

PHOTOPRODUCTION MEASUREMENTS OF  $\pi^+$  NEAR THRESHOLD ON  ${}^3\text{He}$  AND  ${}^2\text{H}$

P. Argan, G. Audit, A. Bloch, N. de Botton, C. Schuhl,  
G. Tamas, C. Tsara, E. Vincent  
DPH-I, JEN Saclay, BP 2, 91191 Gif-sur-Yvette, France

J. Deusch, D. Favart, R. Prieels, B. Van Oystaeyen  
Institut de Physique Corpusculaire, Université Catholique de Louvain  
B-1348 Louvain-la-Neuve, Belgium

The two- and three-nucleon systems offer the ground for a significant study of nuclear structures generated by the nucleon-nucleon interaction. We present here a contribution to the knowledge of these systems, based on the fact that charged pion photoproduction at threshold is a weak probe measuring the spin-flip form factor of the nuclear target. We have thus measured  $\sigma^{\pi^+}$  the cross-sections for reactions :

$$\gamma + {}^3\text{He} \rightarrow {}^3\text{H} + \pi^+, \quad (1)$$

$$\gamma + {}^2\text{H} \rightarrow n + n + \pi^+, \quad (2)$$

relatively to the elementary one :

$$\gamma + p \rightarrow n + \pi^+, \quad (3)$$

at energies ranging from 1 to 5 MeV above threshold.

From these data, we extract the ratio of the slopes of the cross sections at threshold for reactions (1) and (3) :

$$\frac{a({}^3\text{He})}{a(p)} = \lim_{q \rightarrow 0} \frac{\sigma({}^3\text{He}) k/q_s}{\sigma(p) k, q} = 0.62 \pm 0.02.$$

As for reaction (2), because of the three-body final state, the cross section has a more complicated shape, but it can nevertheless be characterized by one number playing a role equivalent to the slope  $a(A)$ , and which was measured with a  $\pm 2.5\%$  accuracy.

In order to extract from the data a form factor, it was necessary to estimate correctly the modification to a plane wave calculation due to the pion rescattering. Using a multiple scattering expansion<sup>3</sup>, we found a small effect ( $-1\%$ ) for reaction (1) and a larger one ( $+10\%$ ) for reaction (2). By applying the correction to the data on reaction (1), we obtain

$$|F(Q^0)|^2 = 0.52 \pm 0.02. \quad (4)$$

This is 20% lower than the form factor squared extracted from the magnetic electron scattering on  ${}^3\text{He}$  and  ${}^2\text{H}$ . Calculations based on realistic wave functions<sup>4</sup> yield form factors in general agreement

with the experimental one (4). These calculations corrected by the effect of the nucleon fermi-motion should result in a slightly lower threshold cross section improving the overall agreement with the experiment.

The cross section for reaction (2) has been computed<sup>5</sup> including the effect of the fermi motion of the nucleon, but with less scrutiny concerning the 2N wave functions. Once corrected for the pion multiple scattering, they fall 10 % above the measured cross section. The cross section for the threshold electrodisintegration :

$$e + d \rightarrow e' + n + p$$

(which measures the spin flip transition form factor). When computed in impulse approximation is about 20 % lower than the measured one.

Thus we observe in the case of the two- and three-nucleon targets that their spin-flip form factor as measured by threshold photoproduction is 11-15 % lower than the one measured by magnetic electron scattering, and agrees better with impulse approximation calculations. It indicates the following trend :

- pion photoproduction measures essentially the one-body spin flip form factor ;
- magnetic electron scattering is more sensitive to many-body corrections.

#### REFERENCES

1. P. Argan et al., to be published.
2. E.C. Booth et al., Phys. Lett. 66B, 236 (1977).  
G. Audit et al., Phys. Rev. C16, 1517 (1977).
3. N. de Botton and C. Tzara, Rapport interne DPh-N/HE/78-06.
4. A.C. Phillips and F. Roig, Nucl. Phys. A234, 378 (1974)  
B. Gouliard, A. Laverne and J.D. Vergados, Phys. Rev. C18, 944 (1978).
5. J.V. Noble, Phys. Lett. 67B, 39 (1977).  
J.M. Lager, Nucl. Phys. A296, 388 (1977).