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Lawrence Livermore Laboratory

PHOTODISSOCIATION RATE CALCULATIONS

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INTRODUCTION

Stratospheric and tropospheric models that contain photochemically active species have almost universally invoked an assumption of pure absorption for computing solar fluxes and photodissociation rates. This has been due primarily to nothing more than limitations involving computer running time and capacity. For example, even with the simplest representations of atmospheric kinetics, models are often beset with computational constraints that normally lead to invoking the assumption of a purely absorbing atomic and molecular atmosphere for determining solar fluxes and the corresponding photodissociation rates. These limiting factors have been (and continue to be) sufficiently overriding that the roles of multiple scattering, the earth's surface reflection, clouds, and aerosols have remained a side issue in evaluating photochemical rates in atmospheric photochemical

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trace gas models; yet it is well-known that these factors can be significant in determining stratospheric and tropospheric radiative intensities at photodissociative wavelengths.

We will briefly review the approach used in pure molecular absorption calculations, and we will discuss the effect of including multiple scattering effects in a stratospheric model.

RADIATIVE TRANSFER FORMULATIONS IN ATMOSPHERIC PHOTOCHEMICAL MODELS

The solution of the one-dimensional purely absorbing source-free radiative transfer equation at a particular altitude z_p (km), solar zenith angle θ_n , and atmospheric composition {N_A(z)} is given by

$$F_{\lambda}(z_{p},\theta_{o},\{N_{A}\}, t) = F_{\lambda}(\infty) \exp[-\tau_{\lambda}(z_{p},\theta_{o},\{N_{A}\},t)]$$
(1)

where $F_{\lambda} d\lambda$ is the flux of photons (in number per square centimeter per second) in the wavelength interval $d\lambda$ about λ . $F_{\lambda}(\infty)$ represents the solar flux at 1 AU, and the optical depth τ_{λ} is given by

$$\tau_{\lambda}(z_{p},\theta_{o}, [N_{A}], t)$$

$$\equiv \int_{z_{p}}^{\infty} dz \sum_{A} N_{A}(z, t) \sigma_{T}^{A} [\lambda, T(z)] \sec \theta_{o}(t) \qquad (2)$$

In Eq. (2) the summation on A includes all atmospheric absorbers, each having number density $N_A(z,t) \text{ cm}^{-3}$ and a total absorption cross section $\sigma_T^A[\lambda,T(z)] \text{ cm}^2$. Most generally, $\sigma_T^A[\lambda,T(z)]$ is a function of the temperature T(z).

The photodissociation rate for transforming species i to species j is denoted by

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$$J_{i \neq j}(z_p, \theta_o, t) \land (z_p, t)$$

where

$$J_{i \rightarrow j}(z_{p}, \theta_{o}, t) \equiv \int_{all \lambda} d\lambda \sigma_{D}^{i}[j, \lambda, T(z_{p})] F_{\lambda}(z_{p}, \theta_{o}, \{N_{A}\}, t) \quad . \quad (3)$$

The microscopic photodissociation cross section $\sigma_D^{i}[j,\lambda,T(z_p)] \text{ cm}^2$ is often written in terms of the so-called quantum yield $Q_{\lambda}(i+j)$ as

$$\sigma_{\mathsf{D}}^{i}[j,\lambda,\mathsf{T}(z_{p})] = \sigma_{\mathsf{T}}^{i}[\lambda,\mathsf{T}(z_{p})] Q_{\lambda}(i+j) \qquad (4)$$

Given microscopic cross-section data, $F_{\lambda}(\infty)$, $\{N_{A}(z)\}$, and ε_{o} , it is a straightforward matter to compute photodissociation coefficients.

It should be noted that uncertainties still remain in the data used to calculate the photodissociation coefficients in the models. For example, major uncertainties remain in the branching of 0_3 photolysis near 310 nm to either $0({}^{3}P)$ or $0({}^{1}D)$, in the branching and quantum yield for N0₃ photolysis, and in the methodology for calculating the photolysis of species having banded or line absorption cross-section structures, such as $0_2 \simeq N0$.

