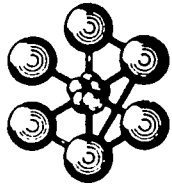


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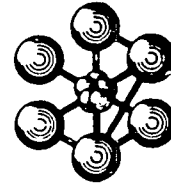
I.W.G.F.R Meeting
on Flow Induced Vibration
21-24 October 1986 - PARIS

VIBRATIONS OF THE SHAFTS OF THE
PRIMARY PUMPS :

"ROTOR" : A FINITE ELEMENT PROGRAM
FOR COMPUTING THE VIBRATIONS OF
ROTATING SHAFTS IMMERSSED IN A LIQUID

by F. AXISA (C.E.A / DENT)

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AIMS OF "ROTOR" IS TO
COMPUTE :

- The complex flexural eigenmodes of the shaft
- The critical angular velocities
- The instability thresholds
- The vibration response to a given unbalance

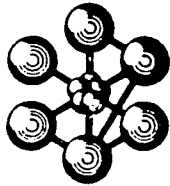
TAKING INTO ACCOUNT :

- The gyroscopic effect
- The various fluid effects
- The support conditions of fluid bearings

FUTURE EXTENSIONS

- Dynamical coupling of the shaft with the non rotating parts of the pump
- Time histories for transient excitation both in the linear and non linear domains

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GENERAL FEATURES OF THE PROGRAM ROTOR

PHYSICAL PROBLEM

A shaft supported by liquid film bearings and rotating in a liquid is a non conservative mechanical system as flexural vibrations are concerned:

Eigen Frequencies ω are complex in nature and depend on the shaft angular velocity Ω

$$\omega_c = \omega_R(\Omega) + i\omega_I(\Omega)$$

UNACCEPTABLE VIBRATION CAN OCCUR AT :

- Critical angular velocities : $\omega_R(\Omega) = \Omega$
- Static instability : $\omega_I(\Omega) = 0$
- Dynamic instabilities : $\omega_I(\Omega) = 0$

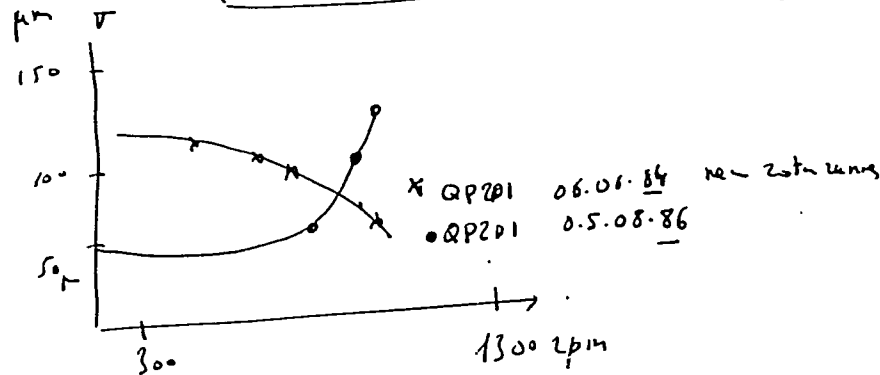
KBG: Company KandK. reactor Pump.

Continuum increase in vibration from 275 μ m

100 μ m \rightarrow 300 μ m

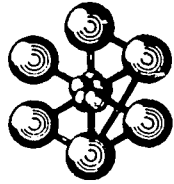
2 level.

stem vibration max!
hand receiver unbalance

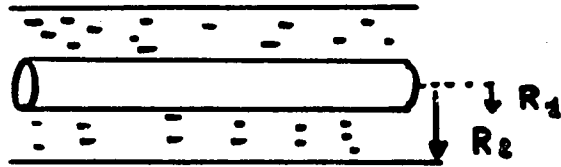


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NON ROTATING SHAFT :
FLUID EFFECTS



Inertial effect

$$m_a = \rho \pi R_1^2 \frac{R_1^2 + R_2^2}{R_2^2 - R_1^2}$$

For small annular gap $R_2 > R_2 \approx R_1$, $R_2 - R_1 = e$

$$m_a \approx m_f \frac{R}{e}$$

Compressibility effect

negligible as far as $\omega R / c \ll 1$

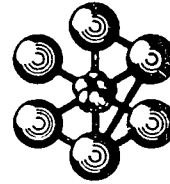
Viscosity effect

small as far as $\omega e^2 / \nu \gg 1$

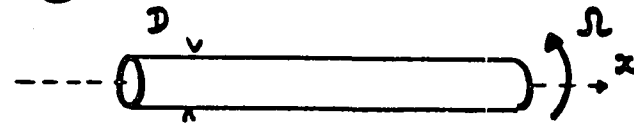
Leading to a modal damping

$$\zeta = \frac{1}{2} \frac{m_a}{M + m_a} \sqrt{\frac{2\nu}{\omega e^2}}$$

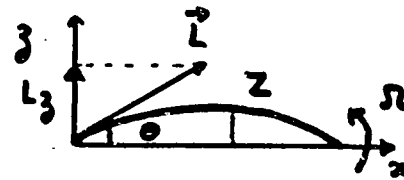
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ROTATING SHAFT :
GYROSCOPIC EFFECT



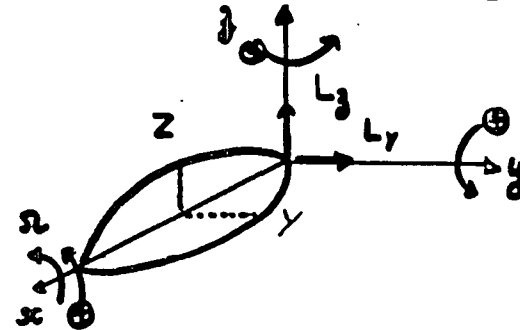
Angular momentum $\vec{L} = J \vec{\Omega}$; $J = MD^2/8$



$$L_3 = J \Omega \theta = J \Omega \frac{\partial z}{\partial x}$$

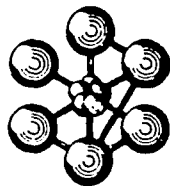
Equilibrium $\Rightarrow \mu$ (flexion in oxy plane) $= \dot{L}_3$

A.S. COUPLED FLEXURAL EQUATIONS



$$EI \frac{d^4 z}{dx^4} - \omega^2 \rho S z + i \omega \Omega \rho J \frac{d^2 y}{dx^2} = 0$$

$$EI \frac{d^4 y}{dx^4} - \omega^2 \rho S y - i \omega \Omega \rho J \frac{d^2 z}{dx^2} = 0$$



