

REFERENCE

IC/88/213



**INTERNATIONAL CENTRE FOR
THEORETICAL PHYSICS**

**NON-EXISTENCE OF STABLE HARMONIC MAPS
FROM SUFFICIENTLY PINCHED SIMPLY CONNECTED RIEMANNIAN MANIFOLDS**

Pan Yanglian



**INTERNATIONAL
ATOMIC ENERGY
AGENCY**



**UNITED NATIONS
EDUCATIONAL,
SCIENTIFIC
AND CULTURAL
ORGANIZATION**

1988 MIRAMARE-TRIESTE



International Atomic Energy Agency
and
United Nations Educational Scientific and Cultural Organization

INTERNATIONAL CENTRE FOR THEORETICAL PHYSICS

NONEXISTENCE OF STABLE HARMONIC MAPS
FROM SUFFICIENTLY PINCHED SIMPLY CONNECTED RIEMANNIAN MANIFOLDS *

Pan Yanglian **

International Centre for Theoretical Physics, Trieste, Italy.

ABSTRACT

It is proved that for $n \geq 3$ there exists a constant $\delta(n)$ with $\frac{1}{4} \leq \delta(n) < 1$ such that if M is a simply connected Riemannian manifold of dimension n with $\delta(n)$ -pinched curvatures then for every Riemannian manifold N every stable harmonic map $\phi: M \rightarrow N$ is constant. Together with Howard's result, it shows that a simply connected sufficiently pinched Riemannian manifold is weakly E-unstable.

MIRAMARE - TRIESTE

August 1988

* Submitted for publication.

** Permanent address: Institute of Mathematics, Fudan University, Shanghai, People's Republic of China.

§1 Introduction

A harmonic map is a critical point of the energy functional. A harmonic map is said to be stable if for any deformation vector field, its second variation is always non-negative. For simplicity, we give

Definition 1. Let M be a compact Riemannian manifold. M is said to be weakly E-unstable, if the following two conditions are fulfilled :

- (A) For any compact Riemannian manifold N there are no nonconstant stable harmonic maps from N to M ,
- (B) For any Riemannian manifold N there are no nonconstant stable harmonic maps from M to N .

Several classes of weakly E-unstable manifolds have been founded in recent years. A typical case is the Euclidean sphere S^n with $n \geq 3$. It is due to a combination of Xin's result for (B) [8] and Leung's result for (A) [5]. Taking S^n as model manifold, one might expect the weakly E-unstability for sufficiently pinched Riemannian manifolds, i.e., for compact Riemannian manifolds whose sectional curvatures are between the interval $[\delta K, K]$ with constants $K \geq 0$ and $1 \geq \delta > 0$. But Urakawa [7] shows that the identity maps of any non-simply connected manifold with positive constant curvatures are stable. So the condition of simply connectness is necessary. In 1985, Howard proved the following

Theorem 1 (Howard). Let $n \geq 3$. There is a number $\delta_1(n)$ with δ such that if M is simply connected with $\delta_1(n)$ -pinched curvatures then for every compact Riemannian manifold N every stable harmonic map $\phi: N \rightarrow M$ is constant.

It means for such a manifold the condition (A) is satisfied. This is a theorem determined only by the intrinsic geometry of the manifold.

In this paper, we establish the following

Theorem 2. Let $n \geq 3$. There is a number $\delta_2(n)$ with $\frac{1}{4} \leq \delta_2(n) < 1$ such that if M is a simply connected Riemannian manifold of dimension n with $\delta_2(n)$ -pinched curvatures then for any Riemannian manifold N every stable harmonic map $\phi: M \rightarrow N$ is constant.

Combining the above two theorems, letting $\delta(n) = \max(\delta_1(n), \delta_2(n))$ we obtain

Theorem 3. Let $n \geq 3$. There is a number $\delta(n)$ with $\frac{1}{4} \leq \delta(n) < 1$ such that if M is simply connected Riemannian manifold of dimension n with $\delta(n)$ -pinched curvatures then M is weakly E-unstable.