RADIATIVE TRANSFER CALCULATIONS INCLUDING MOLECULAR MULTIPLE SCATTERING AND SURFACE ALBEDO

When molecular multiple scattering and surface albedo are included in the radiative transfer calculation, (1) is no longer the solution of the radiative transfer equation. However, the flux in the direct solar beam F_{λ}^{s} is given by an equation similar to (1):

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$$F_{\lambda}^{s}(z_{p},\theta_{o}, \{N_{A}\}, t) = F_{\lambda}(\infty) \exp\left[-\tau_{\lambda}^{s}(z_{p},\theta_{o}, \{N_{A}\}, t)\right]$$
(5)

where the optical depth
$$t_{\lambda}$$
 along the shaft path is given by

$$\tau_{\lambda}^{S}(z_{p}, \theta_{0}, \{N_{A}\}, t)$$

$$\equiv \int_{z_{p}}^{\infty} dz \sum_{A} N_{A}(z, t) \sigma_{T}^{A}[i, T(z)] \sec \theta_{0}(t)$$

$$+ \int_{z_{p}}^{\infty} dz \sum_{i} N_{i}(z, t) \sigma_{R}^{i}(\lambda) \sec \theta_{0}(t) \qquad (6)$$

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In (6) the summation on i includes cill atmospheric species, and $\frac{1}{R}^{i}$ is the Rayleigh scattering cross section for species i. $F_{\lambda}^{S}(z_{p}, \theta_{o}, \{N_{A}^{i}, t\})$ differs from $F_{\lambda}(z_{p}, \theta_{o}, \{N_{A}^{i}, t\})$ as defined by (1) in that attenuation due to both absorption and scattering is included in (5), whereas only absorption is included in (1).

The photodissociation coefficient also depends upon the scattered (diffuse) radiation given by

$$J_{i \neq j}(z_{p}, \theta_{0}, t) = \int_{all \lambda} d\lambda e_{0}^{-1}[j_{i}\lambda_{i}\tau(z_{p})]$$

$$+ \left[F_{\lambda}^{-S}(z_{p}, \theta_{0}, (B_{A}), t) + \int_{A_{i}} I_{\lambda}(z_{p}, \omega) d\omega\right]$$
(7)

where z_{λ} is the specific intensity of the diffuse radiation and ω is a solid angle. Including the effect of molecular multiple scattering and surface albedo in the calculation is simply expressed by changing the value of F_{λ} appearing in (3). For clarity we define

$$F_{\lambda}^{\text{MS}}$$
(multiple scattering) $\notin F_{\lambda}^{-5}(z_{p}, v_{\sigma}, \{H_{A}\}, t) + \int_{4} I_{\lambda}(z_{p}, \omega) d\omega$ (8)

Henceforth, F_{λ}^{PA} (pure absorption) will refer to the flux defined by (1). Aside from substituting F_{λ}^{MS} in place of F_{λ}^{PA} in (3), all other aspects of the photodissociation rate calculation are the same as for the purely absorbing molecular atmosphere. Since $F_{\lambda}^{S}(z_{p}, \theta_{o}, \{N_{A}\}, t)$ is less than F_{λ}^{PA} , F_{λ}^{MS} may be greater than or less than F_{λ}^{PA} , depending upon the intensity of the diffuse radiation.

MODEL DESCRIPTION

The effect of molecular scattering and surface albedo on ambient species concentrations and on model sensitivity has been assessed using a one-dimensional transport-kinetics model. The model has been described previously [Chang, 1975; Wuebbles and Chang, 1975], so only the main features of the model will be summarized.

The governing equation regarding the temporal variation in the number density of the i^{th} constituent c_i is given by the continuity equation,

$$\frac{\partial c_i}{\partial t} = P(c) - L(c) c_i - \frac{\partial}{\partial z} \left[K_z \rho \frac{\partial}{\partial z} \left(\frac{c_i}{\rho} \right) \right]$$
(9)

where P(c) is the production of c_i due to photochemical interactions of the other c_j species; L(c) c_i is the loss of c_i due to chemical interaction of c_i with the other c_j species; K_z is the vertical transport coefficient; a = a(z) is the ambient air density; and t and z are time and altitude, respectively.

