

O<sup>+</sup>- STATES IN <sup>102</sup>Pd AND <sup>108</sup>Cd.

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Abstract

A presentation is made of an angular correlation experiment by which we could identify the two fonon O<sup>+</sup>- states in <sup>102</sup>Pd and <sup>108</sup>Cd. In the case of Pd a comparison is shown with IBA calculations and a level structure calculated with a asymmetric rotor vibrator model.

1. Introduction

The even-even Pd and Cd-nuclei have for many years been interpreted as having an anharmonic vibrational structure. A two fonon multiplet can be seen from N=60 for Cd-nuclei and from N=58 for Pd-nuclei (fig. 3.4.). In order to extend the systematics in the behaviour on those multiplet states we looked for the O<sup>+</sup>- states in <sup>102</sup>Pd and <sup>108</sup>Cd.

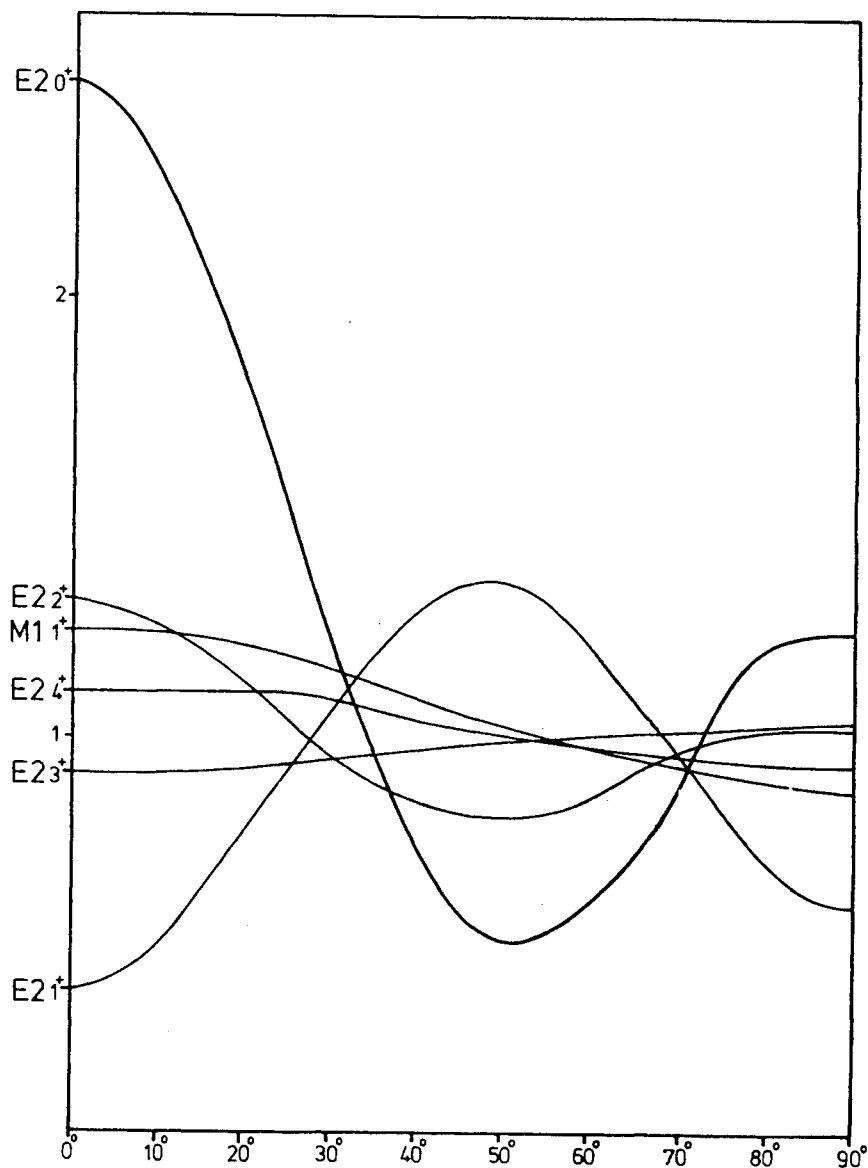


fig.1. Angular correlation pattern of some x<sup>π</sup>→2<sup>+</sup>→0<sup>+</sup> cascades.

