**INTERNATIONAL MEETING ON RESEARCH AND**

**TEST REACTOR CORE CONVERSIONS FROM**

**HEU TO LEU FUELS**

**Argonne National Laboratory**

**November 8-10, 1982**

**NEUTRONICS ANALYSIS OF THE PROPOSED 25-MW LEU TRIGA MULTIPURPOSE RESEARCH REACTOR\***

**Martlas Nurdln**

**Research Centre for Nuclear Techniques National Atomic Energy Agency Republic of Indonesia**

**and**

**M. M. Bretscher and J. L. Snelgrove**

**Argonne National Laboratory United States of America**

The submitted menuscript has been authored by a contractor of the U.S. Government<br>under contract No. W-31-109-ENG-38. Accordingly, the U.S. Government retains a  $\alpha$  consequency, the only-free licenses to publish the space of the published form of this<br>intribution, ceralism others to do so, for œ U.S. Government purpoess.

**\*Work performed under the auspices of the U.S. Department of Energy and the International Atomic Energy Agency.**

This report was prepared as an account of work sponsored by an agency of the United States<br>Government Neither the Huised States Refer-Neither the United States Government nor any agency in the United States<br>also any warranty express Government nor any agency thereof, nor any of their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy contract of the school of the schoo apparatus, product, or mendation, or favoring by the United States Covernment or unply us encorsement, recom-<br>and opinions of authors expressed herein do not necessarily state or reflect those of the<br>United States Government of the present of th ence herein to any specific commercial product, process, or service by trade name, trademark, Refermant and the metal of the service of the service of the service of the manufacturer. or otherwise does not all the service manufacturer, or otherwise does not necessarily constitute or imply its endorsement, recom-<br>mendation, or favoring by the Heistel exercise of imply its endorsement, recomprocess disclosed, or represents that its use would not infringe privately owned rights.<br>ence herein to any specific containments are would not infringe privately owned rights. bility for the accuracy, completeness, or usefulness of any information, United States Government or any agency thereof. Government.

CONF-821155--5

DE83 007726

MASTER

**DISCLAIMER** 

 $CONF - 821155 - -5$ 

### 1. INTRODUCTION

More than two years ago the government of Indonesia announced plans to purchase a research reactor for the Fuspiptek Research Center in Serpong, Indonesia to be used for isotope production, materials testing, neutron physics measurements, and reactor operator training. Reactess using lowenriched uranium (LEU) plate-type and rod-type fuel elements were considered. This paper deals with the neutronic evaluation of the rod-type 25-MW LEU TRIGA Multipurpose Research Reactor (MPRR) proposed by the General Atomic Company of the United States of America.<sup>1</sup>

### 2. indian FOR AND FUEL ELEMENT DESCRIPTIONS

The 25-MW TRIGA Multipurpose Research Reactor is of the swimming pool type and is fueled by U-ZrH-Er rods arranged In 36-rod clusters or fuel elements. The core, which has an active height of 22 in., is water-moderated and berylling reflected.

An 11 x 12 aluminum grid plate is used to position the nominal core consisting of 40 fuel elements, a central 14.7-cm-square cavity and six 7.34-cmsquare in-core irradiation positions. The active core is surrounded by 50 beryllium reflector blocks, half of which contain central irradiation holes. Six natural boron carbide rods are used to control the reactor. Figure 1 shows the arrangement of the fuel, control, irradiation, and beryllium reflector elements. Not shown In the figure are three eight-inch-diameter radial beam tubes and one through tube tangent to the core at the lower flat face. Normally the in-core irradiation spaces would contain experiments or dummy experiments to reduce, for safety reasons, power peaking in adjacent fuel rods. For calculational purposes, however, **these** In-core irradiation positions were assumed to be water-filled.

The fuel elements consist of 36 fuel rods arranged in a  $6 \times 6$  square array within a square aluminum shroud 7.34 cm on a side. Table 1 describes the fuel pins, which are clad in Incoloy 800, and the 36-rod fuel cluster.