SHAFT AT REST : $\Omega = 0$

I CLASSICAL F. E. DISCRETISATION PROVIDES

$[K_s]$: Structural stiffness matrix

$[M]$: Structural + incompressible fluid mass matrix

$[A]$: Damping matrix, in particular for viscosity effect of the fluid

Shaft modeled as a beam

Liquid modeled in 2D in a section perpendicular to the shaft axis (3D corrective coeffs are possible for variation of section along the shaft)

II RESOLUTION OF THE CONSERVATIVE EIGENVALUE PROBLEM

$$\{ [K_s] - \omega^2 [M] \} X^{(1)} = 0$$

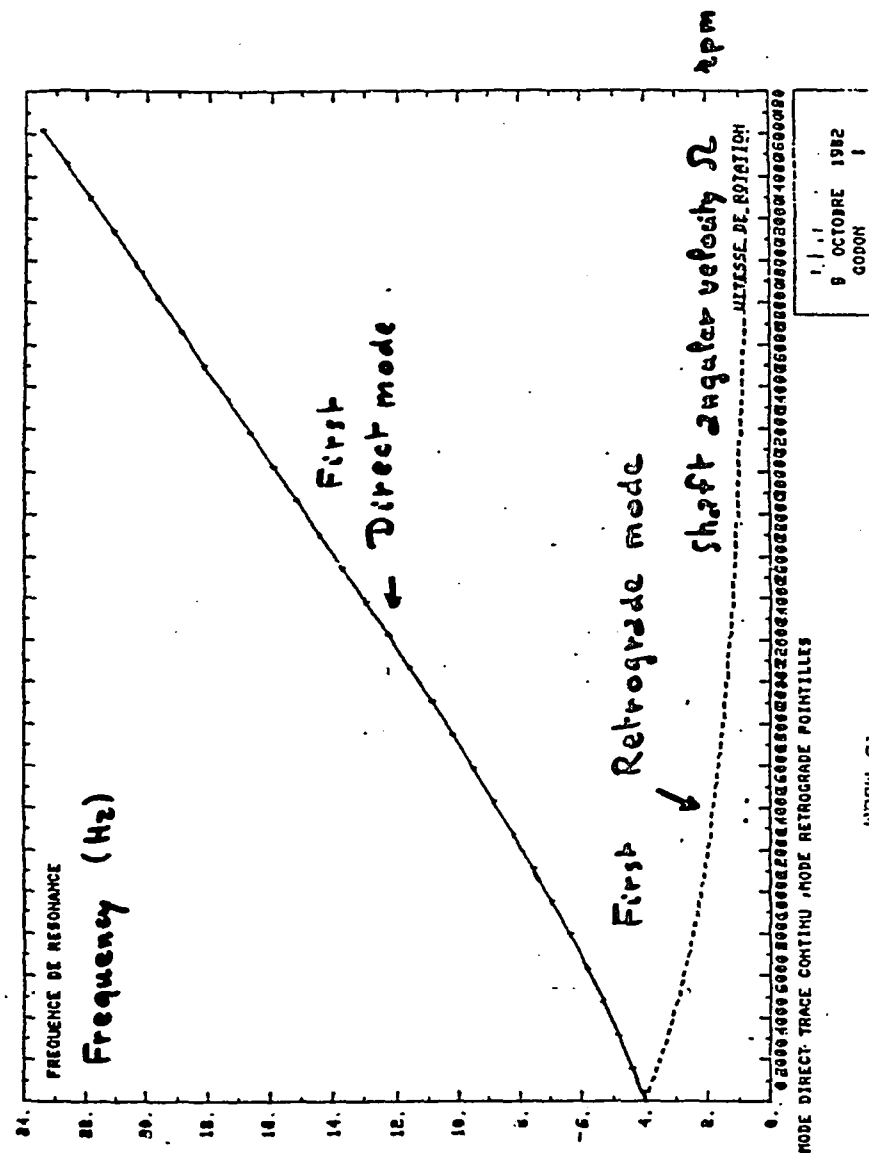
$\Rightarrow \omega_n > 0$ and $\varphi_n^{(1)}$ (real quantity)

III MODAL DAMPING BY MODAL PROJECTION

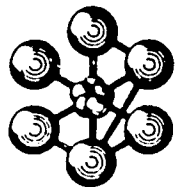
$$A_n = [\varphi_n]_T [A] [\varphi_n]$$

$$\zeta_n = A_n / 2 M_n \omega_n \quad M_n = [\varphi_n]_T [M] [\varphi_n]$$

ROTOR COMPUTATION : GYROSCOPIC EFFECT



GYROSCOPIC COUPLING



CHANGE OF VARIABLES

$$\Phi_D = Y + iZ \quad \Phi_R = Y - iZ$$

For a uniform shaft one obtains the decoupled system

$$\left[EI \frac{d^4}{dx^4} - \omega \Omega J \frac{d^2}{dx^2} - \omega^2 \rho S \right] \Phi_D = 0$$

$$\left[EI \frac{d^4}{dx^4} + \omega \Omega J \frac{d^2}{dx^2} - \omega^2 \rho S \right] \Phi_R = 0$$

Direct mode $\Phi_D = \Phi(x) (1+i)$

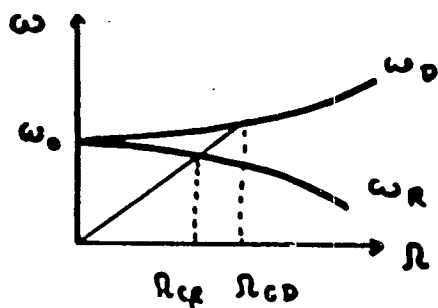
Flexural plane rotates at $+\Omega$

Retrograde mode $\Phi_R = \Phi(x) (1-i)$

Flexural plane rotates at $-\Omega$



ω_D et ω_R remains real and positive



$$\frac{\omega_D - \omega_R}{\omega_0} \approx \frac{1}{2} \frac{\rho(\Omega R)^2}{E}$$

Pinned-pinned uniform beam

**ROTATING SHAFT
COROTATING FLUID EFFECT**



To the 1st order in respect with fluctuating quantities:

$$e \frac{\partial w}{\partial \theta} - \frac{1}{2} \Omega R \frac{\partial \vec{X} \cdot \vec{n}}{\partial \theta} - R (\vec{X} \cdot \vec{n}) = 0$$

$$e R \frac{\partial w}{\partial t} + \frac{e}{\rho} \frac{\partial F}{\partial \theta} + \Omega R e \frac{\partial w}{\partial \theta} - \frac{1}{2} (\Omega R)^2 \frac{\partial \vec{X} \cdot \vec{n}}{\partial \theta} - \frac{1}{2} \Omega R (\vec{X} \cdot \vec{n}) = 0$$