The model atmosphere extends from the ground to 55 km and is divided into 44 layers. A total of 92 chemical (and photochemical) reactions are included in the model; 45 reactions describe 0_{y} , $N0_{y}$, and $H0_{y}$ chemistry, and 47 reactions describe ClO_{x} chemistry. These reactions are used to dynamically describe the stratospheric vertical distributions of 20 minor atmospheric species $(0({}^{3}\text{P}), 0_{3}, N0, N0_{2}, N_{2}0, HN0_{3}, 0H, H0_{2}, H_{2}0_{2}, Cl, Cl0N0_{2}, Cl0, Cl0_{2}, 0Cl0, ClN0, ClN0_{2}, HCl, CCl_{4}, CF_{2}Cl_{2} and CFCl_{3}).$ Three species (H, N, $0({}^{1}\text{D})$) are assumed to be in instantaneous equilibrium, and the vertical distributions of N₂, 0₂, H₂0, CH₄ and H₂ are assumed constant throughout the calculations.

The vertical transport in the model is parameterized through the socalled "eddy" mixing coefficient K_z (see Wuebbles and Chang, 1975; Chang, 1975). The numerical technique, which is a variable order, multistep, implicit method, used to solve Equation (1), is that described by Chang et al. (1974). The boundary conditions are species dependent and are assumed to be either fixed or time varying source dependent concentrations at the surface and flux conditions at 55 km, the upper boundary.

The chemical reaction rates used in the model are shown in Table 1. Other reaction rates in the model are based on the review by Hampson and Garvin (1975). Detailed spectral data used in the calculations are as described by Gelinas (1974).

The effect of multiple scattering was incorporated into the photodissociation rate calculation by applying correction factors to the flux F_{λ} used in the pure absorption calculation (Eq. (1)). These correction factors, which are given by the ratio $F_{\lambda}^{MS}/F_{\lambda}^{PA}$, were computed for the unperturbed species profiles using a highly detailed solar radiation model (Luther and Gelinas, 1976). A separate factor was computed for each of the 44 levels and for each of the 148 wavelength intervals between 133 and 735 nm. Different sets of correction factors

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were computed for each assumed value of surface albedo A_s using a solar zenith angle of 45°. The correction factors were assumed to be constant in the calculations.

Examples of correction factors for a solar zenith angle of 60° are shown in Figure 1 as functions of wavelength for selected altitudes and various values of A_s . Since no sclar radiation reaches the earth's surface at wavelengths less than 290 nm, there is no dependence on surface albedo in this spectral region. The results for a surface albedo of zero demonstrate the effect of molecular multiple scattering alone, which is shown to have its maximum effect near 330 nm. The correction factors are nearly constant with height above 20 km, but they may vary significantly with height in the region below 20 km, which is where most scattering events occur. Correction factors for a solar zenith angle of 45° are somewhat larger than those shown for 60° , but they are qualitatively similar [Luther and Gelinas, 1976].

RESULTS

The effects of including multiple scattering in the photodissociation rate calculation may be of three types: (1) the effect on photodissociation coefficients, (2) the effect on ambient species concentration profiles, and (3) the effect on model sensitivity to perturbations.

1. Photodissociation Coefficients

As shown by Figure 1, the effect of multiple scattering should be largest for those species having strong absorption cross-sections at wavelengths greater than 300 nm. Table 2 compares photodissociation

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coefficients for pure absorption and for multiple scattering computed for the ambient model atmosphere. Only those photodissociation reactions significantly affected by multiple scattering are shown in Table 2. In an attempt to account for the diurnal variation of photodissociation in the model, the solar flux was halved in these calculations.

The importance of the significant changes in reaction rate coefficients is reflected in the species concentration profiles and in the model sensitivity.