### 3. CALCULATIONAL METHODS

The calculational methods used in this study are those which have been described in detail in Appendix A of Ref. 2. A brief description of these basic methods is given below.

The EPRI-CELL code<sup>3</sup> was used to generate broad group, burnup-dependent cross sections and atom densities for subsequent diffusion-theory and transporttheory calculations. EPRI-CELL combines a GAM-1\* resonance treatment in the epithermal energy range with a THERMOS<sup>5</sup> heterogeneous, integral-transport treatment in the thermal energy range. As input, EPRI-CELL utilizes a 68-group epithermal GAM library updated with ENDF/B-IV data processed using the integraltransport option of the  $MC^{2}-2^{6}$  code to account for resonance self-shielding and a 35-group thermal THERMOS library updated with ENDF/B-IV data processed using either the XLACS1 module of **the AMPX** system<sup>7</sup> or **the** NJOY code.<sup>8</sup> Spatial self-shielding factors, calculated using  $MC<sup>2</sup>-2$ , were also used as input to EPRI-CELL.

**The MC^-2 code has a more rigorous resonance treatment than does the EPRI-CELL code. Therefore, energy self-shielding factors were used for the 238(j resonance capture region and the <sup>235</sup> U resonance absorption region.** Use of these energy self-sheilding factors effectively replaces the <sup>238</sup>U and **235u resonance cross sections generated by EPRI-CELL with those calculated by the MC<sup>2</sup>-2 code.**

An il-group cross section set consisting of four fast, three epithermal, **and four thermal groups was used for most of the calculations in this study. Some calculations were also perforaed using the standard five-group structure commonly used at Argonne National Laboratory for MTR plate-type reactor studies. It was found, however, that the five group set did not adequately account for thermal neutron upscattering from excited ZrH energy states in TRIGA fuel. The broad group energy boundaries for both group structures are listed in Table 2.**

**Different cell models were needed to generate appropriate EPRI-CELL cross sections for the various reactor regions. Once generated, these cross sections sets were combined into one master set and used for multigroup diffusion and transport calculations. Burnup-dependent cross sections were calculated only for isotopes in the fuel pin. Separate unit cell calculations were made for the fuel rod, beryllium reflector, water radial reflector, water axial reflector, internal water-filled flux traps, control rod, and control rod follower.**

**Most of the results of this study are based on XY raultigroup diffusion calculations. Energy-independent axial extrapolation distances derived from flux profiles calculated in an RZ diffusion-theory model were used to account for the axial leakage and power profile. The extrapolation distances used are given in Table 3. Studies have shown that axial buckling values are, to all practical purposes, independent of core temperatures. This is because temperature effects the thermal portion of the neutron spectrum whereas leakage is due mostly to high energy neutrons.**

**The REBUS-2 fuel cycle analysis code^ was used to perform burnup calculations. The XY model shown in Fig. 1 was used for the diffusion—theory calculations in REBUS-2. The water thickness outside the beryllium reflector was taken to be 14.68 cm. Except for the outside water pool, each grid position was represented by a uniform 5x 5 mesh. The use of more mesh intervals was studied briefly. The chief effect was an increase in the eigenvalue by ~0.7Z Ak for a doubling of the number of mesh intervals used. The mesh spacing had little effect on calculated flux profiles or control rod worths. Therefore, even though**  $k_{\text{eff}}$  **may have been underestimated by up to 1%**  $\Delta k$ **, the 5**  $\times$  **5 mesh spacing was used to conserve computer resources. Both non-equilibrium and equilibrium fuel cycle calculations were performed. In the non-equilibrium calculations a core of 40 fresh fuel elements was allowed to burn down with no fuel replacement. In the equilibrium calculations a fixed number of fuel elements were replaced at the end of each cycle and the remaining fuel elements were moved to new locations in the core. After some preliminary studies, a five-path fuel management scheme as given in Table 4, was selected. Cross sections representative of the middle-of-cycle burnup were used in the REBUS-2 calculations. The fuel temperature was assumed to be 800K.**