The proof of Theorem 2 goes in a way similar to that of Howard. That is, we make an integral average for the second variation formula over a continuous family of deformations and show the result is negative.

2 Second variation formula

Let M and N be Riemannian manifolds with dimension n and m respectively. M is compact without boundary, and ∇, ∇' represent the Riemannian connections of M and N respectively. Suppose that $\phi: M \rightarrow N$ is a harmonic map, $\phi_*: TM \rightarrow TN$ is the induced map, where TM and TN are the tangent bundle of M and N respectively. We also can consider ϕ_* as a ϕ^*TN valued 1-form $d\phi$, i.e., $d\phi(X) = \phi_*X$, for $X \in TM$. The induced bundle $\phi^*TN \rightarrow M$ possesses the induced Riemannian connection as follows

$$\tilde{\nabla}_X S = \nabla'_{\phi_*X} S, \quad (2.1)$$

where $X \in TM, S \in \Gamma(\phi^*TN)$.

Choose local fields of orthonormal frames $\{e_i\}$ and $\{e'_\alpha\}$ in M and N , respectively, and let $\{w_i\}$ and $\{w'_\alpha\}$ be the fields of dual forms. we shall make the following convention on the ranges of indices: $1 \leq i, j, k, \dots \leq n, 1 \leq \alpha, \beta, \dots \leq m$, and use the summation convention.

Let $W \in (\phi^*TN)$ be a deformation vector field, ϕ_t the one parametric family of maps generated by W , $\phi_0 = \phi$. It is well known that the second variation of the energy functional $E(\phi_t)$ is given by

$$\frac{d^2}{dt^2} E(\phi_t) \Big|_{t=0} = I(W, W) = - \int_M \langle \tilde{\nabla}^* \tilde{\nabla} W + R(\phi_* e_i, W) \phi_* e_i, W \rangle_N \Omega_M, \quad (2.2)$$

where $\tilde{\nabla}^* \tilde{\nabla}$ is the trace Laplacian with respect to $\tilde{\nabla}$, R is the curvature operator of N , and Ω_M is the volume element of M [6].

If we take $\phi_* V$, where $V \in TM$, as the deformation vector, then, using Weitzenböck formula, we can rewrite (2.2) as [5]

$$I(\phi_* V, \phi_* V) = \int_M \langle d\phi(\nabla_{e_i} \nabla_{e_i} V) - 2 \tilde{\nabla}_{e_i} (d\phi(\nabla_{e_i} V)) - \phi_* (\text{Ric}^M(V)), \phi_* V \rangle_N \Omega_M, \quad (2.3)$$

where Ric^M is the Ricci curvature operator of M , $\text{Ric}^M(e_j) = R_{ij} e_j$.

Under the map ϕ , suppose the pull back of w'_α is $\phi^*(w'_\alpha) = a_{\alpha i} w_i$. Then the energy density of ϕ is $e(\phi) = \frac{1}{2} \sum a_{\alpha i}^2$, the energy of ϕ is $E(\phi) = \frac{1}{2} \int_M \sum a_{\alpha i}^2 \Omega_M$, and the tension field of ϕ is $\tau = \sum_{\alpha, i} a_{\alpha i} e'_\alpha$, where $a_{\alpha i}$ is the covariant derivative of $a_{\alpha i}$. For harmonic map ϕ , $\tau = 0$.

Let $V = V_i e_i$, we compute the quantities in (2.3) as follows.