Elimination of fluctuating velocity

$$\rho \frac{e}{\rho} = R^2 \ddot{X} \cos \theta - \Omega R^2 \dot{X} \sin \theta - \frac{\Omega^2}{4} R^2 X \cos \theta$$

Inertia

gyroscopic coupling

stiffness

Integration of pressure leads to the coupled eigenvalue problem:

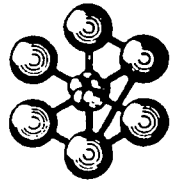
$$EI \frac{d^4 X}{dx^4} - m_2 \frac{\Omega^2}{4} X + i \omega m_a \Omega Y - \omega^2 (\rho S + m_n) X = c$$

$$EI \frac{d^4 Y}{dx^4} - m_2 \frac{\Omega^2}{4} Y - i \omega m_a \Omega X - \omega^2 (\rho S + m_n) Y = 0$$

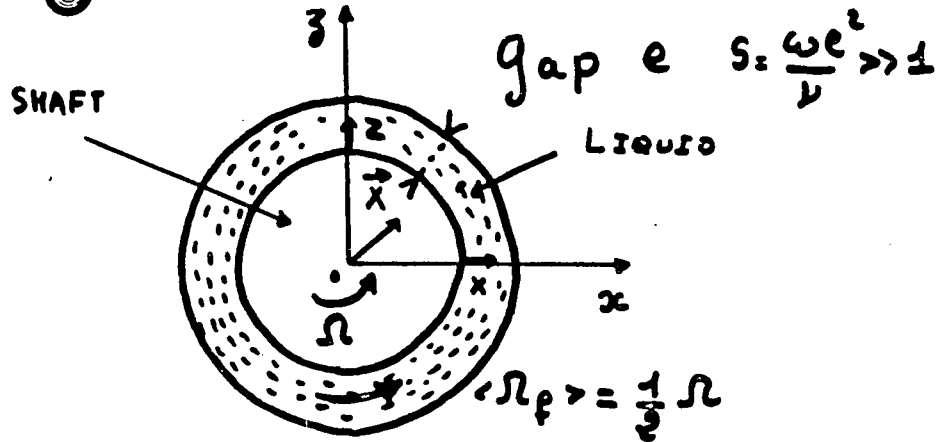
$$\begin{bmatrix} K(\Omega) - \omega^2 M & +i\omega C \\ -i\omega C & K(\Omega) - \omega^2 M \end{bmatrix} \begin{bmatrix} X \\ Y \end{bmatrix} = 0$$

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ROTATING SHAFT



COROTATING FLUID EFFECT

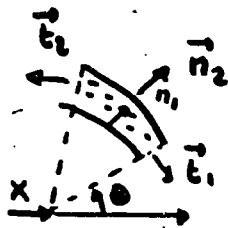


Flow velocity in the presence of a lateral vibration of the shaft

$$\vec{V} = \left(\frac{1}{2} \Omega z + w \right) \vec{e} + u \vec{n}$$

Pressure field

$$P = P_0 + p(t, z, \theta)$$

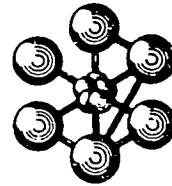


Continuity equation $\iint_{\Sigma} \vec{V} \cdot \vec{k} d\sigma = 0$

Momentum equation

$$\iiint \frac{\partial \vec{V}}{\partial t} dz + \iint_{\Sigma} \vec{V} (\vec{V} \cdot \vec{k}) d\sigma + \frac{1}{\rho} \iint_{\Sigma} P \vec{k} d\sigma = 0$$

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MODAL ALGORITHM IN THE
COMPLEX PLANE :
DISCRETISED LINEAR OPERATOR

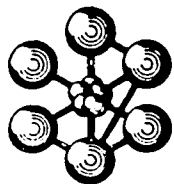
$$\mathcal{L} [\Omega, \omega] =$$

$$\left\{ [K_s + K_B(\Omega)] - \Omega^2 K_F \right\} + 2i\omega [A_{FS} + A_B(\Omega)] + \Omega [G_s + G_F] - \omega^2 [M_s + M_F] \} \Phi = c$$

PROBLEM :

COMPUTE THE EIGENMODES Φ_n, ω_n VERSUS Ω

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ROTATING SHAFT :
COROTATING FLUID EFFECTS