Therefore we performed a  $\gamma$ - $\gamma$  angular correlation experiment on the  $(\beta^+ + EC)$ -decays of  $^{102m}\text{Ag}$ - and  $^{108m}\text{In}$  nuclei. In  $^{102}\text{Pd}$  we found a  $0^+$  state very close to the long living anomalous  $0^+$  on which we reported earlier <sup>1)</sup>. The electromagnetic features of the new  $0^+$ -state are more in agreement with a two phonon description whereas the other  $0^+$ -state seems to be some kind of intruder which can not be accounted for in a normal collective (IBA, triaxial rotor vibrator etc.) model. The  $0^+$ -state in  $^{108}\text{Cd}$  is lower in energy than the expected trend, but seems to have normal electromagnetic properties.

## 2. The angular correlation experiment.

The produced radio activity and the available measuring time for a ISOL setup at a cyclotron are normally not sufficient to perform a decent  $\gamma$ - $\gamma$  angular correlation experiment. However, the huge anisotropy of a  $0^+ \rightarrow 2^+ \rightarrow 0^+$  cascade makes it possible to distinguish it from other  $x^\pi \rightarrow 2^+ \rightarrow 0^+$  cascades, even with poor statistics as can be seen in fig.1. This means that the angular correlation technique can still be a useful tool when looking for  $0^+$ -states in low activity- and short living sources.

In our experiment we only measured the correlation at the three crucial angles  $0^\circ$  ( $180^\circ$ ),  $50^\circ$  and  $90^\circ$ . The  $^{102m}\text{Ag}$  and  $^{108m}\text{In}$  sources were produced with a  $\text{nat}_{\text{Mo}}(^{14}\text{N}(90\text{MeV}), \text{xpy})$  reaction and mass-separated using the LISOL facility at Louvain-la-Neuve <sup>2)</sup>. An example of the resulting spectra gated on the  $2^+ \rightarrow 0^+_g$  transition is shown in fig.2. The spectrum at  $180^\circ$  shows the  $(1101.7 \pm 0.5)\text{keV}$  gamma which proved to have the good anisotropy for a  $0^+ \rightarrow 2^+ \rightarrow 0^+$  cascade. From the coincidences we could also deduce that the  $0^+$  level at 1658 keV is almost completely fed by a 729 keV gamma coming from a known level at 2390 keV ( $J^\pi 1, 2^+$ )

In the same way we observed a  $0^+$  state at  $(1375.0 \pm 0.9)\text{keV}$  which is mainly fed by 1244, 1306 and 827-keV transitions depopulating known levels in  $^{108}\text{Cd}$  <sup>3)</sup>. Also mini-orange spectra were taken from both sources in order to measure the  $E0(0^+ \rightarrow 0^+_g)$  transition strength. From these

spectra only an upper limit could be deduced for the parameter.

$$X = \frac{B(E0, 0^+ \rightarrow 0^+_g)}{B(E2, 0^+ \rightarrow 2^+_1)}$$

For  $^{102}\text{Pd}$  we have  $X \leq 2.7 \cdot 10^{-1}$  and for  $^{108}\text{Cd}$   $X \leq 4$ .