$\nabla_{e_i} V = V_{j,i} e_j$, $\nabla_{e_i} \nabla_{e_i} V = V_{j,i,i} e_j = (\Delta V_j) e_j$, where $V_{j,i}$, $V_{j,i,i}$ are covariant derivatives and Δ is the Laplacian of M .

$$\begin{aligned} d\phi(\nabla_{e_i} \nabla_{e_i} V) &= a_{\alpha j} (\Delta V_j) e'_\alpha \\ \tilde{\nabla}_{e_i} (d\phi(\nabla_{e_i} V)) &= \tilde{\nabla}_{e_i} (a_{\alpha j} V_{j,i} e'_\alpha) = (V_{j,i,i} a_{\alpha j} + V_{j,i} a_{\alpha j,i}) e'_\alpha \\ &= (a_{\alpha j} (\Delta V_j) + a_{\alpha j,i} V_{j,i}) e'_\alpha, \\ \phi_* (\text{Ric}^M(V)) &= a_{\alpha j} R_{ij} V_i e'_\alpha. \end{aligned}$$

Thus (2.3) becomes

$$I(\phi_* V, \phi_* V) = - \int_M (a_{\alpha j} a_{\alpha i} (\Delta V_i) V_j + 2 a_{\alpha j,i} a_{\alpha i} V_{j,i} V_i + a_{\alpha j} a_{\alpha i} R_{ij} V_i V_i) \Omega_M. \quad (2.4)$$

§ 3 Some estimates

In this section, we shall list some useful estimates obtained by Howard. Details can be found in [3]. Suppose V is the gradient of a smooth function f on M . For V we define a smooth field of linear endomorphisms of the tangent spaces to M by

$$\mathcal{Q}^Y(X) = \nabla_X V, \text{ for } X \in TM_x, x \in M. \quad (3.1)$$

For the gradient vector field of f a straightforward calculation shows

$$\langle \mathcal{Q}^Y X, Y \rangle_M = D^2 f(X, Y), \quad (3.2)$$

where $D^2 f$ is the Hessian of f . So \mathcal{Q}^Y is self-adjoint and has real eigenvalues $\lambda_1, \dots, \lambda_n$.

From now on we always suppose M is a compact simply connected with curvatures between δ and 1 , where $1 > \delta > \frac{1}{4}$. Due to Klingenberg's result, the injective radius of such a manifold is greater than π .

If $x \in M$, let $\rho_x(y)$ denote the geodesic distance at y from x . Define a function $f: \mathbb{R} \rightarrow \mathbb{R}$ by

$$f(t) = \begin{cases} -\cos t & |t| \leq \pi \\ 0 & |t| \geq \pi \end{cases} \quad (3.3)$$

Then f and f' are continuous. For any $x \in M$, let V^x be the vector field defined by

$$V^x = \nabla (f \circ \rho_x) = f'(\rho_x) \nabla \rho_x. \quad (3.4)$$

Then V^x is continuous and smooth off the locus defined by $\rho_x = \pi$, and $V^x = 0$ on the set defined by $\rho_x > \pi$.

Noting (3.2) and using the Hessian comparison theorem of Greene-Wu[2], Howard gives the following estimates for the eigenvalues of \mathcal{Q}^{V^x} (λ_i):

At any point y at a geodesic distance from $x = \rho$, with $\rho < \pi$, we have

$$\cos \rho < \lambda_i < \sqrt{\delta} \sin \rho \cot \sqrt{\delta} \rho, \quad 1 \leq i \leq n. \quad (3.5)$$

Set

$$\begin{aligned} \tilde{g}_1(t, \delta) &= \text{middle value of } \{ \cos t, 0, \sqrt{\delta} \sin t \cot \sqrt{\delta} t \}, \\ g_1(t, \delta) &= (\tilde{g}_1(t, \delta))^2, \quad 0 \leq t \leq \pi, \\ g_2(t, \delta) &= \max\{ \cos^2 t, \delta \sin^2 t \cot^2 \sqrt{\delta} t \}, \quad 0 \leq t \leq \pi, \\ g_1(t, \delta) &= g_2(t, \delta) = 0, \quad t > \pi. \end{aligned} \quad (3.6)$$

Then we have

$$g_1(\rho, \delta) \leq \lambda_i^2 \leq g_2(\rho, \delta). \quad (3.7)$$

The other comparison theorem needed is due to Bishop and Grittenden[1]. If $y \in M$ then let UM_y be the unit sphere in TM_y . Then, letting $\rho = \rho_y$, and $u \in UM_y$, we view (ρ, u) as polar coordinates on M near y in the obvious way. Let Ω_{UM_y} be the volume density on UM_y . Then on the open set $M \setminus \text{cut}(y)$, where $\text{cut}(y)$ is the cut locus of y ,

$$\sin^{n-1} \rho d\rho \Omega_{UM_y}(u) \leq \Omega_M \leq \left(\frac{\sin \sqrt{\delta} \rho}{\sqrt{\delta}}\right)^{n-1} d\rho \Omega_{UM_y}(u), \quad (3.8)$$

where the lower bound only holds up to π [1].