2. Species Concentration Profiles

The concentration profiles for selected 0_x , $H0_x$, $N0_x$, and ClX species are shown in Figures 2-5 for the ambient atmosphere prior to inclusion of multiple scattering and prior to the release of CFM's. These figures are included in order to define the reference conditions for assessing the fractional change in concentration caused by multiple scattering.

The changes in concentration of chemical species due to multiple scattering relative to the pure absorption calculation were computed for cases with and without ClONO_2 . Inclusion of ClONO_2 significantly affected the results for the chlorine containing species, but it had a small effect on the other species (< 5% change in NO₂ and much less for other species). The following figures correspond to the case including ClONO_2 and with A_e = 0.25 unless otherwise specified.

The effect of multiple scattering on 0_x species is shown in Figure 6. The large percentage increase in $0({}^{1}D)$ near 10 km occurs where the ambient concentration is very small; nevertheless, it is significant in terms of stratospheric chemistry. The increases in $0({}^{3}P)$ and $0({}^{1}D)$ in the region 20-30 km are due to increased photolysis of v_3 . Because of differences in ambient concentrations, a small percentage decrease in 0_3 causes large percentage increases in the other species. The increase in $0({}^{3}P)$ near 40 km is due

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primarily to increased photolysis of NO₂. The increase in O₃ at this height occurs because the chemical production of O₃ by 0 + O₂ is increased more than the photolysis of O₂.

Figure 7 shows the effect of multiple scattering on HO_x species. The increase in OH results from HNO₃ $\stackrel{h.y}{\rightarrow}$ OH + NO₂ and H₂O₂ $\stackrel{h.y}{\rightarrow}$ 2OH. The peak concentration in H₂O₂ occurs near 28 km, so a small percentage decrease in H₂O₂ in this region can cause a large percentage increase in OH. The increase in HO₂ is chemically linked to the increase in OH.

Figure 8 shows the effect of multiple scattering on NO_x species. There is a very large increase in NO near 20 km due to increased photolysis of NO_2 . NO_2 increases because of increased photolysis of HNO_3 . There is very little HNO_3 above 30 km, so NO_2 decreases in this region because it is essentially the only source of NO.

The effect of multiple scattering on ClONO₂ is shown in Figure 9. The concentration is reduced 20-40% between 20 and 30 km, which is the region of maximum ClONO₂ concentration. Photolysis of ClONO₂ affects several other chlorine containing species through a complex chain of reactions. Certain key reactions are:

 $clono_{2} \xrightarrow{h_{+}} clo + no_{2}$ $clo + no \rightarrow cl + no_{2}$ $cl + ch_{4} \rightarrow hcl + ch_{3}$

The peak concentration of $C10NO_2$ occurs near 25 km, so photolysis of $C10NO_2$ acts as a strong source of C10 in this region. The large increase in NO between 20-30 km tends to destroy C10, leading to a net decrease in C10 in this region (Figure 10) and an increase in C1 (Figure 11). The increase in

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Cl leads to an increase in HCl (Figure 12) through reaction with CH_{4} . These results also indicate that inclusion of $ClONO_{2}$ has a more significant effect on ClX species than the choice of surface albedo.

Ratios of total column abundances computed with multiple scattering to that with pure absorption were 0.94 (0_3) , 0.975 $(N0_2)$ and 1.21 (N0).

The above analysis is overly simplified considering the complexity and interaction of the various chemical cycles. Nevertheless, we have attempted to summarize the major mechanisms by which multiple scattering affects species concentrations.

3. Model Sensitivity

Model sensitivity was tested for two types of perturbations: release of CFM's at constant production levels and the stratospheric injection of NO_{χ} . In each case calculations were made with and without $CIONO_{\chi}$. The CFM calculation to steady state assumes that CFCI₃ and CF₂CI₂ are released at 1973 rates as estimated by McCarthy (1974). The ozone reductions computed at steady state are shown in Table 3 for the various cases.

Multiple scattering significantly reduces the sensitivity of the model without C10NO_2 . The ozone reduction changed from -15.28% to -12.45% when multiple scattering was included with $A_{\rm S}$ = 0.25.