**Because of the very strong absorbing quality of the B4C control rods, the conditions for the valid application of diffusion theory are severely violated, and, therefore, diffusion theory cannot accurately predict controlrod worths. However, approximate control-rod worths can be obtained using diffusion theory with suitable internal boundary conditions calculated from transport theory. The internal group-dependent boundary condition is just the ratio of the neutron current to flux at the surface of the control rod cell. The internal boundary conditions were calcculated for a cell consisting of a control rod surrounded by homogenized fresh fuel using the one-dimensional transport-theory code ONEDANT.l" To account for flux and scattering anisotropics, the calculations were performed in the P1S4 approximation. For the case of a fully-withdrawn rod, the aluminum follower and water were homogenized and normal diffusion theory was used. In order to validate these control-rod-worth calculations, the control rod cell, both with a rod inserted and a rod withdrawn, were calculated using ONEDANT, the VIM Monte Carlo code,11 and diffusion theory. The results of all three calculations were in excellent agreement.**

#### **4. CALCULATIONAL RESULTS**

### **4.1 Fuel Cycle**

**A plot of keff vs. integrated reactor power is shown in Fig. 2.** The graph shows that once equilibrium concentrations of <sup>135</sup>Xe and <sup>149</sup>Sm have **been reached, keff increases with burnup until a maximum value is obtained at about 4000 MWd and thereafter decreases with burnup. This behavior results from the fact that the burnable poison, <sup>167</sup>Er, burns out faster than <sup>235</sup> U. In order to determine the length of such a fuel cycle it is necessary to decide how much excess reactivity is required at the end of cycle. Since the calculations are performed with cross sections representative of hot fuel and with equilibrium xenon and samarium, one needs only that excess reactivity at end of cycle (EOC) sufficient to compensate for the absorption of experiments and to provide for xenon override. In this study it has been assumed that 2X excess is required. Therefore, the non-equilibrium cycle length is calculated to be 8000 MWd (320 full-power days).**

**The equilibrium fuel cycle is much more economical and results in much smaller reactivity swings and power shifts. For the five-path fuel management scheme described earlier, Fig. 3 shows the end-of-cycle keff as a function of cycle length. For a fuel temperature of 800K and a 62 day cycle length the beginning-of-cycle (BOC) and EOC eigenvalues were found to be 1.0293 and 1.0194, respectively. For this case, the BOC and EOC <sup>235</sup> U and <sup>167</sup>Er burnup levels, power generated, and <sup>239</sup>Pu generation are tabulated for each fuel element in Tables 5 and 6. For example, the discharged fuel element from the first path generated a power of 0.622 MW, contained 33.3 g<sup>239</sup>Pu, and had <sup>235</sup> U and <sup>167</sup>Er burnup levels of 43.42 and 84.OZ, respectively. Of course, if less excess reactivity were to be needed at end of cycle, as, for example, with a small experiment load, a longer cycle length can be obtained. It should be noted that the power sums to ~0.5Z less than 25 MW in Tables 5 and 6 because some power is produced by captures in non-fissile regions.**

### **4.2 Neutron Flux Distribution**

**Thermal (E<sub>n</sub> < 0.625 eV), epithermal (0.625 <**  $E_n$  **< 5.53 keV), and fast (5.53 keV**  $\leq$  **E<sub>n</sub>**  $\leq$  **10.0 MeV) neutron flux distributions for the beginning of equilibrium cycle (BOC) are plotted in Figs. 4 to 8 for several axial midplane traverses through the core and reflector regions of the MPRR 25 MW LEU TRIGA reactor. Figure 1 shows the location of these traverses which are between columns F and G, at the center of columns H and J, and at the center of row 9. These BOC flux distributions were taken from the two-dimensional XY full-core REBUS calculation corresponding to the 62 day cycle length for a fuel temperature of 800 K with the core operating at 25 MW.**

**These figures also show EOC/BOC neutron flux ratio distributions. These plots show how the thermal flux in the core regions Increases with burnup. The higher average burnup of the core requires that the EOC fluxes be increased relative to the BOC in order to maintain the 25 MW power level.**

