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AN INVESTIGATION OF ASSOCIATIVE MULTIPLICITY IN THE REACTION $\overline{p}p \longrightarrow p$ (SLOW) + X AT 22.4 GeV/c

Alma-Ata - Dubna - Helsinki - Kosice - Moscow -Prague Collaboration

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AN INVESTIGATION OF ASSOCIATIVE MULTIPLICITY IN THE REACTION pp → p (SLOW) + X AT 22.4 GeV/c

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Исследование ассоциативной множественности в реакции <u>pp</u> — р (медл.)+ X при 22,4 ГэВ/с

В работе исследуется поведение ассоциативной множественности как функции недостающей массы к идентифицированному протону в реакции $\overline{pp} \longrightarrow p(\text{медл.}) + X$ при 22,4 ГэВ/с. Показан различный характер поведения ассоциативной множественности для дифракционных ($M_{\chi}^2/s \le 0.1$) и иедифракционных событий. В рамках двухкомпонентиой модели качественно объяснено поведение отношения $< n(M_{\chi}^2) > /D$ как функции M_{χ}^2 . При использовании дастабной переменной

 $z' = (n_{ch} - 1 - \alpha)/(< n(M_{\chi}^2) > -\alpha), \quad \alpha = -1.04$ для всех значений M_{χ}^2 был получен аналог скейлинга KNO в системе X. Отмечается подобие реакций $\overline{pp} \longrightarrow p(мед.) + X$ и $pp \longrightarrow p(мед.) + X$.

Работа выполнена в Лаборатории высоких энергий ОИЯИ.

Препринт Объединенного института ядерных исследований. Дубна 1979

Batyunya B.V. et al.

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An Investigation of Associative Multiplicity in the Reaction $\overline{pp} \longrightarrow p(slow) + X$ at 22.4 GeV/c

The associative multiplicity as a function of the missing mass squared to the identified proton in the reaction $\bar{p}p \rightarrow p(slow) + X$ at 22.4 GeV/c is studied. A different behaviour of the associative multiplicity for diffraction $(M_{\chi}^2/s \le 0.1)$ and nondiffraction events is observed. The M_{χ}^2 dependence of the ratio $\langle n(M_{\chi}^2) \rangle / D$ is qualitatively explained on the basis of the two-component model.

An analogue of KNO scaling in the system X is obtained for all values of M_X^2 with the help of the scaling variable.

 $z' = (n_{xh} - 1 - a)/(\langle n(M_{\chi}^2) \rangle - a), \quad a = -1.04.$

A similar behaviour of the reactions $\overline{p}p \longrightarrow p(slow) + X$ and $pp \longrightarrow p(slow) + X$ is pointed out,

The investigation has been performed at the Laboratory of High Energies, JINR.

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1. The associative characteristics of recoil system X to identified particle c in the reaction aN - c + X have been extensively studied for various particles a and $c^{/1/}$.

In this paper we study the reaction

$$\overline{p}p \longrightarrow p(slow) + X \tag{1}$$

at 22.4 GeV/c. About 21000 inelastic events were used in this analysis. Protons with a laboratory momentum of ≤ 1.5 GeV/c were identified by ionization thus yielding about 5600 events of the type (1). The questions concerning the separation of elastic events, the corrections for losses of inelastic 2-prong events with slow recoil protons, etc., have been analyzed in ref. $^{/2}/$.

2. The average charged multiplicity of system X is defined through the topological (inclusive) differential cross sections $d\sigma_{x}/dM_{y}^{2}(d\sigma/dM_{y}^{2})$ of reaction (1),

$$\langle n(M_{\chi}^2) \rangle = \sum_{n} (n-1) d\sigma_n / dM_{\chi}^2 / d\sigma / dM_{\chi}^2$$
, (2)

as a function of the missing mass squared M_{χ}^2 to the identified proton. In *fig.* 1 we compare the M_{χ}^2/s dependence of $\langle n(M_{\chi}^2) \rangle$ both for the reaction (1) and the reaction $pp \longrightarrow p(slow) + X$ at 19 $GeV/c^{/9/}$. We have fitted our data points by the logarithmic expression

$$< n(M_{\chi}^2) > = a + b \ell n(M_{\chi}^2 / M_0^2), \qquad M_0 = 1 \text{ GeV}^2.$$
 (3)

in an interval of $4 \le M_X^2 \le 22$ GeV² and by the power law

$$< n(M_{\chi}^2) > = a_1 + b_1 \sqrt{M_{\chi}^2}$$
 (4)