The following lemma can be found in [3].

Lemma 1. Let $Q : \mathbb{R}^n \rightarrow \mathbb{R}$ be a quadratic form, then

$$\int_{S^{n-1}} Q(u) \Omega_{S^{n-1}}(u) = \frac{1}{n} \text{vol}(S^{n-1}) \text{trace}(Q). \quad (3.9)$$

§ 4 Proof of the Theorem 2

Without loss of generality, we can assume the curvatures of M are between $[\delta, 1]$. Let x be a point of M and V^x the deformation vector field defined by (3.4). Write the second variation determined by V^x as $I_x(\psi_x V^x, \phi_x V^x)$, and let $\tilde{V}^x(y) = V_i^x(y) e_i$. Then from (2.4) we have

$$I_x(\phi_x V^x, \phi_x V^x) = \int_{M(y)} \left\{ a_{ij}(y) a_{kl}(y) (\Delta V_l^x(y)) V_j^x(y) + 2a_{ij}(y) a_{kl}(y) V_{ji}^x(y) V_l^x(y) + a_{ij}(y) a_{kl}(y) R_{ij}(y) V_l^x(y) V_k^x(y) \right\} \Omega_M(y). \quad (4.1)$$

In the following, we compute the integration

$$\int_{M(y)} I_x(\phi_x V^x, \phi_x V^x) \Omega_M(x)$$

and show there exists a certain constant $\delta_2(n)$ such that for $1 > \delta > \delta_2(n)$, the integration is negative. Hence, at least for a certain point x , $I_x(\phi_x V^x, \phi_x V^x)$ is negative.

First of all, we need to transform (4.1) into another suitable form. For simplicity, we omit the variable y , then no confusion is caused. We have

$$2a_{ij} a_{kl} V_{ji}^x V_l^x = 2(a_{ij} a_{kl} V_{ji}^x V_l^x)_{,i} - 2a_{ij} a_{kl} V_{ji}^x V_l^x - 2a_{ij} a_{kl} V_{ji}^x V_l^x - 2a_{ij} a_{kl} V_{ji}^x V_l^x. \quad (4.2)$$

Thus by divergence theorem (4.1) becomes

$$I_x(\phi V^x, \phi V^x) = \int_{M(y)} \{ a_{\alpha j} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x + 2a_{\alpha j} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x + 2a_{\alpha j} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x - a_{\alpha j} a_{\alpha l} R_{ij} V_{j\alpha}^x V_{l\alpha}^x \} \Omega_M(y) . \quad (4.3)$$

Noting $a_{\alpha j} = a_{\alpha j}$ and $V_{j\alpha}^x = V_{j\alpha}^x$, from

$$a_{\alpha j} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x = (a_{\alpha j} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x)_{, \alpha} - a_{\alpha j \alpha} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x - a_{\alpha j} a_{\alpha l \alpha} V_{j\alpha}^x V_{l\alpha}^x - a_{\alpha j} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x ,$$

we have

$$2 \int_{M(y)} a_{\alpha j} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x \Omega_M(y) = - \int_{M(y)} (a_{\alpha j} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x + a_{\alpha j} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x) \Omega_M(y) . \quad (4.4)$$

Using Ricci identities and noting $R_{ij} = -R_{i,jk}k$, we have

$$\begin{aligned} a_{\alpha j} a_{\alpha l} V_{j\alpha}^x V_{l\alpha}^x &= a_{\alpha j} a_{\alpha l} (V_{,ij}^x - V_k^x R_{kij}) V_{l\alpha}^x \\ &= a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x + a_{\alpha j} a_{\alpha l} R_{ij} V_k^x V_{l\alpha}^x , \end{aligned} \quad (4.5)$$

