sub>2</sub>O<sub>2</sub> "+ 2OH. The p $\bar{z}$ concentration in  $H_50_2$  occurs near 28 km, so a small percentage decrease in  $H_2O_2$  in this region can cause a large percentage increase in OH. The increase in  $HO_2$  is chemically linked to the increase in OH.

Figure 8 shows the effect of multiple scattering on NO<sub>v</sub> species. There is a very large increase in NO near 20 km due to increased photolysis of NO<sub>2</sub>.  $NO<sub>2</sub>$  increases because of increased photolysis of  $HNO<sub>2</sub>$ . There is very little  $HNO<sub>3</sub>$  above 30 km, so  $NO<sub>2</sub>$  decreases in this region because it is essentially the only source of NO.

The effect of multiple scattering on CIONO<sub>2</sub> is shown in Figure 9. The concentration is reduced 20-40£ between 20 and 30 km, which is the region of maximum C10NO<sub>2</sub> concentration. Photolysis of C10NO<sub>2</sub> affects several other chlorine containing species through a complex chain of reactions. Certain key reactions *are:* 

> CIONO, → CIO + NO,  $C10 + NO \rightarrow C1 + NO_2$  $CH + CH_{L} \rightarrow HCH + CH_{2}$

The peak concentration of C10NO<sub>2</sub> occurs near 25 km, so photolysis of C10NO<sub>2</sub> acts as a strong source of CIO in this region. The large increase in NO between 20-30 km tends to destroy C10, leading to a net decrease in C10 in this region (Figure 10) and an increase in CI (Figure 11). The increase in

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**CI leads to an increase in HCi (Figure 12) through reaction with CH<sub>1</sub>.** These results also indicate that inclusion of CIONO<sub>2</sub> has a more significant **effect on C1X species than the choice of surface albedo.** 

**Ratios of total column abundances computed with multiple scattering to**  that with pure absorption were  $0.94$   $(0_2)$ ,  $0.975$   $(N0_2)$  and 1.21  $(N0)$ .

**The above analysis is overly simplified considering the complexity and interaction of the various chemical cycles. Nevertheless, we have attempted to summarize the major mechanisms by which multiple scattering affects species concentrations.** 

**3. Model Sensitivity** 

**Model sensitivity was tested for two types of perturbations: release of CFM's at constant production levels and the stratospheric injection of NO**  In each case calculations were made with and without CIONO<sub>2</sub>. The CFM calculation to steady state assumes that CFC1<sub>3</sub> and CF<sub>2</sub>C1<sub>2</sub> are released at **1973 rates as estimated by McCarthy (1974). The orone reductions competed at steady state are shown in Table 3 for the various cases.** 

**Multiple scattering significantly reduces the sensitivity of the model without CIONO . The ozone reduction changed from -15-28/ to -12.45>**  when multiple scattering was included with  $A_{\rm s} = 0.25$ .

**The model with CI0N0\_ is the more complete model according to our present understanding. When multiple scattering is included in this**  model, there is a negligible change in model sensitivity for A<sub>s</sub> = 0.25 in **spite of significant changes in photodissociation rates and species concentrations. The decrease in sensitivity for the model without CI0N0 occurs because multiple scattering shifts the balance between NO<sub>2</sub> and NO** 

toward NO. This increases the importance of the reaction CIO + NO relative to C10 + 0, thus decreasing the effectiveness of the chlorine catalytic cycle. The effect of including CIONO<sub>2</sub> is also reduced by multiple scattering, This occurs primarily because of increased C10NO<sub>2</sub> photolysis reducing the C10NO<sub>2</sub> concentration. Thus, while the O<sub>2</sub> destructiveness of the CIO<sub>v</sub> catalytic cycle is reduced by multiple scattering, the ameliorating effect of CIONO<sub>2</sub> is also reduced. The net effect is that the CFM calculations including CIONO<sub>2</sub> with and without multiple scattering give approximately the same reduction in ozone. The model sensitivity depends upon the choice of A<sub>c</sub>, as indicated by the reduced model sensitivity for A<sub>c</sub> = 0.75.