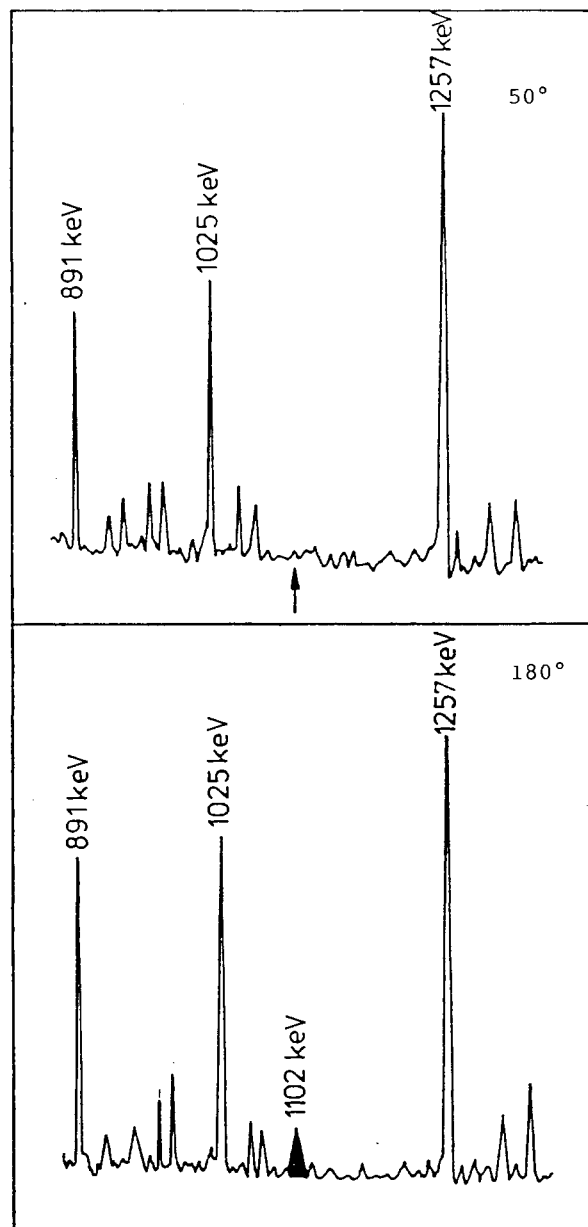


fig. 2. Spectra at  $50^\circ$  and  $180^\circ$  gated on the  $2^+ \rightarrow 0^+_g$  transition

### 3. Discussion.

In fig. 3 we can see that the  $0^+$ - state in  $^{102}\text{Pd}$  is following the systematical trend of rising when going to lesser neutrons. The  $2^+$  and  $4^+$  do not go up as fast and in  $^{102}\text{Pd}$  the two fonon triplet is completely broken. The same systematical behaviour can be reproduced by IBA-2 calculations <sup>4)</sup> as shown in fig. 4.

In table 1 a comparison is shown of the low energy levels in  $^{102}\text{Pd}$  and a theoretical calculation with an anharmonic vibrator model <sup>5)</sup>.

Table 1.

$I^\pi$	Exp. (keV)	Theory: $E_{\text{INn}}$	INn
$0^+$	0	0	011
$2^+$	556	552	211
$4^+$	1275	1333	411
$2^+$	1534	1492	211
$0^+$	<u>1658</u>	<u>1658</u>	012
$6^+$	2111	2119	611
$4^+$	2138	2177	421

In this calculation the experimental energy of the  $0^+$ - state was not used in fitting the model parameters, however, the agreement with experiment is excellent.

The first excited  $0^+$ -state in  $^{108}\text{Cd}$  (fig. 5) does not seem to follow this systematical trend of rising when going more neutrondeficient. In this case it would be interesting to find the second excited  $0^+$ -state which appears to be low in the heavier Cd isotopes. These  $0^+$ - states can probably be explained in terms of 2-proton particle - 2 hole excitations through the  $Z=50$  shell closure as in the case of the even Sn isotopes <sup>6)</sup>.

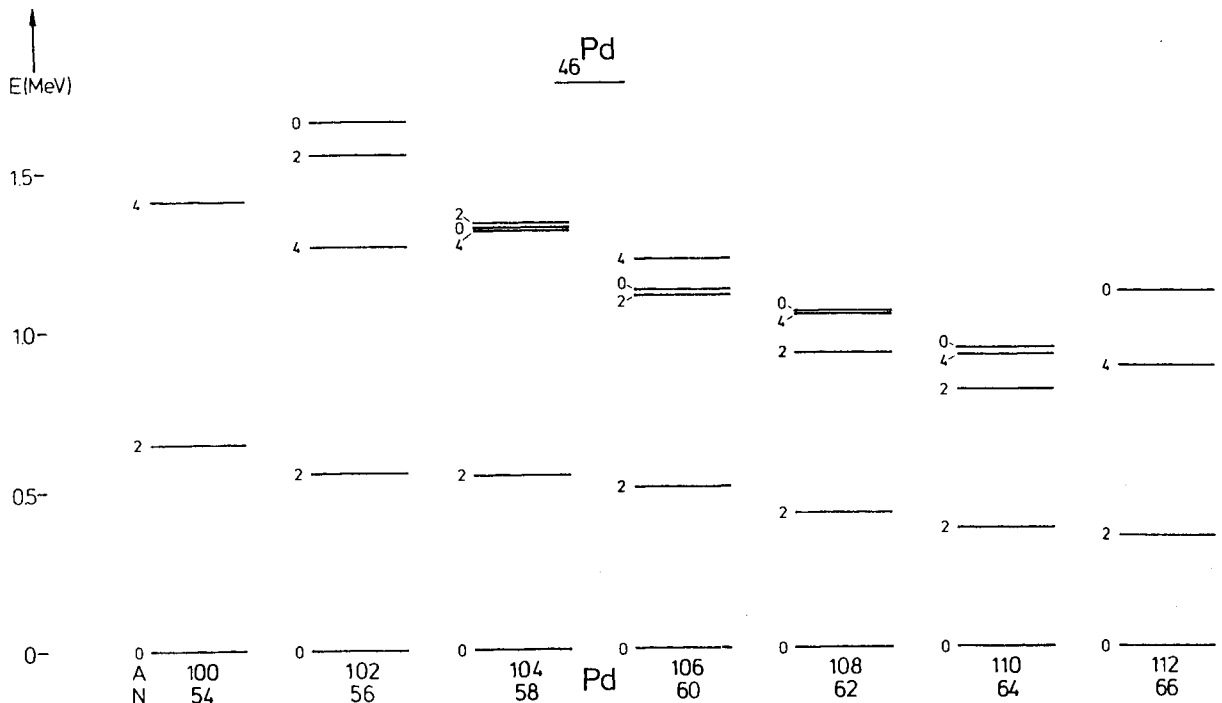


fig. 3 Positive parity states in even Pd nuclei.