Noting $a_{\alpha j} = 0$ due to the harmonicity of ϕ , we have

$$\begin{aligned} a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x &= (a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x)_{, j} - a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x - a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x - a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x \\ &= (a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x)_{, j} - a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x - a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x . \end{aligned} \quad (4.6)$$

From (4.4) - (4.6), it follows that

$$I_x(\phi V^x, \phi V^x) = \int_{M(y)} \{ -a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x - 2a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x - a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x + 2a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x \} \Omega_M(y) . \quad (4.7)$$

Now since

$$a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x = e(\phi)_{, \alpha} V_{,ij}^x V_{l\alpha}^x = (e(\phi) V_{,ij}^x V_{l\alpha}^x)_{, \alpha} - e(\phi) V_{,ij}^x V_{l\alpha}^x - e(\phi) V_{,ij}^x V_{l\alpha}^x ,$$

(4.7) becomes

$$I_x(\phi V^x, \phi V^x) = \int_{M(y)} \{ e(\phi) V_{,ij}^x V_{l\alpha}^x + e(\phi) V_{,ij}^x V_{l\alpha}^x - 2a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x - a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x + 2a_{\alpha j} a_{\alpha l} V_{,ij}^x V_{l\alpha}^x \} \Omega_M(y) \quad (4.8)$$

Since v^x is the gradient vector of the function $f \circ \rho_x$, at the

point y with $\rho_x(y) = \rho$ and $\rho < \pi$, we have $v^x(y) = \sin \rho \frac{\partial}{\partial \rho}$ and

since $\nabla_{\frac{\partial}{\partial \rho}} \frac{\partial}{\partial \rho} = 0$,

$$\nabla_x v^x(y) = \cos \rho v^x(y) . \quad (4.9)$$

So at the point y , with respect to the frame $\{e_i\}$, we have from

(4.9)

$$V_{,i}^x V_{,i}^x = \cos \rho V_j^x . \quad (4.10)$$

Differentiating (4.10), we get

$$V_{,i}^x V_{,j}^x + V_{,i}^x V_{,j}^x = -V_{,i}^x V_{,j}^x + \cos \rho V_{,j}^x . \quad (4.11)$$

It follows

$$V_{,j}^x V_{,i}^x = -V_{,i}^x V_{,j}^x + \cos \rho V_{,j}^x - V_{,i}^x V_{,j}^x . \quad (4.12)$$

Now, using Ricci identities, from (4.12) we have

$$V_{ji} V_i^x = -V_i^x V_j^x + \cos \rho V_{ji}^x - V_{il}^x V_{jl}^x - V_k^x V_i^x R_{kj} e_i \quad (4.13)$$

Summing on the indices j and l , we get

$$V_{li} V_i^x = -\sin^2 \rho + \cos \rho V_{li}^x - V_{il}^x V_{li}^x - R_{ik} V_i^x V_k^x \quad (4.14)$$

Substituting (4.13) and (4.14) into (4.8), (4.8) becomes

$$\begin{aligned} I_x(\phi_s V^x, \phi_s V^x) = \int_{M(y)} \{ & -e(\phi) \sin^2 \rho + e(\phi) \cos \rho V_{li}^x - e(\phi) V_{li}^x V_{li}^x \\ & - e(\phi) R_{ik} V_i^x V_k^x + e(\phi) V_{li}^x V_{li}^x - 2a_{ij} a_{kl} V_{li}^x V_{kj}^x + a_{ij} a_{kl} V_i^x V_j^x \\ & - \cos \rho a_{ij} a_{kl} V_{ij}^x + 3a_{ij} a_{kl} V_{ij}^x V_{li}^x + a_{ij} a_{kl} V_k^x V_l^x R_{kj} e_i \} \Omega_M(y) \end{aligned} \quad (4.15)$$

At the point y , we take the unit eigenvectors of \mathcal{K}^{V^x} as $\{e_i\}$, and we have

$$V_{ij}^x = \lambda_i \delta_{ij}, \quad i, j, = 1, 2, \dots, n, \quad (4.16)$$

where λ_i is the eigenvalue of \mathcal{K}^{V^x} .