**Maximum and region-averaged thermal fluxes for several irradiation positions are shown in Table 7. Values are given for the BOC and EOC configurations.**

**Figure 8 shows the flux distribution through row 12 in the pool water region 1.835 cm from the beryllium reflector. Although not modeled in the XY calculations, this is the region where the beam tubes are to be located. Fluxes shown in this figure should be used with caution since they can be expected to be significantly smaller when leakage through the beam tubes is taken iato account.**

### **4.3 Safety-Related Parameters**

**This section presents the results of calculations on safety-related neutronic parameters needed for transient analyses of the 25 MW MPRR LEU TRIGA** reactor. These parameters include kinetic parameters ( $\beta$ -effective and the **prompt neutron lifetime), prompt negative temperature coefficients, isothermal feedback coefficients and power peaking factors. In most cases calculations were performed for fresh fuel, beginning-of-equilibrium-cycle (BOC) and endof-equilibrium cycle (EOC) cores for the 62-day cycle-length case.**

### **A.7.1 Kinetic Parameters**

The prompt-neutron lifetime  $(\ell_p)$ , the neutron generation time  $(\Lambda)$ , and **the effective delayed-neutron fraction (6eff) were calculated for fresh fuel, BOC, and EOC equilibrium cores using the two-dimensional diffusion theory perturbation capability of the ARC System.\*2 Table 8 shows the results of these calculations. For these calculations burnup-dependent atom densities were taken from the REBUS calculation for a cycle length of 62 days. The delayed neutron data, used in the calculation of Seff» were taken from Version V of ENDF/B. Delayed neutron constants for BOC and EOC equilibrium cores are si own in Table 9. All calculations were performed using 11-group cross sections.**

### **4.3.2 The Prompt Negative Temperature Coefficient**

**One of the characteristics of U-ZrH-Er TRIGA fuel Is its large prompt negative teraperature coefficent. For small diameter fuel pins, such as those proposed for the MPRR 25 MW TRIGA reactor, the primary contribution to the prompt negative temperature coefficient is a hardening of the thermal neutron spectrum resulting from an increase in the fuel temperature. The binding of the ZrH molecule is described in terns of a harmonic oscillator potential with excited states separated in energy by about 0.14 eV. Thus, the population of excited oscillator states increases with fuel temperature. Thermal neutrons scattered from excited ZrH molecules receive a boost in energy with a subsequent hardening of the neutron spectrum. With this spectral shift toward** higher energy, increased absorption in the  $\sim$ 0.5 eV double resonance of  $^{167}$ Er **occurs, resulting in a negative reactivity effect. Since the fuel pin is a solid uniform mixture of U-ZrH-Er, the negative reactivity effect as a function of temperature is prompt. This characteristic of TRIGA fuel provides a builtin safety feature in the event of an unplanned power transient.**

**To evaluate the prompt negative temperature coefficient, 11-group core cross sections were generated at various temperatures using the EPRI-CELL code which was described earlier. Cross sections for H in ZrH were created for temperatures of 296, 500, 800, 1000 and 1200 K using temperature-dependent S(a,@) data. Doppler broadening of the <sup>238</sup> U resonances, as well as those for the other uranium and plutonium isotopes, was determined by a resonance calcu**lation at each of the above temperatures. However, the EPRI-CELL code does **not permit an interpolation on temperature for resonances In the thermal neutron energy range, which is the case for <sup>167</sup>Er. Therefore, <sup>166</sup> Er and <sup>167</sup> Er resonances have been Ooppler broadened only at those temperatures for which these cross sections exist in the EPRI-CELL library, namely 293, 564, 886, 1100, and 1200 R. Thus, a mismatch exists between the temperatures for which the erbium resonances have been Doppler broadened and the temperatures at which the H (in ZrH), U and Pu cross sections apply. This mismatch is summarized below.**