Model sensitivity to stratospheric injections of NO<sub>2</sub> is shown in Table 4. For these calculations, NO<sub>x</sub> was injected globally at a rate of 2.5 x 10<sup>9</sup> kg/yr as NO<sub>2</sub> uniformly discributed over a 1-km thick shell centered at either 17 or 20 km altitude. Multiple scattering significantly increased the sensitivity of the model without C10NO<sub>2</sub>, but it had only a small effect  $\langle$ <52) on the of the model without C1ON0-, but it had only a smaii effect *(<\$%)* on the

#### DISCUSSION

The results above provide a partial assessment of effect of including multiple scattering in the photodissociat on rate calculation of transportkinetics models. Multiple scattering significantly affects photodissociation rates, particularly at wavelengths greater than 290 ma, and this is reflected in the species concentration profiles. The degree to which multiple scattering affects model sensitivity to various perturbations depends upon model chemistry, solar zenith angle, and surface albedo. The assessment has not yet been extended to diurnal or two-dimensional calculations.

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TABLE 1

1976 Chemistry:  $0_x$ , NO<sub>x</sub> and HO<sub>x</sub> Chemistry

Reaction Rate  $0^{3} + \mu^{3} + 0 + 0$ **0 )**   $0_3$  + hv + 0 + 0<sub>2</sub> ( 2 )  $0^{3} + \mu \wedge \nu (0) + 0^{5}$  (n) + 0<sup>3</sup>  $N_2$ 0 + 0( $\binom{1}{D}$  + 2N0 **NO + hv + N** + 0  $N + 0$ <sub>2</sub> + NO + 0  $N + NO + N_2 + O$  $o(^{1}D) + H_2O + 20H$  $o(^{1}v) + cH_4 + oH + CH_3$  $HNO_3 + hV + OH + NO_2$  $0_3$  + OH  $+$  HO<sub>2</sub> + C<sub>2</sub>  $0 + 0_2 + M + 0_3 + M$  $0 + 0_3 + 20_2$  1.9 x 10<sup>-11</sup> exp(-2300/T)  $N0<sub>2</sub> + h v + N0 + 0$  QJ(4)  $0_3$  + NO + NO<sub>2</sub> + 0<sub>2</sub>  $0 + N0<sub>2</sub> \rightarrow NO + 0<sub>2</sub>$  $N_2$ 0 + hv +  $N_2$  + 0(<sup>1</sup>0) QJ(5)  $N_2$ 0 + 0(<sup>1</sup>D) +  $N_2$  + 0<sub>2</sub>  $7 \times 10^{-11}$ **Q.)(6) <U(7)**   $0 + 0H + 0<sub>2</sub> + H$  $0_3$  + H0<sub>2</sub> + OH + 20<sub>2</sub> 1.0 x 10<sup>-13</sup> exp(-1250/T)  $0 + H0<sub>2</sub> + OH + 0<sub>2</sub>$   $3 \times 10^{-11}$  $H + 0$ <sub>2</sub> + M + H0<sub>2</sub> + M 2.08 x 10<sup>-32</sup> exp(290/T)  $0_3$  + H + OH +  $0_2$  **1.23** x 10<sup>-10</sup> exp(-562/T)

1.07 x  $10^{-34}$  exp(510/T) 9.0 x  $10^{-13}$ exp(-1200/T)  $9.1 \times 10^{-12}$  $7 \times 10^{-11}$ **1.1 x 10""tTexp(-3150/T) 2.7 x lo"1' 2.i x I0"'° 1.3 x 10"<sup>1</sup> <sup>0</sup>** 1.6  $\times$  10<sup>-12</sup>  $exp(-1000/T)$ <br>4.2  $\times$  10<sup>-11</sup>