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Fig. 1. The average charged multiplicity $\langle n(M^2) \rangle$ of system X as a function of M^2_{χ}/s for the reactions $X \ \overline{p} p \longrightarrow p(slow) + X$ (22.4 GeV/c) and $pp \longrightarrow p(slow) + X$ (19 GeV/c)⁹⁹.

in the diffraction region $(M_X^2 < 4 \text{ GeV}^2)$, as suggested by multiperipheral models (see, e.g., ref. $^{/3/}$) and by the fragmentation model "NOVA", incorporating cluster production $^{/4/}$, respectively. The fitted parameters are shown in *table 1* together with similar parameters for the reaction pp \longrightarrow p (slow) + X.

3. In fig. 2 we show the M_{χ}^2 -dependence of the ratio $< n(M_{\chi}^2) > / D$, where

$$D = (\langle n^2(M_{\chi}^2) \rangle - \langle n(M_{\chi}^2) \rangle^2)^{1/2}.$$
 (5)

A clear dip structure seen for $2 \le M_{\chi}^2 \le 4 \text{ GeV}^2$ can be explained, e.g., due to overlapping of two different mechanisms

Table 1

The parameters in eqs. (3), (4) and corresponding

Reaction momentum (GeV/c)	8	Ъ	x ² /ND	⁸ 1	^b l	₹ ² /ND
pp (22.4) pp (36)	0.69±0.11 0.78±0.45	0.94±0.05 0.85±0.16	7.8/7 1.4/5	0.63±0.09 0.66±0.18	0.61±0.06 0.55±0.13	5/2 0.3/I
pp (102- 405)	-0.55±0.39	1.37±0.09	-	0.7I±0.I7	0.72±0.07	-

 $X^2/ND \cdot pp(102-405)^{/8/}$, $pp(36)^{/10/}$.

in the diffraction region. In such a two-component model $^{/5/}$ the dispersion can be written in the form

$$\mathbf{D} = (\beta_1 \mathbf{D}_1^2 + \beta_2 \mathbf{D}_2^2 + \beta_1 \beta_1 (< \mathbf{n}_1 (\mathbf{M}_{\chi}^2)) - < \mathbf{n}_2 (\mathbf{M}_{\chi}^2) >)^2)^{1/2},$$
(6)



where the probabilities β_i of the two processes satisfy the normalization condition $\beta_1 + \beta_2 = 1$. The diffraction mechanism $(D=D_1)$ dominates at $M_{\chi}^2 \sim 1 \ GeV^2$ and the nondiffraction mechanism is responsible for $M_{\chi}^2 > 4 \ GeV^2$ $(D=D_2)$, while in the intermediate region the "interference" term in (6) increases the dispersion thus yielding the observed dip structure.

4. We analyse also the possibility of an analogue of KNO-scaling for the system $X^{/6/}$. In *table 2* we show the normalized moments $c_q = <n^q (M_X^2) > / < n (M_X^2) >^q$. It is seen that c_q values are different in the diffraction $(M_X^2 / s \le 0.09)$ and nondiffraction regions. From this it follows that the usual KNO-scaling for the system X is not valid, i.e., the function

$$\Psi(\mathbf{z}, \mathbf{M}_{\chi}^{2}) = \frac{\langle \mathbf{n} (\mathbf{M}_{\chi}^{2}) \rangle d\sigma_{\mathrm{N}} / d\mathbf{M}_{\chi}^{2}}{d\sigma / d\mathbf{M}_{\chi}^{2}}$$
(7)

Table 2

The normalized associative moments for various intervals of $M_v^2/s \cdot pp(205)^{/6/}$