**Except for Zr, all other core materials were assumed to remain at room temperature.**

**The core-isothermal prompt negative temperature coefficients were calculated for fresh fuel atom densities and cross sections and for REBUS atom densities corresponding to the BOC and EOC configurations for the 62-day cyclelength case. For these calculations it was assumed that the changes in core temperature are independent of position. The effect of this approximation on the value of the temperature coefficient is thought to be small, but has not been investigated.**

**Diffusion theory calculations of keff at each ZrH temperature were made** using the XY model of the 25 MW MPRR TRIGA Reactor (Fig. 1) and the appro**priate 11-group temperature-dependent cross sections. It was assumed that all control rods are fully withdrawn, experiment regions are water-filled, the fuel pin composition is at the specified temperature, and all other materials are at room temperature. It was also assumed that the axial bucklings are independent of temperature.**

**The calculated values of keff were fitted by the least squares process to a 3rd degree polynomial in temperature and the prompt negative temperature coefficient (Op) wa s evaluated as the derivative of the polynomial. The** prompt negative temperature coefficient decreases as a function of burnup **because of the depletion of <sup>167</sup>Er in the fuel.**

# **4.3.3 Equilibrium <sup>135</sup>Xe and 1>>9Sm Worths**

The reactivity worths of equilibrium concentrations of <sup>135</sup>Xe and <sup>149</sup>Sm **were evaluated for the BOC configuration using REBUS-calculated atom densities for the case of a 62 day cycle length at a fuel temperature of 800K. Table 10 gives the results.**

### **4.3.4 Isothermal Reactivity Feedback Coefficients**

**Isothermal feedback coefficients were evaluated for the combined effects of temperature and density changes in the water moderator. These reactivity changes are the results of two physical effects:**

- **1. The hardening of the thermal neutron spectrum resulting from an increase in the water temperature.**
- **2. The increase in neutron leakage resulting from a reduction in the density of the water as it heats (or boils).**

**Using 11-group cross sections generated for various water temperatures in the core, XY diffusion calculations were performed with fresh fuel atom densities to evaluate the feedback coefficients. Table 11 shows the feedback coefficients for the combined effects of temperature and water density** changes. In this table,  $\delta \rho = (k_2 - k_1)/k_1k_2$  is the change in reac**tivity related to changes in core water temperature and density. The temperature and density of the reflector and flux trap water were not allowed to vary. RZ calculations were performed at each water density to determine the axial extrapolation distances needed for the XY calculations. ' '**

### **4.3.5 Power Peaking Factors**

**Radial, axial, and local power peaking factors have been calculated for the MPRR LEU TRIGA for the beginning-of-cycle (BOC) and end-of-cycle (EOC) equilibrium core. The radial peaking factor, Fr, is the ratio of the power** density at the hot spot on the axial midplane to the average midplane power **density, as calculated in XY diffusion theory problems. The axial peaking factor, Fa, is just the peak-to-average value of the chopped cosine axial shape. The local peaking factor, F£, is the radial peak-to-average power density in the local fuel element. Finally, the total peaking factor is the product of these three components.**

**Peaking factors were evaluated for BOC and EOC equilibrium cores. The** hot spot is in fuel element FE9G (see Fig. 1) adjacent to the water-filled **irradiation hole ES9H. Power peaking factors are given in Table 12. These are over-pessimistic values since in actual practice the irradiation positions would be fille d with aluminum or beryllium blocks containing small holes to accomodate samples so as to minimize power peaking effects. Figure 11 shows** how the power density varies in the X and Y directions across fuel element **FE9G. The Y-traverse is 0.734 cm from the core-water interface. Note the very large power peak in fuel next to the water-filled irradiation hole.**

