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ON TWO - PHOTON WIDTH
OF δ - MESON

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Abstract

It is shown that a natural magnitude of the radiative width of the δ meson must be $\Gamma_{\delta\gamma\gamma} \approx 1$ KeV.

The reason for which the published experimental number for $\Gamma_{\delta\gamma\gamma}$ seems to be too low is presented.

The preliminary result of determination of the two-photon width of the scalar resonance δ (980) from the data on the process $e^+ e^- \rightarrow e^+ e^- \eta \pi^0$ is^{11/}.

$$\Gamma_{\delta\gamma\gamma} \cdot BR(\delta \rightarrow \eta\pi) = (0.100 \pm 0.025 \pm 0.100) \text{ KeV} \quad (1)$$

In a literature, one can find the estimates according to which $\Gamma_{\delta\gamma\gamma} \approx 5 \text{ KeV}^{12,31}$ if the δ meson is a bound quark-antiquark system, and $\Gamma_{\delta\gamma\gamma} \approx 0.27 \text{ KeV}^{13/}$ if the δ is a four-quark system $\bar{q}^2 q^2$. On a base of these estimates, the result (1) was interpreted as one testifying a four-quark nature of the δ meson^{14/}. In our opinion such a conclusion is not seem to be perfectly grounded.

Firstly, a quantitative evaluation in the framework of the $\bar{q}q$ scheme, based on using of an effective chiral Lagrangian gives the result^{15/}

$$\Gamma_{\delta\gamma\gamma}(m_\delta) = \frac{m_\delta}{\pi} \left(\frac{\alpha m_\delta}{12\pi F_\pi} \right)^2 \approx 1.3 \text{ KeV} \quad (2)$$

As to the estimate $\Gamma_{\delta\gamma\gamma} \approx 0.27$ in $\bar{q}^2 q^2$ scheme, it has been obtained without of computing any definite diagrams and it can be easily increased up to the value 1-2 KeV.

Really, evaluating the contribution of the diagrams of Fig.1 in $\Gamma_{\delta\gamma\gamma}$, we shall get (in approximation $m_\delta = 2 m_\pi$)

$$\Gamma_{\delta\gamma\gamma}^{(K\bar{K})} = \frac{\alpha^2 (\pi^2/4 - 1)^2}{16 \pi^2 m_\delta} \cdot \frac{g_{\delta K\bar{K}}^2}{4\pi} \quad (3)$$

Substituting the values $g_{\delta K\bar{K}}^2/4\pi = 2.3 \text{ GeV}^2$ or $g_{\delta K\bar{K}}^2/4\pi = 3 \text{ GeV}^2$ used in the $\bar{q}^2 q^2$ scheme¹⁸⁾, one gets $\Gamma_{\delta\gamma\gamma}^{(K\bar{K})} = 1.7 \text{ KeV}$ or $\Gamma_{\delta\gamma\gamma}^{(K\bar{K})} = 2.2 \text{ KeV}$ respectively.

In the $\bar{q}q$ scheme, where $g_{\delta K\bar{K}}^2/4\pi = 0.82 \text{ GeV}^2$, we should get $\Gamma_{\delta\gamma\gamma}^{(K\bar{K})} = 0.61 \text{ KeV}$.

Of course, a calculation of a loop diagram with one definite physical intermediate state does not permit yet to fix an exact value of amplitude, but such calculation determines a scale of natural values of the amplitude under consideration. We see that a scale of values of the amplitude $\delta \rightarrow \gamma\gamma$ is the same in $\bar{q}^2 q^2$ and $\bar{q}q$ schemes, and, consequently, the result (1) is too small to be explained in these schemes.

Let's show now that the data¹¹⁾ are compatible with the values of $\Gamma_{\delta\gamma\gamma}$ considerable larger than the value (1).

The result (1) was obtained under assumption that the total width of the δ meson is small: $\Gamma_\delta = 54 \text{ MeV}$ ¹⁹⁾. But in both schemes ($\bar{q}q$ and $\bar{q}^2 q^2$) a natural value of Γ_δ approximately is by six time more, and this fact does not contradict (see Ref.^{16-8, 10)} to observation of a narrow peak in the process $K\bar{p} \rightarrow \Sigma^+(1385) \eta \pi^-$.

If the δ resonance is wide ($\Gamma_\delta \sim 300 \text{ MeV}$), it must be described (see Ref.^{16-8, 10)} by formulae

$$\sigma_{\gamma\gamma \rightarrow \gamma\pi}(W) = 8\pi \Gamma_{\delta\gamma\gamma}^3(W) \Gamma_{\delta\eta\pi}^2(W) \cdot \left| W^2 - m_\delta^2 + i m_\delta \Gamma_\delta(W) \right|^{-2} \quad (4)$$

where

$$\Gamma_\delta \cong \Gamma_{\delta\eta\pi^0} + \Gamma_{\delta K\bar{K}}$$

and $\Gamma_{\delta\gamma\gamma}(W) = (W/m_\delta)^3 \Gamma_{\delta\gamma\gamma}(m_\delta)$,

$$\Gamma_{\delta\eta\pi^0} = \frac{g_{\delta\eta\pi^0}^2}{16\pi W} \left[1 - \frac{2(m_\eta^2 + m_\pi^2)}{W^2} + \frac{(m_\eta^2 - m_\pi^2)^2}{W^4} \right]^{\frac{1}{2}},$$

$$\Gamma_{\delta K\bar{K}}(W) = \frac{g_{\delta K\bar{K}}^2}{16\pi W} \begin{cases} \sqrt{1 - 4m_K^2/W^2}, & W \geq 2m_K \\ i\sqrt{4m_K^2/W^2 - 1} \left(1 - \frac{2}{\pi} \arctan \sqrt{4m_K^2/W^2 - 1} \right), & W < 2m_K \end{cases}$$