The model with $llono_2$ is the more complete model according to our present understanding. When multiple scattering is included in this model, there is a negligible change in model sensitivity for $A_s = 0.25$ in spite of significant changes in photodissociation rates and species concentrations. The decrease in sensitivity for the model without $llono_2$ occurs because multiple scattering : hifts the balance between NO₂ and NO toward NO. This increases the importance of the reaction C10 + NO relative to C10 + 0, thus decreasing the effectiveness of the chlorine catalytic cycle. The effect of including $C10NO_2$ is also reduced by multiple scattering. This occurs primarily because of increased $C10NO_2$ photolysis reducing the $C10NO_2$ concentration. Thus, while the O_3 destructiveness of the $C10_x$ catalytic cycle is reduced by multiple scattering, the ameliorating effect of $C10NO_2$ is also reduced. The net effect is that the CFM calculations including $C10NO_2$ with and without multiple scattering give approximately the same reduction in ozone. The model sensitivity depends upon the choice of A₂, as indicated by the reduced model sensitivity for A₂ = 0.75.

Model sensitivity to stratospheric injections of NO_x is shown in Table 4. For these calculations, NO_x was injected globally at a rate of 2.5 x 10^9 kg/yr as NO_2 uniformly discributed over a 1-km thick shell centered at either 17 or 20 km altitude. Multiple scattering significantly increased the sensitivity of the model without ClONO₂, but it had only a small effect (<5%) on the sensitivity of the model with ClONO₂.

DISCUSSION

The results above provide a partial assessment of effect of including multiple scattering in the photodissociat on rate calculation of transportkinetics models. Multiple scattering significantly affects photodissociation rates, particularly at wavelengths greater than 290 nm, and this is reflected in the species concentration profiles. The degree to which multiple scattering affects model sensitivity to various perturbations depends upon model chemistry, solar zenith angle, and surface albedo. The assessment has not yet been extended to diurnal or two-dimensional calculations.

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TABLE 1

1976 Chemistry: 0_x, NO_x and HO_x Chemistry

Rate Reaction QJ(1) $0_{2} + h_{v} \rightarrow 0 + 0$ $0_3 + hv \rightarrow 0 + 0_7$ QJ(2) $0_3 + hv + 0(10) + 0_2$ QJ (3) 1.07 x 10⁻³⁴ exp(510/T) $0 + 0_2 + M + 0_3 + M$ $1.9 \times 10^{-11} \exp(-2300/T)$ $0 + 0_3 + 20_2$ QJ(4) $NO_2 + hv \rightarrow NO + O$ 9.0 x 10⁻¹³exp(-1200/T) $0_3 + NO + NO_2 + 0_2$ 9.1 × 10⁻¹² $0 + NO_2 \rightarrow NO + O_2$ $N_2^0 + hv + N_2 + 0(^1 D)$ QJ (5) 7 × 10⁻¹¹ $N_{2}0 + 0(^{1}D) + N_{2} + 0_{2}$.7 × 10⁻¹¹ $N_{0}0 + 0(^{1}D) + 2NO$ Q.J(6) $NO + hv \rightarrow N + O$ 1.1 x 10⁻¹⁴ T exp(-3150/T) $N + 0_2 + N0 + 0$ 2.7×10^{-11} $N + NO \rightarrow N_2 + 0$ 2.1 × 10⁻¹⁰ $0(^{1}D) + H_{2}O \rightarrow 20H$ 1.3 × 10⁻¹⁰ $O(^{1}D) + CH_{4} \rightarrow OH + CH_{3}$ QJ (7) HNO3 + hv → OH + NO2 $1.6 \times 10^{-12} \exp(-1000/T)$ $0_3 + 0H \rightarrow H0_2 + C_2$ 4.2 × 10⁻¹¹ $0 + 0H \div 0_2 + H$ 1.0 x 10⁻¹³ exp(-1250/T) $0_3 + H0_2 \rightarrow OH + 20_2$ 3 × 10⁻¹¹ $0 + HO_2 \rightarrow OH + O_2$ 2.08 x 10⁻³² exp(290/T) $H + 0_2 + M \rightarrow H0_2 + M$ 1.23 x 10⁻¹⁰ exp(-562/T) $0_3 + H \rightarrow 0H + 0_2$