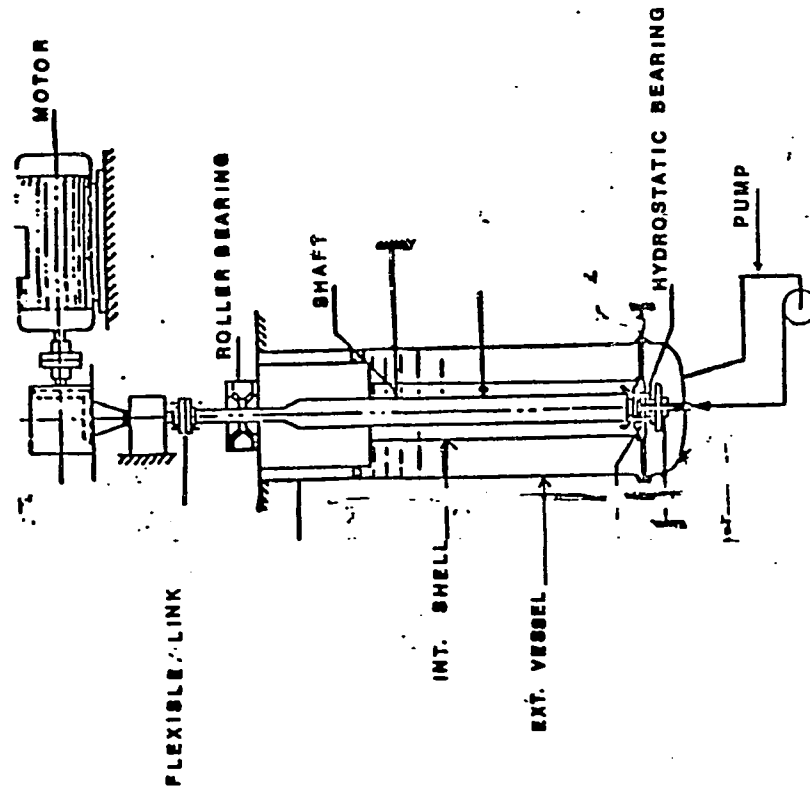
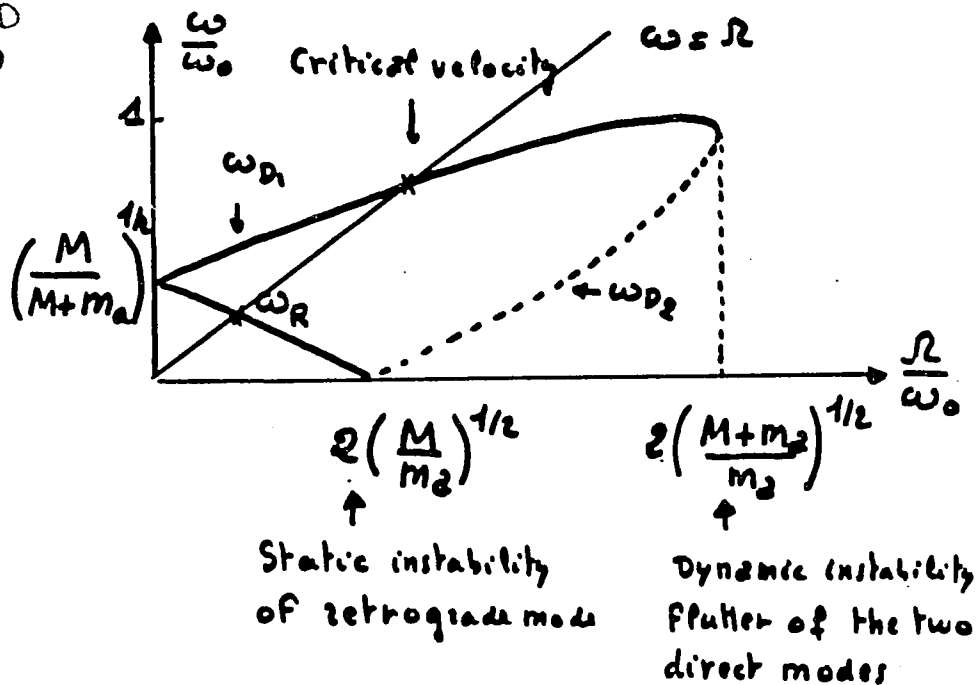
For uniform shaft and fluid annulus decoupling is readily achieved

Direct mode $\varphi_D = X + iY$

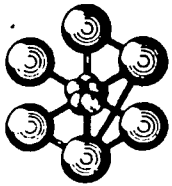
Retrograde mode $\varphi_R = X - iY$

ω_0 natural pulsation of the shaft at rest and in air

M: structural mass, m_a added fluid mass/unit length



MALVINA TEST RIG



MODAL ALGORITHM IN THE COMPLEX PLANE :
STEP BY STEP IN Ω ITERATIVE PROCEDURE



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next step

Velocity step $n \quad \Omega = n \Delta \Omega$
 (starting from $n=1 \quad \Delta \Omega \ll \Omega_{final}$)

Iteration k provides ω_k, Φ_k

Approximation $k+1$ of Φ satisfies:

$$\mathcal{L}[\Omega, \omega_k] \Phi_{k+1} = \Phi_k$$

then approximation $k+1$ of ω satisfies:

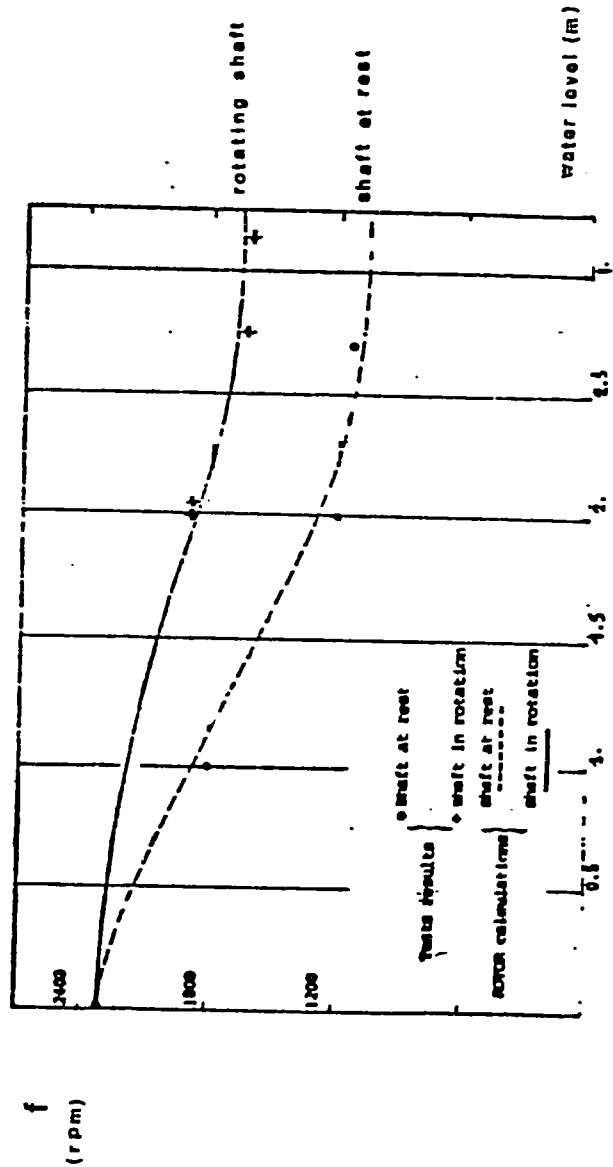
$$\langle \Phi_{k+1}^*, \mathcal{L}[\Omega, \omega_{k+1}] \Phi_{k+1} \rangle = 0$$

CONVERGENCE :