Now we can estimate each term of the integration (4.15) just in the same way as in [3]. Using (3.5), (3.7), (3.8) and (4.16) we get

$$\int_{M(x)} \int_{M(y)} (-e(\phi) \sin^2 \rho) \Omega_M(y) \Omega_M(x) \leq \text{vol}(S^{n-1}) \int_{M(y)} e(\phi) \int_0^{\pi} (-\sin^2 \rho) d\rho \Omega_M(y), \quad (4.17)$$

$$\begin{aligned} \int_{M(x)} \int_{M(y)} e(\phi) \cos \rho V_{li}^x \Omega_M(y) \Omega_M(x) \leq \text{vol}(S^{n-1}) \int_{M(y)} e(\phi) \left\{ \int_0^{\frac{\pi}{2}} n \sqrt{\delta} \sin \rho \cos \rho \cot \rho \left(\frac{\sin \sqrt{\delta} \rho}{\sqrt{\delta}} \right)^{n-1} d\rho \right. \\ \left. + \int_{\frac{\pi}{2}}^{\pi} n \cos^2 \rho \left(\frac{\sin \sqrt{\delta} \rho}{\sqrt{\delta}} \right)^{n-1} d\rho \right\} \Omega_M(y), \end{aligned} \quad (4.18)$$

$$\int_{M(x)} \int_{M(y)} (-e(\phi) V_{li}^x V_{li}^x) \Omega_M(y) \Omega_M(x) \leq \text{vol}(S^{n-1}) \int_{M(y)} e(\phi) \int_0^{\pi} (-ng_1(\rho, \delta) \sin^2 \rho) d\rho \Omega_M(y) \quad (4.19)$$

$$\int_{M(x)} \int_{M(y)} e(\phi) V_{li}^x V_{li}^x \Omega_M(y) \Omega_M(x) \leq \text{vol}(S^{n-1}) \int_{M(y)} e(\phi) \int_0^{\pi} ng_2(\rho, \delta) \left(\frac{\sin \sqrt{\delta} \rho}{\sqrt{\delta}} \right)^{n-1} d\rho \Omega_M(y), \quad (4.20)$$

$$\begin{aligned} \int_{M(x)} \int_{M(y)} (-\cos \rho a_{ij} a_{kl} V_{ij}^x) \Omega_M(y) \Omega_M(x) \leq \text{vol}(S^{n-1}) \int_{M(y)} e(\phi) \left\{ \int_0^{\frac{\pi}{2}} -2 \cos^2 \rho \sin^2 \rho d\rho \right. \\ \left. + \int_{\frac{\pi}{2}}^{\pi} -2 \sqrt{\delta} \sin \rho \cos \rho \cot \rho \left(\frac{\sin \sqrt{\delta} \rho}{\sqrt{\delta}} \right)^{n-1} d\rho \right. \\ \left. + \int_{\frac{\pi}{2}}^{\pi} -2 \sqrt{\delta} \sin^2 \rho \cos \rho \cot \rho \left(\frac{\sin \sqrt{\delta} \rho}{\sqrt{\delta}} \right)^{n-1} d\rho \right\} \Omega_M(y), \end{aligned} \quad (4.21)$$

and

$$\int_{M(x)} \int_{M(y)} a_{ij} a_{kl} V_{li}^x V_{ij}^x \Omega_M(y) \Omega_M(x) \leq \text{vol}(S^{n-1}) \int_{M(y)} e(\phi) \int_0^{\pi} 2g_2(\rho, \delta) \left(\frac{\sin \sqrt{\delta} \rho}{\sqrt{\delta}} \right)^{n-1} d\rho \Omega_M(y). \quad (4.22)$$