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Reaction $HO_2 + HO_2 \rightarrow H_2O_2 + O_2$ $HO_2 + OH \rightarrow H_2O + O_2$ $OH + NO_2 + M \rightarrow HNO_3 + M$ $OH + HNO_3 + H_2O + NO_3$ $H_20_2 + hv \rightarrow 20H$ $H_20_2 + OH \rightarrow H_20 + H0_2$ $N_2 + O(10) + M + N_20 + M$ $N + NO_2 \rightarrow N_2O + O$ $NO + O + M + NO_2 + M$ $NO + HO_2 \rightarrow NO_2 + OH$ $H_2 + O(^{\dagger}D) \rightarrow OH + H$ $OH + OH \rightarrow H_2O + O$ $N + 0_3 \rightarrow NO + 0_2$ $NO_2 + O_3 \rightarrow NO_3 + O_2$ $H0_2 + hv \rightarrow 0H + 0$ $OH + CH_4 \rightarrow H_2O + CH_3$ $OH + OH + M + H_2O_2 + M$ $H_2^0_2 + 0 \rightarrow 0H + H_2^0_2$ $0 + CH_4 \rightarrow OH + CH_3$ $CO + OH \rightarrow H + CO_2$ $O(^{1}D) + M \rightarrow O + M$ $NO_3 + hv \rightarrow NO_2 + 0$ + NO + 02 0.34

Rate $1.7 \times 10^{-11} \exp(-500/T)$ 2.0×10^{-11} $\frac{2.76 \times 10^{-13} \exp(880/T)}{1.166 \times 10^{18} \exp(222/T) + M}$ 8.9 × 10⁻¹⁴ QJ(8) $1.7 \times 10^{-11} \exp(-910/T)$ 2.8 × 10⁻³⁶ 1.4×10^{-12} $3.96 \times 10^{-33} \exp(940/7)$ 2.0×10^{-13} 2.9×10^{-10} 1.0 x 10⁻¹¹ exp(-550/T) 5.7×10^{-13} $1.2 \times 10^{-13} \exp(-2450/T)$ QJ (9) 2.36 × 10⁻¹² exp(-1710/T) 2.5 x 10⁻³³ exp(2500/T) $2.75 \times 10^{-12} \exp(-2125/T)$ $3.5 \times 10^{-11} \exp(-4550/T)$ 1.4×10^{-13} $2.2 \times 10^{-11} \exp(92/T)$ 0.66

-15-TABLE 1 1976 CHEM CONT

Reaction

 $c_1 + 0_3 \rightarrow c_{10} + 0_2$ C1 + OC10 + 2C10 $C1 + 0_2 + M + C10_2 + M$ $c_1 + c_{4} + Hc_1 + c_{3}$ $c_1 + c_{10_2} + c_{1_2} + o_2$ $c1 + c10_2 + 2c10$ C1 + NO + M + CINO + M $c1 + c1N0 - c1_2 + N0$ $C1 + NO_2 + M + CINC_2 + M$ $c_1 + c_{1NO_2} \rightarrow c_2 + NO_2$ $C10 + 0 + C1 + 0_{2}$ $NO + CIO \rightarrow NO_2 + CI$ $c_{10} + 0_3 + c_{10}2 + 0_2$ $c_{10} + 0_3 + 0c_{10} + 0_2$ $c_{10} + no_2 \stackrel{M}{\rightarrow} c_{1no_3}$ C10 + C10 + C1 + OC10 $c_{10} + c_{10} + c_{12} + 0_2$ $C10 + C10 + C1 + C10_{2}$ $HC1 + O(^{1}D) + C1 + OH$ $C1NO_3 + HC1 + O_2 + HNO_3$ $OH + HC1 \rightarrow H_2O + C1$ 0 + HC1 → 0H + C1 $C10_2 + M + C1 + 0_2 + M$ $0 + 0010 + 010 + 0_{2}$