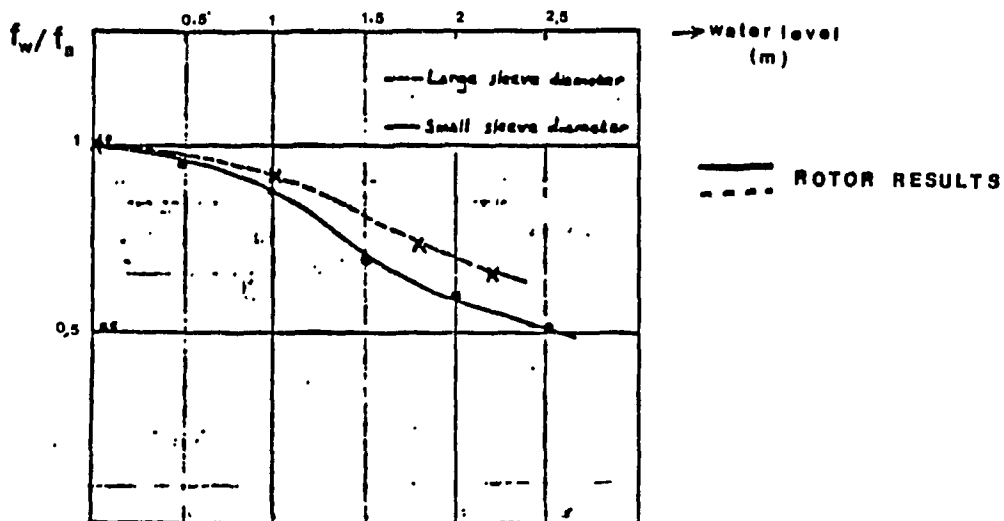
$$|\omega_{k+1} - \omega_k| \rightarrow 0 \quad \text{and} \quad \|\Phi_k\| \rightarrow \infty$$

INITIALISATION :

$$\Omega = 0 \Rightarrow \varphi_0, \omega_0 \rightarrow \begin{cases} \Phi_0 = \varphi_0 (1+i) \\ \Phi_k = \varphi_0 (1-i) \\ \omega_1 = \omega_0 (1+i^3) \end{cases}$$