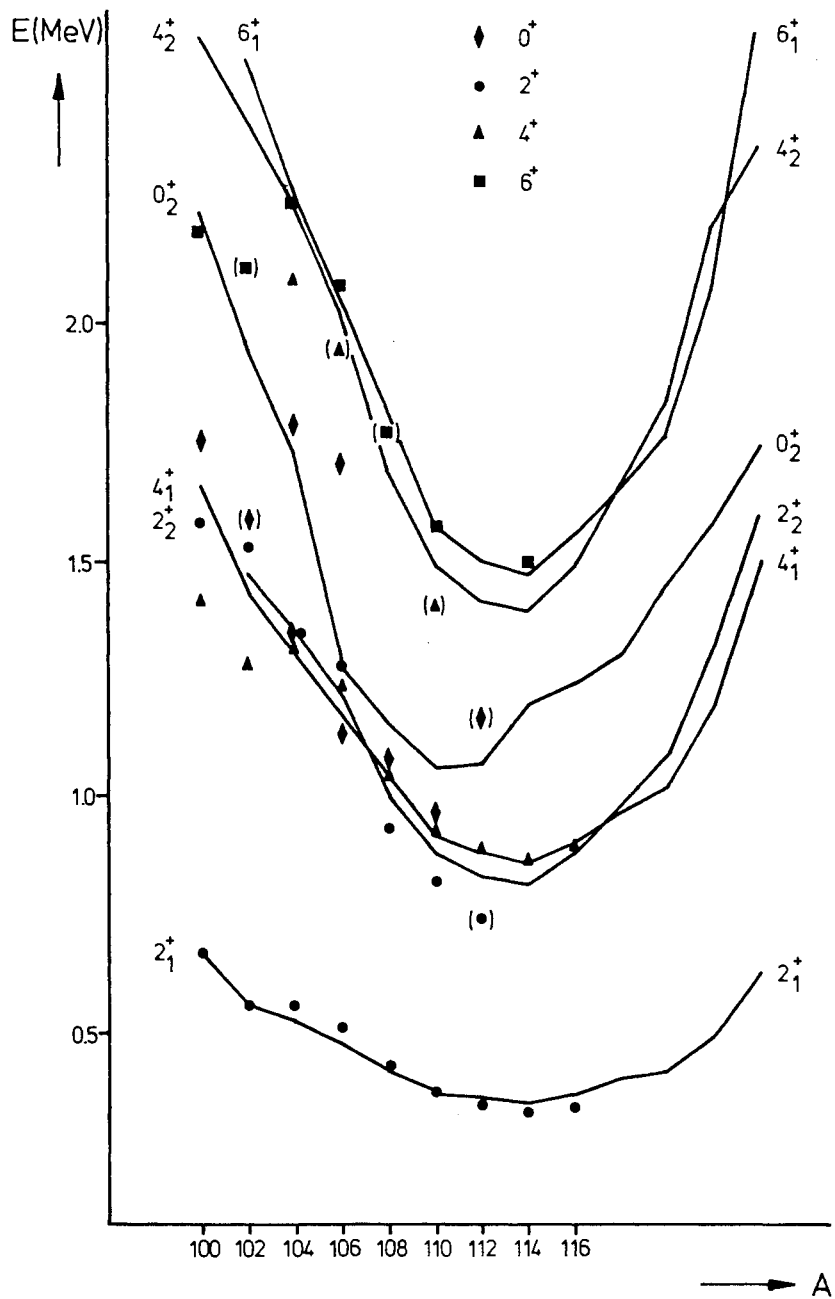


fig. 4 IBA-2 calculation of the Pd isotopes.