Using Lemma 1 and (3.8), we have

$$\begin{aligned}
\int_{M(y)} \int_{M(x)} a_{\alpha_i} a_{\beta_j} V_i^x V_j^x \Omega_M(x) \Omega_M(y) &\leq \int_{M(y)} \left(\int_{M(x)} a_{\alpha_i} a_{\beta_j} V_i^x V_j^x \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho \Omega_M(x) \right) \Omega_M(y) \\
&= \int_{M(y)} \left(\int_0^{\sqrt{\xi}} \frac{1}{\sqrt{\xi}} \text{vol}(S^{n-1}) \sum_{\alpha_i} (a_{\alpha_i}^2) \cdot \sin^2 \rho \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho \right) \Omega_M(y) \\
&= \text{vol}(S^{n-1}) \int_{M(y)} e(\phi) \int_0^{\sqrt{\xi}} \frac{2}{\sqrt{\xi}} \sin^2 \rho \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho \Omega_M(y) .
\end{aligned} \tag{4.23}$$

From the assumption on the curvatures of M , it follows

$$\begin{aligned}
-\int_{M(y)} \int_{M(x)} e(\phi) R_{ik} V_i^x V_k^x \Omega_M(x) \Omega_M(y) &\leq -\int_{M(y)} \int_{M(x)} e(\phi) \cdot (n-1) \delta \sin^2 \rho \Omega_M(x) \Omega_M(y) \\
&< -\text{vol}(S^{n-1}) \int_{M(y)} e(\phi) \int_0^{\sqrt{\xi}} (n-1) \delta \sin^2 \rho d\rho \Omega_M(y) .
\end{aligned} \tag{4.24}$$

and using Lemma 1 and the assumption on curvatures, we have

$$\int_{M(y)} \int_{M(x)} a_{\alpha_i} a_{\beta_j} V_i^x V_k^x R_{kj} \Omega_M(x) \Omega_M(y) \leq \text{vol}(S^{n-1}) \int_{M(y)} e(\phi) \int_0^{\sqrt{\xi}} \frac{2(n-1)}{\sqrt{\xi}} \sin^2 \rho \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho \Omega_M(y) . \tag{4.25}$$

It is also easy to get the following estimate

$$\begin{aligned}
-2 \int_{M(y)} \int_{M(x)} a_{\alpha_i} a_{\beta_j} V_i^x V_j^x \Omega_M(x) \Omega_M(y) &\leq \text{vol}(S^{n-1}) \int_{M(y)} e(\phi) \left\{ \int_0^{\frac{\pi}{2}} -4n \cos^2 \rho \sin^{n-1} \rho d\rho \right. \\
&\quad \left. + \int_{\frac{\pi}{2}}^{\pi} -4n \delta \sin^2 \rho \cot^2 \rho d\rho + \int_{\frac{\pi}{2}}^{\pi} -4n \sqrt{\xi} \sin^2 \rho \cos \rho \cot^2 \rho d\rho \right\} \Omega_M(y) .
\end{aligned} \tag{4.26}$$

By means of the above estimates (4.17)-(4.26) and Fubini theorem, we get

$$\int_{M(y)} I_x(\phi, V^x, \phi, V^x) \Omega_M(x) \leq \text{vol}(S^{n-1}) E(\phi) F(n, \xi) , \tag{4.27}$$

