

FUSION BARRIERS WITH FOLDING POTENTIALS
FOR SYSTEMS WITH $1000 < Z_1 Z_2 < 3000$

by

M. LOZANO, A. MANDLY, and G. MADURGA

Departamento de Física Atómica y Nuclear

Facultad de Física, Universidad de Sevilla, Sevilla 4, Spain.

Following a suggestion repeatedly stated by Vaz, Alexander, and Satchler in their recent study of fusion barriers (1), we have used two different folding models for the construction of the real nuclear potential for very heavy systems in the region $1000 < Z_1 Z_2 < 3000$, and computed the corresponding s-wave barrier. A systematic description of the nuclear densities has been undertaken as a starting point for the foldings.

THE NUCLEAR DENSITY DISTRIBUTIONS have been computed as the sum of squared single-particle wave functions, ψ_i , weighted with occupation numbers. The ψ_i are initially eigenfunctions of a Woods-Saxon potential with a spin-orbit term of the Thomas type and a Coulomb repulsion for protons, later corrected for nonlocality and orthogonalized. The method has been described at length by Malacuti et al (2) and specified for our use in our previous work (3), hereafter referred to as VLM. With respect to the latter we have changed here the constant term of the depth to 50 MeV, instead of 55.7, with two exceptions: 46.7 for ^{152}Sm and 51.5 for ^{209}Bi . Table I shows the results for 17 nuclides used in this work, compared with the experimental rms radii of the charge densities.

FOLDING POTENTIALS of two different types are used here. The SL-folding, described and widely applied by Satchler and Love (4), is computed here with the proton and neutron densities obtained by the method outlined above and with the M3Y effective nucleon-nucleon interaction. The VLM-folding is based on a Gaussian density-dependent effective interaction described in ref. (3),

TABLE I Rms radii (in fm) for the nuclear density distributions computed in this work, and experimental rms charge radius.

Nucleus	p	n	p+n	charge	experimental
³⁵ Cl	3.275	3.213	3.243	3.372	3.351
⁴⁰ Ar	3.326	3.507	3.427	3.419	3.38 - 3.56
⁵⁶ Fe	3.673	3.717	3.697	3.759	3.727 - 3.800
⁶⁵ Cu	3.806	3.927	3.874	3.887	3.947(22)
⁸⁴ Kr	4.068	4.265	4.182	4.143	
⁸⁶ Kr	4.068	4.323	4.215	4.142	
⁹⁰ Zr	4.197	4.300	4.255	4.271	
¹⁰⁹ Ag	4.457	4.612	4.546	4.525	
¹³² Xe	4.732	4.944	4.858	4.795	
¹³⁹ La	4.796	5.017	4.928	4.859	4.85
¹⁴¹ Pr	4.842	5.010	4.940	4.905	
¹⁴⁴ Sm	4.910	5.003	4.963	4.972	
¹⁴⁸ Sm	4.920	5.103	5.027	4.982	4.952 - 5.026
¹⁵⁴ Sm	4.998	5.345	5.208	5.058	5.126
¹⁶⁵ Ho	5.139	5.311	5.242	5.197	5.19 - 5.23
²⁰⁹ Bi	5.503	5.696	5.620	5.556	5.52
²³⁸ U	5.793	6.070	5.965	5.843	5.843

and uses our proton + neutron densities. In both cases only a few folding points in the relevant region of the potential have been needed to fit an analytical approximation, as explained in VLM. This has the form

$$V(r) = \begin{cases} -v_0 (an)^n e^{-n} & \text{if } s < an \\ -v_0 s^n \exp(-s/a) & \text{if } s > an \end{cases}$$

where $s = r - C_1 - C_2$, C_i are the standard half density radii, and the parameters v_0 , n , a are fitted to seven points in the region $0.8 < s < 3.6$ fm. In Table II are listed the values of these parameters for SL-folding (above) and VLM-folding (below) corresponding to the 14 systems considered in this work as well as the range and depth of the (irrelevant) flat constant region introduced in this analytical expression. In the range of distances which are significant for fusion and for elastic scattering VLM-folding potential is almost twice as deep as SL-folding.