MALVINA TESTS

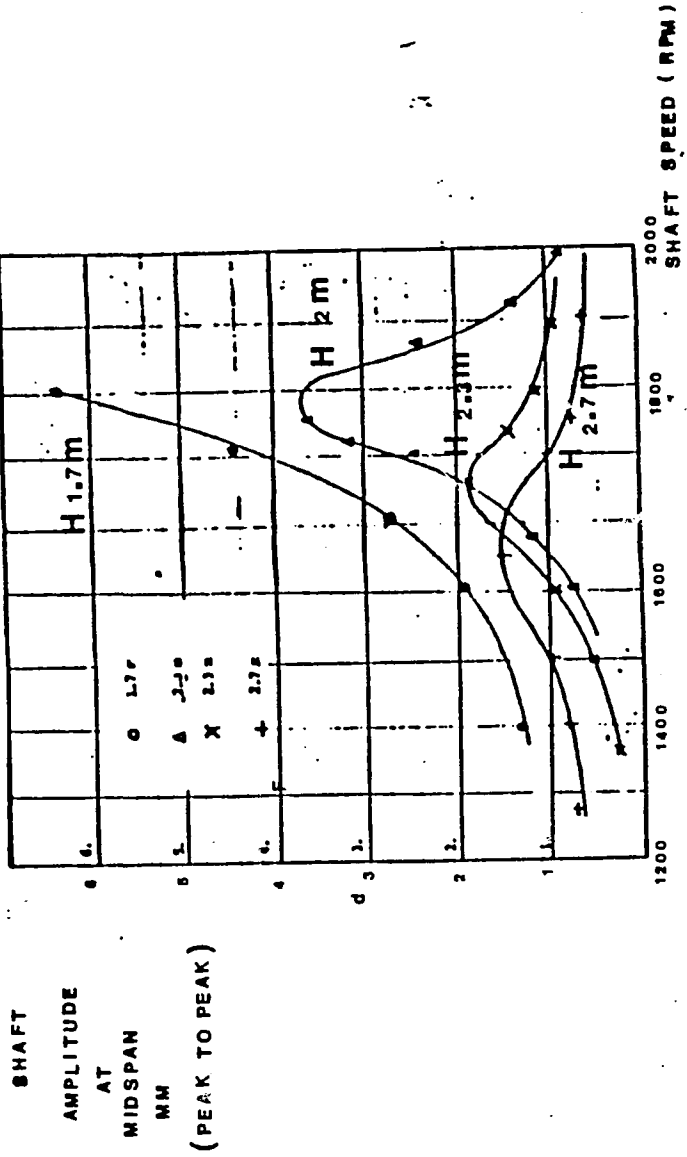


Eigenfrequencies variation for
the shaft at rest

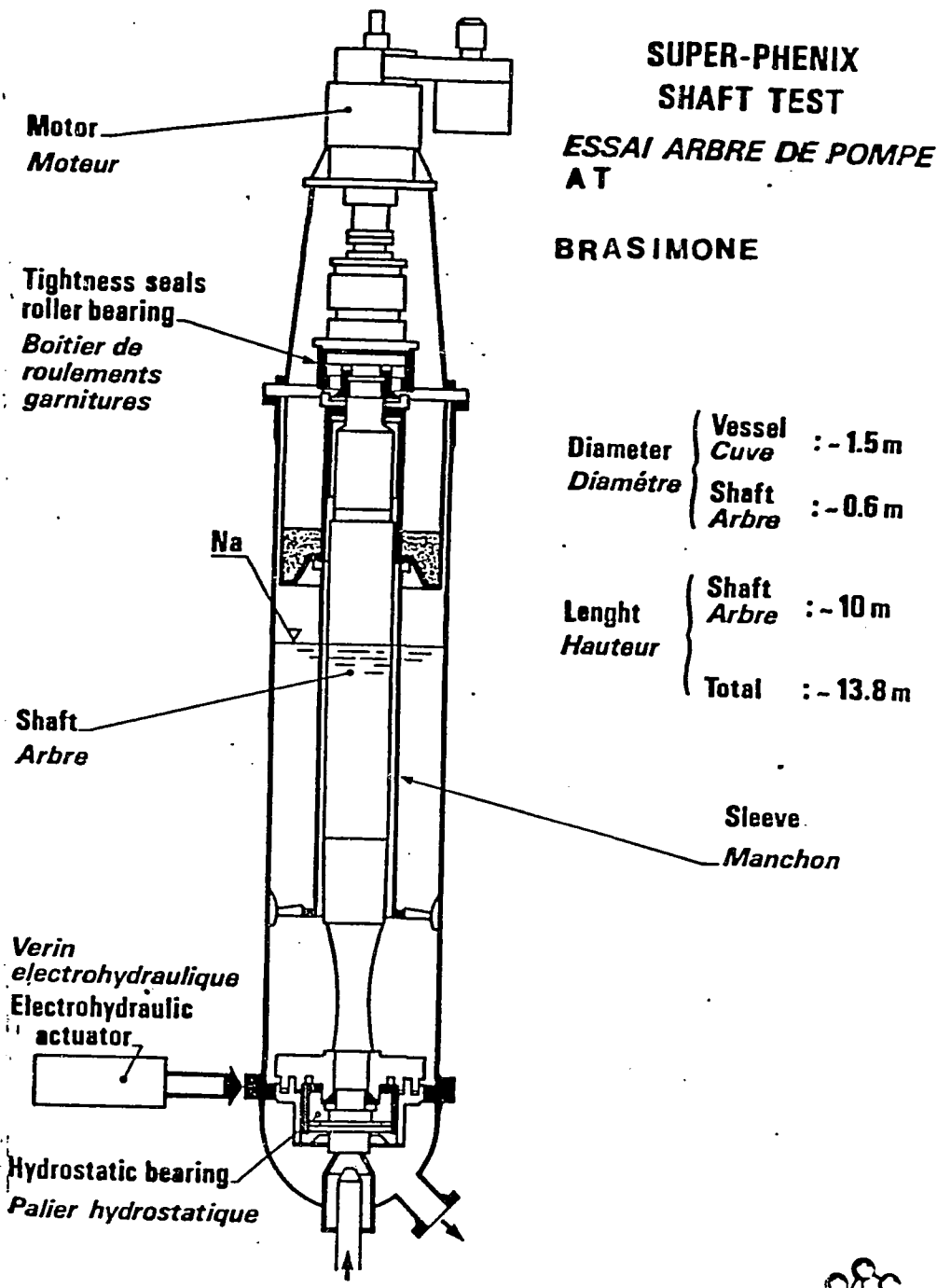
Table 1 : Experimental results

Water level (m)	Critical speed (rpm)	Mean Displacements (mm)		Mean Damping	
		at mid-shaft	at lower bearing	at mid-shaft	at bearing
1.7	>1950	>7.6	>0.50	/	/
2	1875	3.6	0.38	3.9	4.1
2.13	1800	2.0	0.36	/	/
2.30	1740	1.8	0.28	6.2	8.3
2.70	1630	1.3	0.14	7.3	10.3

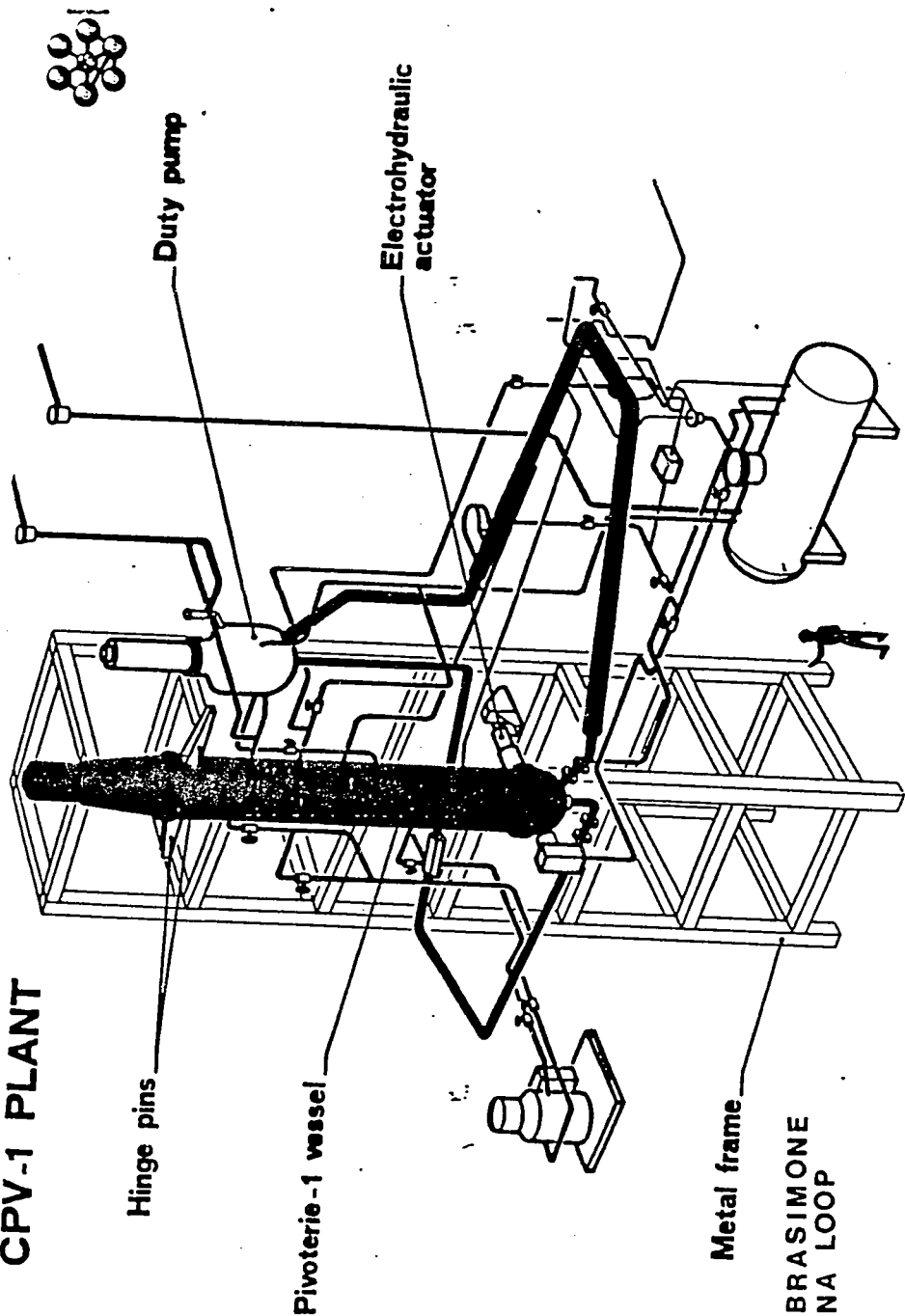
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MALVINA
crossing of critical speed