### **4.4 Control Rod Worths**

**The results of the control rod worth calculations are summarized in Table 13. In the BOC configuration, rod C9F is the most reactive; when i t is stuck out, the worth of the remaining five rods is 6.47% 6p. In the beginning-of-cycle condition the five inserted rods should be able to shut down the reactor with all experiments removed, with all xenon decayed, and with the fuel cold. For the 62-day cycle-length case, the BOC excess reactivity is 2.85\* 6p with an 800 K fuel temperature (Section 4.1), the xenon worth is 2.50% δρ (table 10), and the increase in reactivity upon cooling of the fuel meat to room temperature is 1.92Z 6p (Fig. 9), giving a maximum excess reactivity of 7.27% 6p. Therefore, if the control rod worth calculations are correct, there is not an adequate shutdown margin when one rod is stuck out of the core. Also, the fresh core has an inadequate shutdown margin with one rod stuck. In relation to the accuracy of the control rod worth calculations i t must be emphasized that no comparisons with measured data have been made for control rod worths in LEU TRIGA cores.** However, the same methods for individual borated stainless steel rods in the **LEU core of the Ford Nuclear Reactor at The University of Michigan yielded worths within 0.2%** *&p* **of the measured values.1?**

**Higher-worth control rods of a different design could be considered.** For example, higher worth rods would result if the borated stainless steel poison material were in the shape of a square annulus about 7 cm on a side **and 1 cm thick with a water hole at the center. The water hole serves to thermalize and trap fast neutrons which penetrate the borated stainless stee l annulus. Relative to the cylindical rod, the square shape of the borated stainless stee l absorber provides a greater surface area and this too tends to increase the value of the rod worth. However, no calculations were made for the worth of control rods of this design.**

### 5. **CONCLUSIONS**

**In all aspects except for the shutdown margin, the 25-MW LEU TRIGA Multipurpose Research Reactor performs very well. The high uranium density of the U-ZrH-Er fuel with its burnable poison makes possible a long equilibrium cycle length with a relatively small reactivity swing. Therefore, control rod movement is minimized during the cycle, leading to a stable flux. The lack of adequate shutdown margin can probably be remedied by the use of a higher-worth design of the control rods.**

#### References

- 1. The 25-MW TRIGA Multipurpose Research Reactor for the Bandung Research Center, Bandung, Indonesia, General Atomic Company.
- 2. Guidebook on "Research Reactor Core Conversion from the Use of Highly Enriched Uranium to the Use of **Low** Enriched Uranium Fuels, " Internationa l Atomic Energy Agency Report IAEA-TECDOC-233, (August 1980).
- 3. B. A. Zolotar, et al., "EPRI-CELL Code Description, "Advanced Recycle Methodology Program System Documentation, Part II, Chapter 5, Electric Power Research Institute (September 1977).
- 4. G. D. Joanou and J. S. Dudek, "GAM-I: A Consistent P<sub>1</sub> Multigroup Code for the Calculation of Fast Neutron Spectra and Multigroup Constants," General Atomic Company Report GA-1850 (June 1961).
- 5. H. C. Honeck, "THERMOS, A Thermalization Transport Theory Code for Reactor Lattice Calculations," Brookhaven National Laboratory Report BNL-5826 (September 1961).
- 6. H. Henryson III, B. J. Toppel, and C. G. Stenberg, "MC<sup>2</sup>-2: A Code to Calculate Fast Neutron Spectra and Multigroup Cross Sections, " Argonne National Laboratory Report ANL-8144 (June 1976).
- 7. N. M. Greene, et al., "AMPX: A Modular Code System for Generating Coupled Multigroup Neutron and Gamma Libraries from ENDF/B," Oak Ridge National Laboratory Technical Memorandum ORNL-TM-3706 (November 1974).
- 8. R. E. MacFarlane, et al., "The NJOY Nuclear Data Processing System: User's Manual," Los Alamos Scientific Laboratory Report LA-7584-M (December 1978).
- 9. R. B. Hosteny, "The ARC System Fuel Cycle Analysis Capability, REBUS-2," Argonne National Laboratory Report ANL-7721 (1978).
- 10. R. D. O'Dell, F. W. Brinkley and D. R. Marr, "Users' Manual for ONEDANT: A Code Package for One-Dimensional Diffusion-Accelerated, Neutral-Particle Transport," Los Alamos Scientific Laboratory, October 1980.
- 11. R. E. Prael and L. J. Milton, "A User's Manual for the Monte Carlo Code VIM," Argonne National Laboratory Technical Memorandum FRAM-TM-84 (February 1976).
- 12. T. A. Daly, C. G. Stenberg, D. E. Neal, D. A. Schoengold, and G. K. Leaf, "The ARC System Two-Dimensional Adjunct Calculations, " Argonne National Laboratory Report ANL-7722 (October 1972).
- 13. M. M. Bretscher and J. L. Snelgrove, "Comparison of Calculated with Measured Quantities for the LEU-Fueled Ford Nuclear Reactor," International Meeting on Research and Test Reactor Core Conversions from HEU to LEU Fuels, Argonne, Illinois, November 8-10, 1982 (to be published).