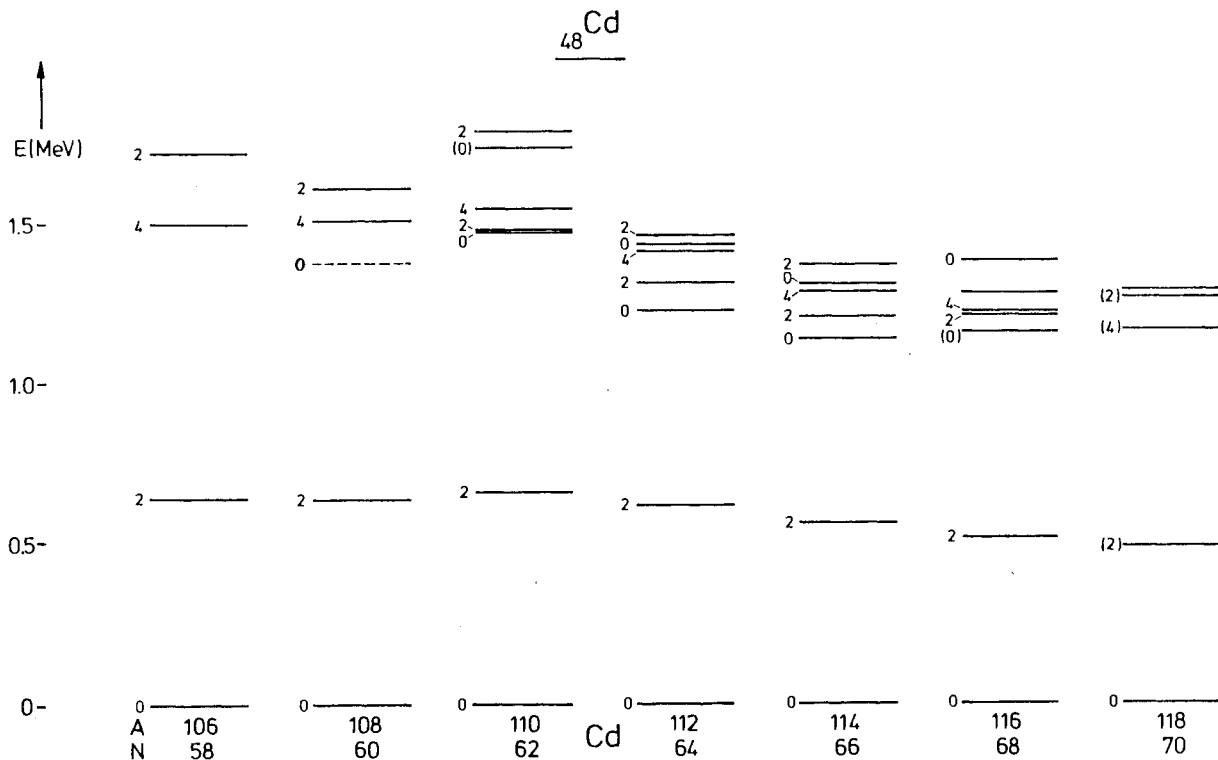


fig. 5 Positive parity states in even Cd nuclei .

### References

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### DISCUSSION

*J. Kantele:* I have some complementary information to give: within the Jyväskylä-Uppsala collaboration, we have studied  $^{102}\text{Pd}$  low-spin states using  $(pp')$  and double Coulomb excitation. The 1648 keV state is seen in the latter reaction which indicates a collective character for this state; it probably contains some two-phonon strength. The 1592 keV level does not seem to be connected to the  $2_1^+$  state, which yields the largest known X-value, over 400. Thus the 1592 keV  $0_2^+$  state has a structure completely different from that of the  $0_3^+$  one at 1648 keV; the wave function is probably dominated by some two-quasiparticle component.

*V. Paar:* How many parameters do you have in fitting spectra in your figure? Do you fit to each nucleus parameters separately?

*K. Cornelis:* The parameters in this fit were allowed to vary in the fit.

*J. Stachel:* In principle the full IBA-2 Hamiltonian has  $\approx 12$  free parameters. The calculations referred to here (P. Van Isacker, G. Puddu, Nucl. Phys. 1980) are a simultaneous fit to all neutron-rich ( $N = 50-66$ ) Pd- and Ru-isotopes. Some of the parameters were set = 0 in this fit, all parameters involving only protons were kept constant over the whole isotope chain, some were taken equal for Ru- and Pd-isotopes. So it is essentially the change in three parameters describing the change in excitation energies and decay properties along each isotope chain.

*D.S. Bremner:* In your IBA-2 calculations which parameters did you vary to obtain your fits to the Pd and Cd isotopes? What criteria were used to determine the "best" fit to these data?

*K. Cornelis:* For this question I refer to P. Van Isacker, for he has been doing the calculations. He certainly can tell you more about the details.