Clo, Chemistry Rate 2.97 x 10⁻¹¹ exp(-243/T) 5.9 × 10⁻¹¹ $1.7 \times 10^{-33} (\frac{300}{7})$ 5.4 x 10⁻¹² exp(-1133/T) 5 x 10⁻¹¹ 1.4×10^{-12} 1.7 x 10⁻³² exp(553/T) 3.0 × 10⁻¹¹ 6.9 x 10⁻³⁴ exp(2115/T) 3.0×10^{-12} 3.38 x 10⁻¹¹ exp(+75/T) 1.13 x 10⁻¹¹ exp(+200/T) 1.0 x 10⁻¹² exp(-2763/T) 1.Q x 10⁻¹² exp(-2763/T) 0.05*HNO₃ formation (incl. diurnal effect) 2.0 × 10⁻¹² exp(-2300/T) 2.0 x 10⁻¹³ exp(-1260/T) 2 x 10⁻¹³ exp(-1260/T) 2×10^{-10} 0.0 2.0 x 10⁻¹² exp(-310/T) 1.75 x 10⁻¹² exp(-2273/T) 1.5 x 10⁻⁸ exp(-4000/T) 5.0×10^{-13}

-16-TABLE 1 1976 CHEM CONT

ClO_x Chemistry

Reaction	Rate
$NO + OCIO + NO_2 + CIO$	3.4×10^{-13}
N + 0C10 + N0' + C10	6.0×10^{-13}
H + OC10 + OH + C10	5.7×10^{-11}
C1 + OH + HC1 + O	2.0 x 10 ⁻¹² exp(-1878/T)
$C1 + HO_2 \rightarrow HC1 + O_2$	3.0×10^{-11}
$C1 + HNO_3 + HC1 + NO_3$	$4.0 \times 10^{-12} \exp(-1500/T)$
$C10_2 + H0_2 + HCi + 20_2$	3.0×10^{-12}
$Cl_2 + hv \rightarrow 2Cl$	QCJ (1)
HCl + $h_v \rightarrow H + Cl$	QCJ (2)
$C10_2 + h_v + C10 + O(^1D)$	QCJ(3)
$C10 + h_v \rightarrow C1 + 0$	QCJ (4)
$C10 + h_v + C1 + 0(^1D)$	QCJ (5)
C1NO + hv + C1 + NO	QCJ (6)
$CINO_2 + h_v + CI + NO_2$	QCJ(7)
$0C10 + h_v \rightarrow C10 + 0(^1D)$	QCJ (8)
$0C10 + h_v + C10 + 0$	QCJ (9)
$CF_{2}Cl_{2} + h_{v} + 2Cl$	QCJ (10)
$CFCl_3 + h_v \neq 2.5 Cl$	QCJ (11)
CCl ₄ + hv → 2Cl	QCJ(12)
$CFC1_3 + O(^1D) \rightarrow 2C1$	5.8×10^{-10}
$CF_2C1_2 + O(^1D) \rightarrow 2C1$	5.3×10^{-10}
C1 + H ₂ + HC1 + H	5.7 x 10 ⁻¹¹ exp(-2400/T)
$C1 + H_20_2 + HC1 + H0_2$	1.0 x 10 ⁻¹¹ exp(-810/T)
$0 + C1NO_3 + C10 + NO_3$	2.1 x 10^{-13}
$0H + CH_3C1 + H_2O + HO_2 + HC1$	1.58 x 10 ⁻¹² exp(-1049/T)