¥2/•	°2	°3	C ₄	Reaction momentum (GeV/c)
0	I.4I ± 0.02	2.71 ± 0.16	6.55 ± 0.85	pp (22.4)
0.05	1.40 ± 0.04	2.58 ± 0.20	5.51 ± 0.82	pp (36)
	1.41 ± 0.03	2.54 ± 0.17	5.54 ± 0.69	pp (205)
	1.26 ± 0.02	I.85 ± 0.09	3.18 ± 0.44	jp (22.4)
0.09 -	1.25 + 0.05	I.60 ± 0.19	2.90 ± 0.60	TTD (36)
0.18	1.22 ± 0.02	1.71 ± 0.06	2.66 ± 0.18	pp (205)
	1.23 ± 0.02	1.72 ± 0.08	2.71 ± 0.18	p (22.4)
0.18 -	I.23 ± 0.04	I.7I ± 0.15	2.63 ± 0.43	pp (36)
0.30	1.23 ± 0.02	1.77 ± 0.07	2.87 ± 0.20	pp (205)
	T 23 ± 0 02	T 75 ± 0.04	2 76 ± 0 72	
0.20 -	T 27 ± 0.05	T 96 ± 0 17	2.10 ± 0.10	m (96)
0.39	1.27 - 0.00	1.00 - 0.17	3.06 - 0.00	pp (305)
	1.23 - 0.02	1.73 - 0.08	2.70 - 0.16	pp (200)
0.39 -	I.23 ± 0.02	1.74 ± 0.04	2.74 ± 0.13	p (22.4)
0,50	1.24 ± 0.04	1.76 ± 0.15	2.73 ± 0.47	pp (36)

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essentially depends on M_{χ}^2 . To get rid of the M_{χ}^2 dependence, we introduce the scaling variable $^{/7/}$

$$z' = (n_{ch} - 1 - \alpha)/(\langle n(M_{\chi}^2) \rangle - \alpha)$$
 (8)

The parameter a = -1.04 in (8) is determined from the condition

$$[(-\alpha)/D]_{d} = [(-\alpha)/D]_{nd}.$$
 (9)

where the l.h.s. (r.h.s.) have been calculated for $0 \le M_{\chi}^2 \le 4 \text{ GeV}^2$ $(4 \le M_{\chi}^2 \le 22 \text{ GeV}^2)$. The corresponding normalized distribution $\Psi(z')$ is shown in *fig. 3* for $0 \le M_{\chi}^2 \le 22 \text{ GeV}^2$. The solid curve is the result of the fit:

$$\Psi(z') = (0.05\pm0.01) \exp[(7.78\pm0.59)z' - (4.17\pm0.58) z'^{2} + (0.13\pm0.18) z'^{3}](10)$$

with $X^2 / ND = 26/20$.

5. CONCLUSIONS

The following results have been obtained in a study of the associative multiplicities in the reaction (1).

(i) We observe striking similarity between the reactions $\tilde{p}p \longrightarrow p(slow) + X$ and $pp \longrightarrow p(slow) + X$ for comparable energies in an interval of $0 < M_X^2/s < 0.5$. This fact is illustrated by the $< n(M_X^2) >$ distribution in *fig. 1* and the corresponding parameters

in *table 1.* Besides, in the diffraction region $M_{\chi}^2/s \lesssim 0.09$ the $< n(M_{\chi}^2) >$ distribution is clearly independent of incident momentum (see the parameters a_1 , b_1 in *table 1*).

tum (see the parameters a_1 , b_1 in *table 1*). (ii) The M_X^2 dependence of the ratio $\langle n(M_X^2) \rangle / D$ is qualitatively explained on the basis of the two-component model.

(iii) An analogue of KNO-scaling in the system X is obtained for all values of M_X^2 with the help of the scaling variable

$$z' = (n_{ch} - 1 - a)/(\langle n(M_X^2) - a), \quad a = -1.04.$$



in an interval of 0 < M_χ^2 < 22 GeV 2 . The solid curve is defined in eq. (10).

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