 $\ddot{\phantom{a}}$ 

# **Table 1. Data For The 36-Rod TRIGA Fuel Cluster**

 $\sim$ 

 $\bar{z}$ 



 $\sim$ 

 $\frac{\partial \mathbf{q}}{\partial \mathbf{r}} \left( \mathbf{r} \right) = \frac{\partial \mathbf{r}}{\partial \mathbf{r}} \left( \mathbf{r} \right) = \frac{\partial \mathbf{r}}{\partial \mathbf{r}} \left( \mathbf{r} \right)$ 

 $\sim 10^7$ 

**Table 2. Group Structure for 5- and 11-Broad-Group Cross Section Sets**

 $\sim$ 

 $\label{eq:2} \frac{1}{2} \sum_{i=1}^n \frac{1}{2} \sum_{j=1}^n \frac{1}{$ 

 $\ddot{\phantom{a}}$ 



 $\Delta$ 

 $\ddot{\phantom{a}}$ 

 $\infty$ 

# Table 3. Extrapolation Distances and Corresponding Bucklings

 $\mathcal{A}^{\mathcal{C}}$  $\mathcal{L}$ 

 $\bar{\mathcal{A}}$ 



 $\epsilon$ 

## Table 4. Fuel Element Positions Arranged in the Order of Increasing Burnup for the Five Path Fuel Management Scheme

 $\overline{\phantom{a}}$ 

\* The five columns give the fuel shuffling sequence from top to bottom.

 $\overline{\phantom{a}}$ 

 $\cdot$ 

**Table 5. BOC Fuel Element 23 5 167- U Mass, \*"'Er Mass, Burnup and Power for the MPRR 25MW LEU TRIGA Reactor**

**Case: Fuel Temp** 800K, Cycle Length = 62 days days,  $k_{\text{eff}}$ (EOC) = 1.019



 $\ddot{\phantom{a}}$ 



**Table 6. EOC Fuel Element <sup>2</sup>3 <sup>5</sup>U Mass, <sup>167</sup>Er Mass, Burnup and Power for the MPRR 25MW LEU TRIGA Reactor** Case: Fuel Temp = 800K, Cycle Length = 62 days days,  $k_{\text{ref}}(EOC) = 1.019$ 



 $\ddot{\phantom{0}}$ 

 $\bar{\mathbf{r}}$ 

 $\mathcal{L}^{\mathcal{L}}$ 

 $\ddot{\phantom{a}}$ 

**Table 7. Thermal Neutron Fluxes (En < 0.625 eV) In the Axial Midplane for the 25 MM LEU TRIGA Reactor**

 $\frac{1}{4}$ 

Quantity	Fresh <b>Fuel</b> $K = 1.086366$	<b>BOC</b>	<b>EOC</b>
Fuel Temperature, <sup>o</sup> K	296	800	800
Neutron Generation Time, us	28.62	31.23	32.15
Prompt Neutron Lifetime, us	31.09	32.11	32.74
Effective Delayed Neutron Fraction	0.00733	0.00696	0.00687

Table 8. MPRR 25 MW LEU TRIGA Reactor Kinetic Parameters

÷.