where $F(n, \xi)$ is a continuous function of ξ given by

$$\begin{aligned}
F(n, \xi) &= \int_0^{\sqrt{\xi}} \sin^2 \rho d\rho + \int_0^{\frac{\pi}{2}} n \delta \sin^2 \rho \cos \rho \cot^2 \rho \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho \\
&\quad + \int_0^{\frac{\pi}{2}} n \cos^2 \rho \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho - \int_0^{\pi} n g_1(\rho, \xi) \sin^{n-1} \rho d\rho \\
&\quad + \int_0^{\frac{\pi}{2}} n g_2(\rho, \xi) \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho + \int_0^{\pi} \frac{2}{\sqrt{\xi}} \sin^2 \rho \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho \\
&\quad - \int_0^{\frac{\pi}{2}} (n-1) \delta \sin^{n-1} \rho d\rho - \int_0^{\frac{\pi}{2}} 2 \cos^2 \rho \sin^{n-1} \rho d\rho \\
&\quad - \int_0^{\frac{\pi}{2}} 2 \sin \rho \cos \rho \cot^2 \rho \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho - \int_0^{\frac{\pi}{2}} 2 \sqrt{\xi} \sin^2 \rho \cos \rho \cot^2 \rho d\rho \\
&\quad + \int_0^{\frac{\pi}{2}} 6 g_2(\rho, \xi) \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho + \int_0^{\pi} \frac{2(n-1)}{\sqrt{\xi}} \sin^2 \rho \left(\frac{\sin \sqrt{\xi} \rho}{\sqrt{\xi}} \right)^{n-1} d\rho \\
&\quad - \int_0^{\frac{\pi}{2}} 4n \cos^2 \rho \sin^{n-1} \rho d\rho - \int_0^{\frac{\pi}{2}} 4n \delta \sin^2 \rho \cot^2 \rho d\rho \\
&\quad - \int_0^{\frac{\pi}{2}} 4n \sqrt{\xi} \sin^2 \rho \cos \rho \cot^2 \rho d\rho .
\end{aligned} \tag{4.28}$$

If $\xi = 1$, then $g_1(\rho, \xi) = g_2(\rho, \xi) = \cos^2 \rho$ and consequently

$$F(n, 1) = -2(n-2) \int_0^{\frac{\pi}{2}} \cos^2 \rho \sin^{n-1} \rho d\rho , \tag{4.29}$$

which is negative when $n \geq 3$.

Set

$$\bar{\epsilon}_2(n) = \inf \{ \delta : \delta > \frac{1}{4} \text{ and } F(n, \delta) \leq 0 \}. \quad (4.30)$$

First note that $F(n, 1) < 0$ and so $\frac{1}{4} \leq \bar{\epsilon}_2(n) < 1$. Then (4.27) shows that for any nonconstant harmonic map $\phi: M \rightarrow N$ there is a certain vector field V^x such that the second variation for V^x is negative. If V^x were smooth we might conclude that ϕ is stable iff ϕ is constant. So the only problem is that the vector field V^x is not smooth. But this difficulty can be overcome just in the same way as in [3]. So the proof is complete.

ACKNOWLEDGMENTS

The author would like to thank Professor Abdus Salam, the International Atomic Energy Agency and UNESCO for hospitality at the International Centre for Theoretical Physics, Trieste. He would also like to thank Professor J. Eells for his helpful conversation. This work was supported by the Science Fund of the Chinese Academy of Sciences.

REFERENCES

- [1] R.L.Bishop and R.J.Grittenden, Geometry of manifolds, Academic Press , New York, 1964.
- [2] R.E.Greene and H.Wu, Function theory on manifolds which possess a pole, Lectures in Math. 699 Springer Berlin, 1979.
- [3] R.Howard, The nonexistence of stable submanifolds, varifolds, and harmonic maps in sufficiently pinched simply connected Riemannian manifolds, Michigan Math.J. 32(1985), 321-334.
- [4] P.F.Leung, On the stability of harmonic mappings, 122-129, Lectures in Math. 949, Springer, New York, 1982.
- [5] Y.L.Pan (Pan Yanglian), Some nonexistence theorems on stable harmonic mappings, Chin. Ann. of Math. 3(4) 1982, 315-318.
- [6] R.T.Smith, The second variation formula for harmonic mappings, Proc. Amer. Math. Soc. 47(1975).
- [7] H.Urakawa, Stability of harmonic maps and eigenvalues of the Laplacian , Trans. Amer.Math. Soc. 301(1987) 557-589.
- [8] Y.L.Xin, Some results on stable harmonic maps, Duke Math J. 47(3), 1980.

Stampato in proprio nella tipografia
del Centro Internazionale di Fisica Teorica