CPV-1 PLANT



7.68

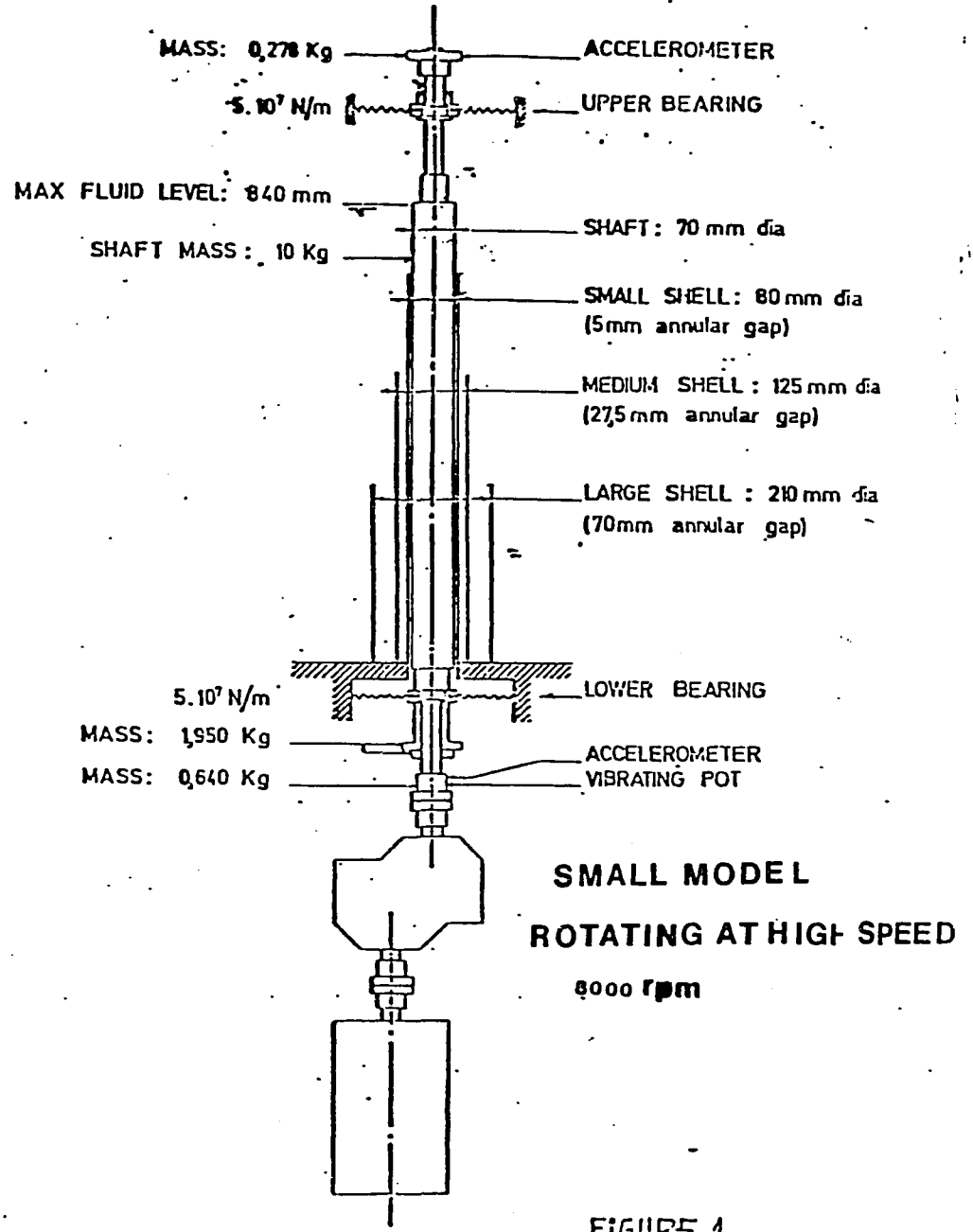
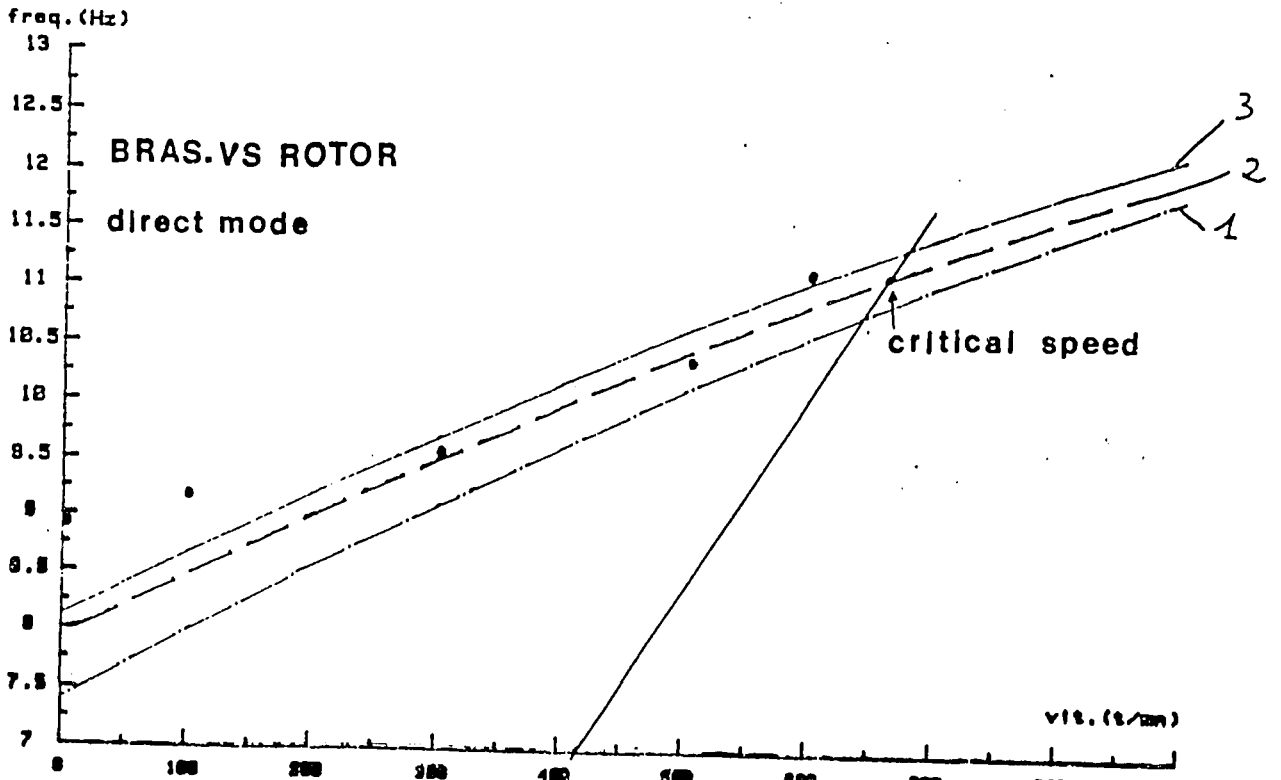