 $\mathcal{L}^{\pm}$ 

Table 9. Delayed Neutron Parameters for the MPRR 25 MM LEU TRIGA Reactor

 $\sim 10^{-10}$ 

 $\sim 10$ 

 $\sim 10^{11}$  km  $^{-1}$ 



 $\mathcal{A}$ 

**Table 10. <sup>135</sup>Xe and Jlt9Sm Reactivity Worths Beginning-of-Equllibrium Cycle MPRR 25 MM LEU TRIGA Reactor**



 $\overline{a}$ 

Table 11. MFRR LEU TRIGA Isothermal Feedback Coefficients for the Combined Effect of Hater Temperature and Density Changes



 $\ddot{\phantom{a}}$ 

 $\mathcal{A}$ 

 $\mathcal{L}$ 

 $\ddot{\phantom{a}}$ 

 $\sim$ 

Table 12. Radial, Axial, and Local Power Peaking Factors for the MPRR 25 MW LEU TRIGA Reactor Equilibrium Core



 $\sim$ 

and the state of the state of

**k l <sup>k</sup>2**

 $\bar{\mathcal{A}}$ 

**Table 13. Control Rod Worths for the MPR 25 MW LEU TRIGA Reactor**

 $\bar{z}$ 

 $\langle \downarrow \rangle$ 

 $\label{eq:2.1} \mathcal{L}(\mathcal{L}(\mathcal{L}^{\mathcal{L}})) = \mathcal{L}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})) = \mathcal{L}(\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}))$ 

 $\sim$ 



Fig. 1..XY Model of the 25-MW TRIGA Multipurpose Research Reactor (Figure taken from Ref. 1)

![](_page_23_Figure_0.jpeg)

Fig. 2  $k_{eff}$  as a function of reactor integrated power for the non-equilibrium fuel cycle.

![](_page_24_Figure_0.jpeg)

Fig. 3 EOC  $k_{eff}$  as a function of fuel cycle length for the five-path fuel management scheme.

![](_page_25_Figure_0.jpeg)

Fig. 4a. MPRR 25 MW LKU TRIGA Core with Water Filled Central Flux Trap, Beginning of Equilibrium Cycle (BOC) Midplane Flux Traverses are Between Columns F and G.

![](_page_25_Figure_2.jpeg)

À,

![](_page_26_Figure_0.jpeg)

![](_page_26_Figure_1.jpeg)

Fig. 5b. E0C/B0C Midplane Flux Ratio Distribution at the Center of Column H.

![](_page_27_Figure_0.jpeg)

![](_page_27_Figure_1.jpeg)

**Fig. 6b. EOC/BOC Midplane Flux Ratio Distribution at the Center of Column J.**

![](_page_28_Figure_0.jpeg)

![](_page_28_Figure_1.jpeg)

![](_page_28_Figure_2.jpeg)

**cT£ 3 H- ro rt g § ^ O 0) M (B • O 9 > rt e** with<br>**Midpla**<br>of Row **CD « s s: • it to x 4 O I-" rt ix <sup>H</sup>**

**O H ( 0) 3 2 5«c « \* (tw**

 $\bullet$ 

 $\bullet$ 

![](_page_29_Figure_0.jpeg)

Fig. 8. B0C Midplane Flux Distributions Through Row 12 in the Water Reflector Region Adjacent to the Beryllium Reflector for the MPRR LEU TRIGA Reactor. This is the Region Where Beam Tubes Would be Located.

![](_page_30_Figure_0.jpeg)

Fig. 9.  $k_{eff}$  as a function of fuel temperature for fresh fuel, BOC, and EOC cores.

 $\Box$ 

![](_page_31_Figure_0.jpeg)

Fig. 10. Prompt negative temperature coefficient as a function of temperature for fresh fuel,  $BOC<sub>p</sub>$  and EOC cores.

![](_page_32_Figure_0.jpeg)

Fig. lla. Power profile in X-direction through element FE9G which contains the core hot spot.

![](_page_33_Figure_0.jpeg)

Fig. lib. Power profile in Y-direction through element FE9G which contains the core hot spot.