Then, it is possible to get a value of $\Gamma_{\delta\gamma\gamma}$ larger than the value (1).

The Fig.2 demonstrates an attempt to describe the data¹¹ in terms of two resonance cross sections, first of which was calculated using the formulae (3) and (4) with $\Gamma_{\delta\gamma\gamma}(m_\delta) = 0.5$ KeV and $g_{\delta\eta\pi^0} = -4.8$ GeV, $g_{\delta K\bar{K}} = -3.2$ GeV^{1/2}. The part of cross section connected with A_2 meson was computed using the standard formula for narrow resonance with $\Gamma_{A_2\gamma\gamma} = 0.6$ KeV and with other parameters from Ref.¹⁹. We have used a slightly less value of $\Gamma_{A_2\gamma\gamma}$ as compared to conventional mean value, because some part of events from area of A_2 meson must be treated as belonging to right branch of δ resonance $m_{\eta\pi^0}$ distribution if δ resonance is wide.

Thus, as it is seen from Fig.2 the data¹¹ do not contradict to $\Gamma_{\delta\gamma\gamma} = 0.5$ KeV. If one will take into account that the systematic mistake in determination of $\Gamma_{\delta\gamma\gamma}$ in Ref.¹¹

is estimated to be 100%, one can not exclude until now that the real value of $\Gamma_{\delta\delta}$ could be close to $\Gamma_{\delta\delta}(m_{\delta}) \approx 1 \text{ KeV}$ - a value predicted by the $\bar{99}$ scheme of the δ meson.

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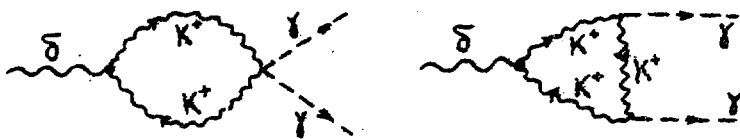


Fig. 1

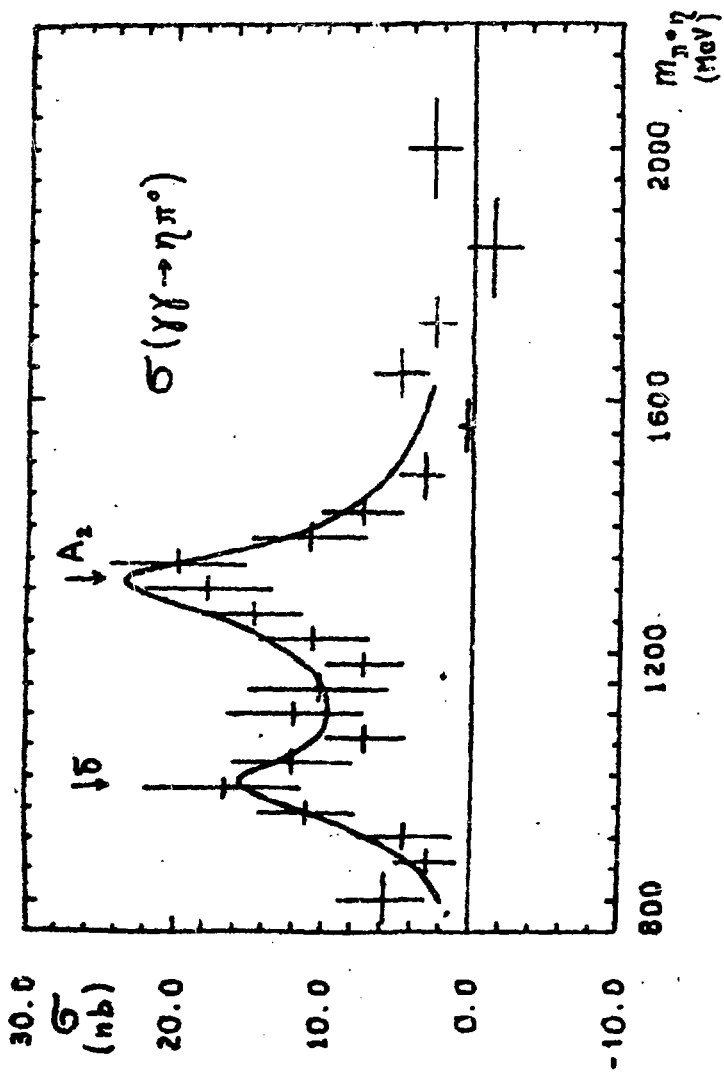


Fig. 2

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