Altitude, km	J _{PA}	J _{MS}	MS ^J MS ^{/J} PA	
$0_3 + hv + 0(^{3}P) + 0_2$				
10	2.02×10^{-4}	3.01×10^{-4}	1.49	
20	2.07	3.00	1.45	
30	2.41	3.21	1.33	
40	2.74	3.49	1.27	
$0_3 + hv \neq 0(^{1}D) + 0_2$				
10	5.83×10^{-6}	1.08 × 10 ⁻⁵	1.85	
20	7.97×10^{-6}	1.22×10^{-5}	1.53	
30	5.78 × 10 ⁻⁵	6.27×10^{-5}	1.08	
40	9.32×10^{-4}	9.14 × 10 ⁻⁴	0.98	
$NO_2 + hv + NO + O$				
10	4.72	7.77×10^{-3}	1.65	
20	4.74	7.76	1.64	
30	4.85	7.68	1.58	
40	4.97	7.73	1.56	
$HNO_3 + hv \rightarrow OH + NO_2$				
10	2.72 × 10 ⁻⁷	4.82×10^{-7}	1.77	
20	3.37×10^{-7}	5.18 × 10 ⁻⁷	1.54	
30	4.98×10^{-6}	5.11 × 10 ⁻⁶	1.03	
40	3.52×10^{-5}	3.52×10^{-5}	1.00	

TABLE 2

COMPARISON OF PHOTODISSOCIATION RATES CALCULATED WITH AND WITHOUT MULTIPLE SCATTERING

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Altitude, km	J _{PA} J _{MS}		J _{MS} /J _{PA}	
$H_20_2 + hv + 20H$				
10	1.90 x 10 ⁻⁶	3.35×10^{-6}	1.76	
20	2.23×10^{-6}	3.46×10^{-6}	1.55	
30	6.57×10^{-6}	7.29 × 10 ⁻⁶	1.11	
40	3.51 × 10 ⁻⁵	3.52×10^{-5}	1.00	
Cl0 + hv → Cl + O			1	
10	1.72×10^{-5}	3.25×10^{-5}	1.89	
20	2.55×10^{-5}	3.82×10^{-5}	1.50	
30	1.92×10^{-4}	2.05×10^{-4}	1.07	
40	1.64×10^{-3}	1.62×10^{-3}	0.99	
C10N0 ₂ + hv → C10 + NO ₂				
10	3.31 × 10 ⁻⁵	5.57 × 10 ⁻⁵	1.68	
20	3.43×10^{-5}	5.61 × 10 ⁻⁵	1.64	
30	5.22×10^{-5}	7.07 × 10 ⁻⁵	1.35	
40	2.24×10^{-4}	2.38×10^{-4}	1.06	

TABLE 3

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CHANGE IN THE OZONE COLUMN DUE TO THE ATMOSPHERIC RELEASE OF CFM'S -- STEADY STATE VALUE AT CONSTANT PRODUCTION

	Change in Ozone Column - %		
CASE	WITHOUT MULTIPLE SCATTERING	WITH MULTIPLE SCATTERING	RATIO
		A _s = 0.25	
WITHOUT CLONO2	-15,28	-12,45	0.81
WITH CLONO2	-8.97	-9.03	1.01
_		A _s = 0,75	
WITH CLONO2	-8.97	-6.91	0.77
_			

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TABLE 4

NO_x INJECTION

2.5 x 10^9 kg/yr as $N0_2$

	Change in Ozone Column - %			
CASE	WITHOUT MULTIPLE SCATTERING	MULTIPLE SCATTERING A _S = 0,25	RATIO	
17-KM INJECTION:				
WITHOUT CLONO2	-0.75	-0.97	1.29	
WITH CLONO2	-1.07	-1.10	1.03	
20-KM INJECTION:				
WITHOUT CLIONO2	-3.23	-3.82	1.18	
WITH CIONO2	-3,93	-4.12	1.35	



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Fig. 1. Flux ratio versus wavelength for a solar zenith angle of 60° and selected values of surface albedo R at altitudes z of (a) 40 km. (b) 30 km, (c) 20 km, (d) 10 km, and (e) 0 km.







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Fig. 4. Ambient concentration profiles of NO₂, NO, and UNO₃.



Fig. 5. Ambient concentration profiles of species containing chlorine.



Fig. 6. The change in 0 species concentrations due to multiple scattering with $A_s = 0.25$.

















Fig. 12. The change in concentration of HCl due to multiple scattering.

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