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Excitation Function and Half-life for the Fission Isomer 240m Pu from the 238 U(a,2n)^{240m}Pu Reaction

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ABSTRACT

The excitation function for the fission isomer $240m_{\text{Pu}}$ produced in the reaction of 238 U with α particles, has been measured at laboratory incident energies ranging from 20.1 to 27.3 MeV. A value of $3.8_{-0.4}^{+0.0}$ nsec was obtained for the halflife of ^{240m}Pu. The results are compared with other measurements performed using different experimental methods. (hullow)

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Int. roduct ion

Z-1Um "J'he Pu fission isoiner has been studied in dctaij by several authors, and relative energies and spins of leveJs in the second potential well have been determined. For this **isomer, six measurements of half-Jifc and excitation function using different techniques were published, with half-lives** ranging from 2.4 to 9 nsec /1-6/. Cross-sections also show **disagreement.**

In this work, new measurements are presented for the excitation function and half-life of the ^{240m}Pu fission isomer, **formed in the reaction** $238_{U(\alpha,2n)} 240$ **m** Pu. Experiments were **performed al u-particle energies of 20.1, 21.1, 21.0, 12.L, 24.2, 2 5.2 and 2 7.3 McV (±0.1 MeV)**

Experimental Procedure

Alpha-particles from the variable energy Cyclotron of the Instituto de lingenharia Nuclear, Rio de Janeiro, with energies up to 2K MeV were used. The energy spread of the beam was about 100 keV. The target consisted of metallic natural U, which was evaporated on an ordinary 4 mg/cm² aluminum foil. Target thicknesses ranged from 24 to 130 μ g/cm². **Due to a cooling system and good thermal contact between target and target holder, currents up to 5 uA could be used** $171.$

To discriminate isomeric from prompt fission fragments, the fission in flight method was employed /8/, the geometry

of which is shown in the inset of fig. 1. The a-particle beam was collumned by two Smm tantalum slits, 200 cm apart, producing a beam spot of 3.5 mm diameter at the target.

The detector, a policarbonate foil (MAKROFOL), 100 pm thick, was placed on the inner wall of a 24.3 mm radius cylinder. Prompt fission fragments were monitored with a Si surface barrier detector, set at 60[°] to the incoming beam.

For some energies, two runs were performed to measure prompt and isomeric fission distributions with an accumulated charge of about 200 μ C and 20.000 μ C, respectively. The short run gave the 90⁰ prompt fission cross-section, which in turn was used to determine the isomeric to prompt fission ratio.

After irradiations the MAKROFOL foil was etched, according to procedures described in /9/. Optical scanning with a microscope was used to measure track densities.

The distribution of tracks produced by the fission fragments on the detector foil, allows the determination of the half-life and the cross sections for isomeric and prompt fission σ_j and $\sigma_{\bf n}$. The extrapolated ratios of isomeric to prompt fission track densities at 90° , yield the cross section ratios σ_i/σ_n . Through this normalization the effect of the Coulomb barrier on the excitation function is eliminated. For thick targets the fraction of isomeric nuclei which are not able to leave the target was taken into account. Monte-Carlo simulated curves /10/ for different half-lives were compared with the experimental ones. From this comparison the half-live and the cross section σ_i were obtained. The simulations took into

account corrections for the velocity distribution and angular spread of the recoil nuclei leaving the target, and for the finite target size.

An isotropic angular distribution for isomeric fission fragments was assumed, consistent with experimental results of ref. $/11/$. In the calculations it was assumed that the recoil ing nuclei, go into a cone with a half-opening of 45^0 and are uniformly distributed $/13/$. The results obtained for the halflife differ at most by o.l ns in comparison with calculations with no angular spread.

It was not necessary to include a minimum detectable angle of incident fragments for the following reasons: i) with optical scanning the observation cut-off for the angle of incidence is nearly 0^0 /6/; ii) for MAKROFOL the "critical angle of etching" is 3^0 /12/; iii) in the geometry used, an angle of incidence lower than or causal to 3^0 is highly unprobable.

Results and Discussion

In fig. I a typical experimental result for the track distribution of fission fragments is shown for a thick target of 130 μ g/cm². Simulated track distributions for half-lives of 3.4 , 3.8 and 4.4 ns are included (curves 1, 11, 111). A least squares fit to the experimental distribution gives a half-life of 3.8 ns with an error of about $\binom{10}{1}$ and $\binom{20}{4}$ and $\binom{20}{4}$ IV was calculated without corrections for the velocity distribution of the recoil nuclei, and for $T_{1/2} = 3.8$ ns).

For the half-life determination thick targets were used, for which the velocity spread ranges from zero to the maximum vaJue, given by momentum conservation. Thus no uncertainty exists for the lower limit of the velocity.

For cross-section measurements mainly thin targets were used, for which the range of the recoils was larger than tne target thicknesses. This ensures that all recoiling nuclei leave the target, thus reducing the experimental error.

In table I a comparison with other half-life measurements is given. Our value of 3.8 $_{-0.4}^{+0.6}$ ns is in good agreement with the results obtained with electronic techniques /2,5/. The measurements performed with different fission in flight experimental arrangements show less agreement with ours.

Cross sections σ_j and ratios $\sigma_j/\sigma_{\mathbf{n}}^{\top}$ have been measured foi energies from 20.1 to 27.3 McV. The results are shown in Table II and Figs. 2 and 3.

The excitation function (σ_i/σ_p) in Fig. 3 was compared with ref. /S/, and a discrepancy of a factor of two exists, at the most. Our results for σ_j (Fig. 2) are in good agreement with the results obtained in a different detection geometry with the recoil method /0/. Disagreement exists with experiments performed with electronic techniques /2, 5/.

We may conclude that our experiment represents a contribution for resolving the existing discrepancies in the experi. mental results, concerning the half-life and excitation function of the fission isomer $240m_{\text{Pu}}$.

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- Comparison of half-life measurements for $240m_{\text{Pu}}$ Table I: fission isomer
- Our experimental results for the Pu^{240m} delayed Table 2: fission cross section $(\sigma_{\frac{1}{2}})$ and the isomeric ratio (σ_i/σ_p)
- Figure 1: Experimental track distribution for prompt and $a)$ delayed fission of 240 Pu, E_{α} =27.3 MeV. Calculated curves are also shown for half-lives of 3.4ns(1), 3.8ns(11) and 4.4ns(111). Curve IV corresponds to a half-life of 3.8ns assuming a unique recoil velocity.
	- Diagram of target and MAKROFOL holders. \mathbf{b}) Isomeric nuclei that recoil into the region in front of the target undergo fission and produce tracks in the MAKROFOL.
- Figure 2: Excitation function for $u^{2.38}$ (⁴He, 2n)^{240m}Pu.
	- (. Namboodiri et al /6/, A Vandenbosch et al /2/, o this work).
- Figure 3: Isomeric ratio σ_{i}/σ_{p} for the reaction 238 U(α , 2n) 240m_{Pu} as a function of bombarding energy. $(\Box B\text{Filt et al}(5),$ o this work).

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Table 2

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Reaction^{-2.58}U^{-(\alpha,2n)^{240m}}Pu
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