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React ion  $H_2$  +  $H_2$  +  $H_2$   $_2$  +  $0$ <sub>2</sub>  $H_0 + H_2 + H_2 + H_2$ OH +  $NO_2$  + M + HNO<sub>3</sub> + M OH + HNO<sub>3</sub> + H<sub>2</sub>O + NO<sub>3</sub>  $H_2 O_2 + h\nu + 20H$  $H_2$ <sup>0</sup><sub>2</sub> + OH +  $H_2$ <sup>0</sup> + H0<sub>2</sub>  $N_2 + O({}^{1}0) + M + N_2 0 + M$  $N + NO_2 + N_2O + O$  $NO + O + M + NO_2 + M$  $NO + HO_2 + NO_2 + OH$  $H_2 + 0(^{7}D) + 0H + H$  $OH + OH \rightarrow H_2O + O$  $N + 0_3 + N0 + 0_2$  $NO_2 + O_3 + NO_3 + O_2$  $H0_{2}$  + hv + OH + 0 OH + CH<sub>4</sub> + H<sub>2</sub>O + CH<sub>3</sub>  $OH + OH + M + H<sub>2</sub>O<sub>2</sub> + M$  $H_2$ <sup>0</sup><sub>2</sub> + 0 + 0H + H0<sub>2</sub>  $0 + CH_4 + OH + CH_3$  $CO + OH + H + CO<sub>2</sub>$  $0(^{1}D) + M \rightarrow 0 + M$  $NO_3 + hv + NO_2 + 0$  $+ N0 + 0,$ 

*Rale*   $1.7 \times 10^{-11}$  cxp(-500/T)  $2.0 \times 10^{-11}$ 2.76 x 10 cxp(880/T) 1.166 x 10 exp(222/T) + H  $8.9 \times 10^{-14}$ QJ(8)  $1.7 \times 10^{-11}$  exp(-910/T)  $2.8 \times 10^{-36}$  $1.4 \times 10^{-12}$ 3.96 x 10 <sup>33</sup> exp(940/7)  $2.0 \times 10^{-13}$  $2.9 \times 10^{-10}$ 1.0 x 10<sup>-11</sup> exp(-550/T)  $5.7 \times 10^{-13}$  $1.2 x 10^{-12} exp(-2450/T)$ **<y(9)**  2.36 x 10 '<sup>-</sup> exp(-1710/T)  $2.5 \times 10^{-33}$  exp(2500/T)  $2.75 \times 10^{-12}$  exp(-2125/T)  $3.5 \times 10^{-11}$  exp(-4550/T)  $1.4 \times 10^{-13}$  $2.2 \times 10^{-11}$  exp(92/T) 0.66

0.34

 $-15-$ TABLE 1 1976 CHEM CONT

Reaction

 $C1 + 0_3 \div C10 + 0_2$  $C1 + OCl0 + 2C10$  $CI + O_2 + M + ClO_2 + M$  $c_1 + c_{H_{\frac{1}{4}}} + Hc_1 + c_{H_{\frac{3}{2}}}$  $c1 + c10<sub>2</sub> + c1<sub>2</sub> + 0<sub>2</sub>$  $C1 + C10<sub>2</sub> + 2C10$  $C1 + NO + M + CINO + M$  $C1 + C1N0 - C1<sub>2</sub> + N0$  $C1 + NO_2 + M + C1NC_2 + M$  $C1 + C1NO<sub>2</sub> + C1<sub>2</sub> + NO<sub>2</sub>$  $C10 + 0 + C1 + 0$ <sub>2</sub>  $NO + C10 + NO_2 + C1$  $10 + 0_3 + 10_2 + 0_2$  $c10 + 0_3 + 0c10 + 0_2$  $C10 + NO_2 \stackrel{M}{+} C1NO_3$  $C10 + C10 + C1 + OCl0$  $C10 + C10 + C1<sub>2</sub> + 0<sub>2</sub>$  $C10 + C10 + C1 + C10$ <sub>2</sub>  $HCl + O({}^{1}D) + Cl + OH$ CINO<sub>3</sub> + HCl + 0<sub>2</sub> + HNO<sub>3</sub>  $0H + HCl + H<sub>2</sub>0 + Cl$  $0 + HC1 \rightarrow OH + Cl$  $C10_2 + M + C1 + 0_2 + M$  $0 + 0010 + 010 + 0$ 