FUSION BARRIERS obtained with these potentials and their location are also listed in Table II. VLM-folding predicts lower barriers by a few percent.

TABLE II Parameters that fit SL-folding (1st line) and VLM-folding (2nd line), range and depth of flat bottom, and location and height of fusion barrier (Units: MeV and fm).

Reaction	$Z_1 Z_2$	V_0	n	a	$C_1 + C_2 + an$	$-V(0)$	R_B	V_B
$^{35}\text{Cl} + ^{141}\text{Pr}$	1003	145.4	0.4479	0.6569	9.607	53.71	11.39	118.3
		360.5	0.9298	0.5591	9.833	77.42	12.05	113.0
$^{86}\text{Kr} + ^{65}\text{Cu}$	1044	150.8	0.3868	0.6711	9.504	60.78	11.33	123.7
		338.9	0.7619	0.5873	9.692	85.73	11.93	118.6
$^{40}\text{Ar} + ^{144}\text{Sm}$	1116	153.8	0.4389	0.6666	9.837	57.82	11.64	128.9
		347.7	0.7564	0.5999	9.998	89.77	12.29	123.1
$^{40}\text{Ar} + ^{148}\text{Sm}$	1116	148.8	0.3805	0.6986	9.872	61.42	11.75	125.5
		324.9	0.6347	0.6452	10.02	98.04	12.44	121.4
$^{40}\text{Ar} + ^{154}\text{Sm}$	1116	145.5	0.2809	0.7817	9.916	71.79	12.00	124.3
		299.5	0.4400	0.7569	10.03	118.9	12.78	117.4
$^{40}\text{Ar} + ^{165}\text{Ho}$	1206	153.1	0.3757	0.7025	10.12	63.74	12.00	135.1
		341.6	0.6253	0.6477	10.26	103.9	12.70	128.6
$^{56}\text{Fe} + ^{132}\text{Xe}$	1404	158.1	0.3952	0.6684	10.12	62.93	11.80	160.1
		351.9	0.7464	0.5965	10.30	91.20	12.44	153.1
$^{86}\text{Kr} + ^{90}\text{Zr}$	1440	158.3	0.3922	0.6634	10.07	63.93	11.72	165.4
		368.8	0.7854	0.5682	10.26	89.24	12.30	159.1
$^{40}\text{Ar} + ^{238}\text{U}$	1656	150.1	0.3183	0.7593	11.04	69.48	12.87	172.9
		333.9	0.4858	0.7137	11.14	122.8	13.63	164.5
$^{86}\text{Kr} + ^{109}\text{Ag}$	1692	158.1	0.3254	0.6917	10.41	70.26	12.04	189.1
		358.0	0.6768	0.6024	10.59	99.14	12.65	181.6
$^{56}\text{Fe} + ^{165}\text{Ho}$	1742	152.0	0.2907	0.7104	10.57	71.84	12.21	192.0
		341.1	0.5552	0.6448	10.72	110.7	12.89	183.3
$^{86}\text{Kr} + ^{139}\text{La}$	2052	158.4	0.2480	0.7101	10.87	80.33	12.47	221.9
		377.9	0.6801	0.5971	11.10	103.7	13.09	213.3
$^{84}\text{Kr} + ^{165}\text{Ho}$	2412	160.3	0.2977	0.7243	11.26	75.36	12.77	254.6
		368.3	0.6208	0.6321	11.44	110.8	13.43	244.1
$^{84}\text{Kr} + ^{209}\text{Bi}$	2988	167.4	0.3214	0.6978	11.86	75.05	13.25	305.4
		473.4	0.9542	0.5465	12.16	97.95	13.93	293.2

not only as compared with SL-folding but also with the "empirical" values reported in ref. (1) for $Z_1 Z_2 < 1400$. However Vaz et al note that their technique for defining an empirical fusion barrier could possibly overestimate the barriers by 3-8% for $Z_1 Z_2 > 800$.

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