FIGURE 1

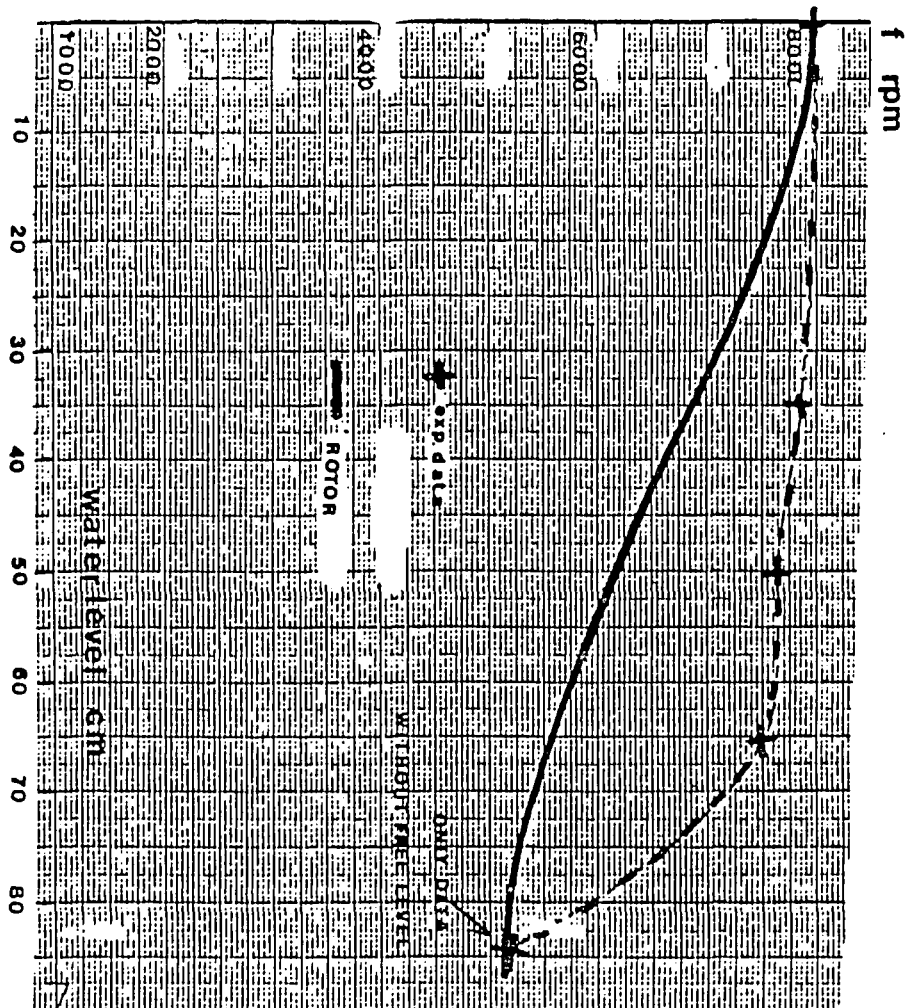
568



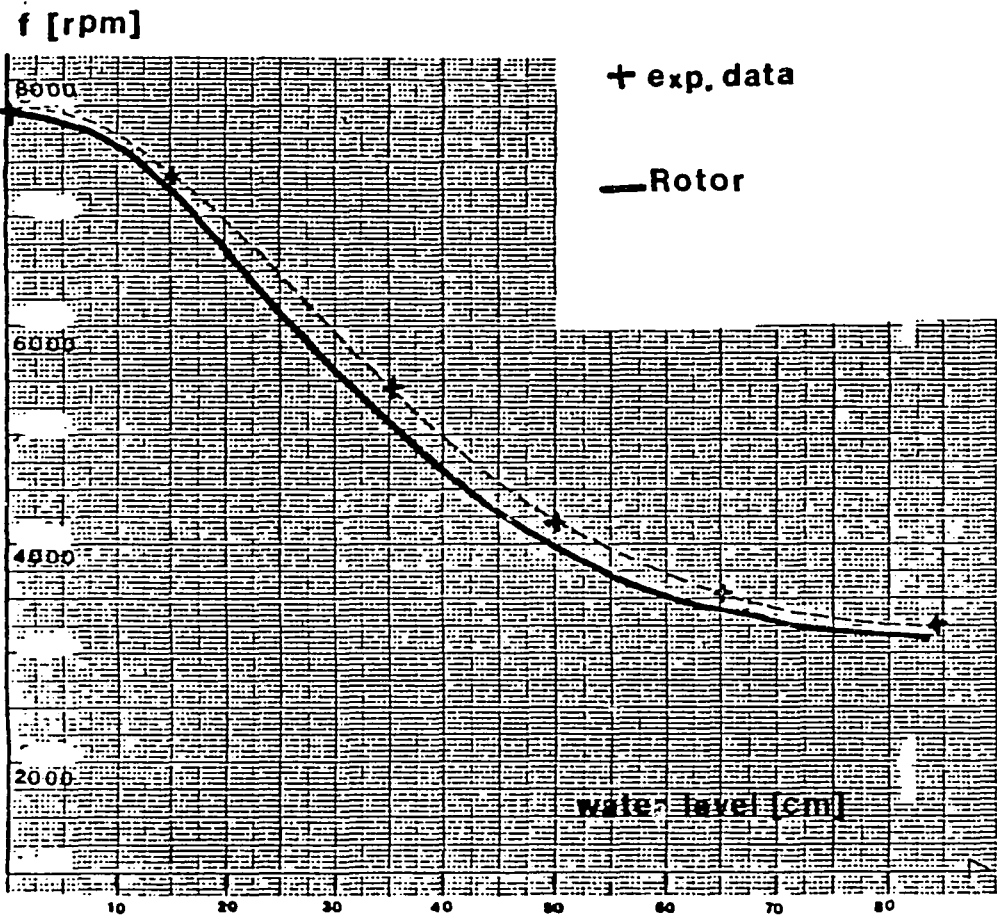
- n°1 : Palier modélisé sur 2 noeuds + labyrinthe représenté par son confinement (§III.6.B)
- n°2 : Palier modélisé sur 2 noeuds (§III.5)
- n°3 : Palier modélisé sur 2 noeuds + labyrinthe modélisé comme palier (§III.6.a)
- : Points de mesure.



ROTATING SHAFT: free level effect



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SHAFT AT REST

small shell , gap/radius: 0.14