ClO<sub>x</sub> Chemistry Rate  $2.97 \times 10^{-11}$  exp(-243/T)  $5.9 \times 10^{-11}$  $1.7 \times 10^{-33} \left(\frac{300}{7}\right)$  $5.4 \times 10^{-12}$  exp(-1133/T)  $5 \times 10^{-11}$  $1.4 \times 10^{-12}$ 1.7 x 10<sup>-32</sup> exp(553/T)  $3.0 \times 10^{-11}$ 6.9 x  $10^{-34}$  exp(2115/T)  $3.0 \times 10^{-12}$  $3.38 \times 10^{-11}$  exp(+75/T) 1.13 x 10<sup>-11</sup> exp(+200/T)  $1.0 \times 10^{-12}$  exp(-2763/T)  $1.0 \times 10^{-12}$  exp(-2763/T) 0.05\*HNO<sub>3</sub> formation (incl. diurnal effect)  $2.0 \times 10^{-12}$  exp(-2300/T)  $2.0 \times 10^{-13}$  exp(-1260/T)  $2 \times 10^{-13}$  exp(-1260/T)  $2 \times 10^{-10}$  $0.0$  $2.0 \times 10^{-12}$  exp(-310/T) 1.75 x  $10^{-12}$  exp(-2273/T) 1.5 x  $10^{-8}$  exp(-4000/T)  $5.0 \times 10^{-13}$ 

### $-16-$ TABLE 1 1976 CHEM CONT

 $C10\text{ }$  Chemistry





### **TABLE 2**

### **COMPARISON OF PHOTODISSOCIATION RATES CALCULATED WITH AND WITHOUT MULTIPLE SCATTERING**





TABLE 3

 $-19-$ 

CHANGE IN THE OZONE COLUMN DUE TO THE ATMOSPHERIC RELEASE OF CFM'S -- STEADY STATE VALUE AT CONSTANT PRODUCTION



 $-20-$ 

## TABLE 4

## NO<sub>x</sub> INJECTION

## $2.5\times10^9$  kg/yr as  $\rm N\!O_{\rm Z}$





Fig. 1. Flux ratio versus wavelength for a solar zenith angle of 60° and selected values of surface albedo R at altitudes z of (a) 40 km, (b) 30 km, (c) 20 km, (d) 10 km, and (e) 0 km.

 $\bf 2.6$  $Z = 40$  km  $= 1.00$  $\overline{0.75}$ 0.50  $\overline{0.75}$  $0.00$  $0.2$  $0.0^{-1}$ żo  $30<sup>c</sup>$ 400 500 600  $700$ Wavelength  $-$ Fig. 1a  $2.8$ L, i.  $2.6$  $Z = 30$  km  $\frac{24}{22}$ Flux(mult\_scat.)/Flux(pure\_4bs.)  $r > 10$  $R = 1.00$  $2.0 0.75$  $1.8$  $0.50$  $1.6 1.4$  $\overline{0.25}$  $1.2$  $0.00$  $1.0$ <br> $0.8$  $0.6$ <br> $0.4$  $0.2$  $0.0$  $300 -$ 400 500 600 700 200 dength Waw  $-$  on

 $2.8$ 

# Fig. 16





**こうしょう** 





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**Fig. 6. The change in 0 species concentrations due to multiple scattering with As = 0.25.** 

· 1992年4月4日 · 1994年4月4日 · 1994年1月



 $\frac{d}{dx}$ 















Fig. 12. The change in concen-<br>tration of HCl due to multiple scattering.

 $\mathcal{A}$  <br>such that  $\mathcal{A}_\text{max}$  and  $\mathcal{A}_\text{max}$