

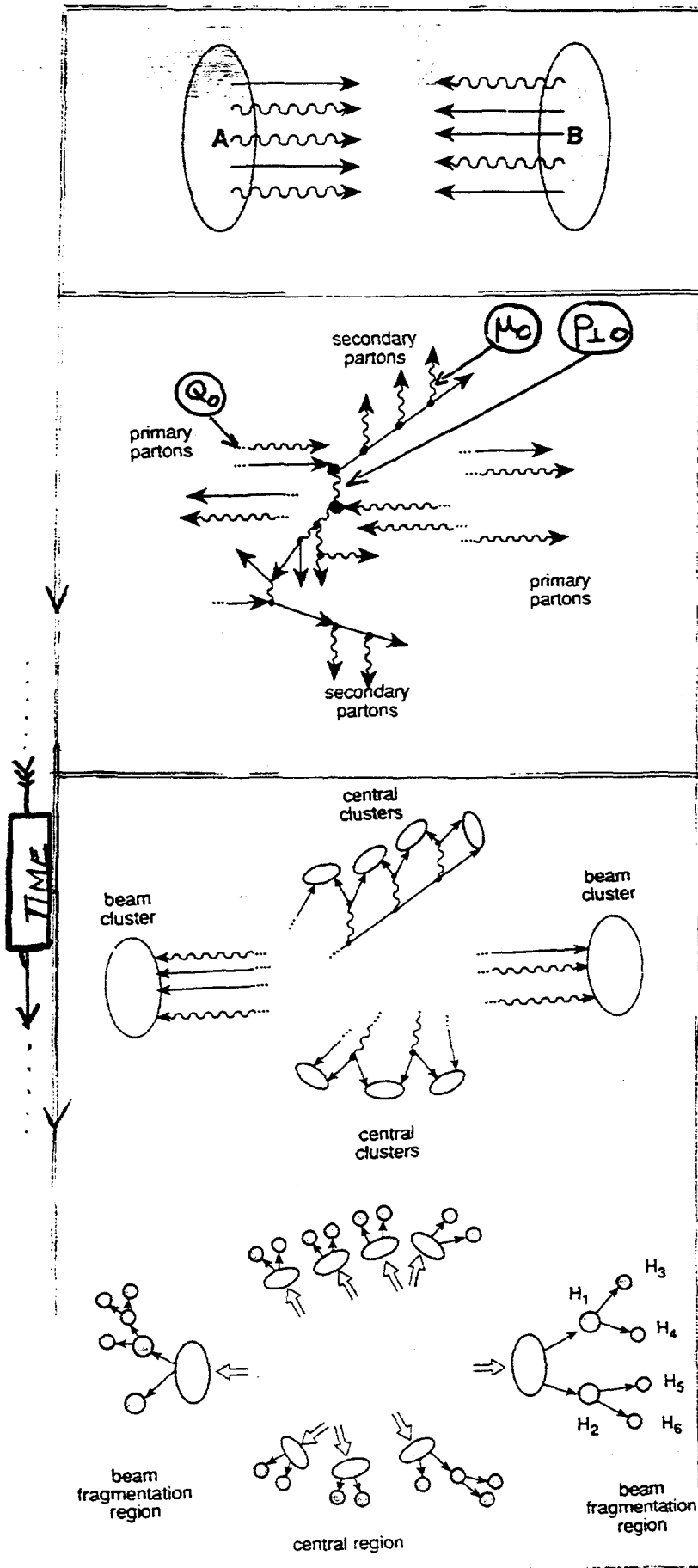
QCD space-time description
of high-density particle matter
in high-energy collisions

... from clean e^+e^- to dirty nuclei....

Klaus Kinder-Geiger

Kay Kay Gee
BNL

- Why at all?
- How to?
- Does it make sense?
- What then?



• INITIAL STATE

SCALE: Q_0^2

• PARTON CASCADE DEVELOPMENT

SCALE: $[Q_0^2, \mu_0^2]$

• PARTON RECOMBINATION (preconfinement)

SCALE: $[\mu_0^2, \Lambda^2]$

• CLUSTER FRAGMENTATION (confinement)

SCALE: Λ^2

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Closed set of 'balance' equations

• partons:

$$k \cdot \frac{\partial}{\partial \tau} F_P = \left[\begin{array}{c} \text{parton-scattering} \\ \text{parton-fusion} \\ - \text{parton-cluster} \end{array} \right] + \left[\begin{array}{c} \text{parton-branching} \end{array} \right]$$

$$k^2 \frac{\partial}{\partial k^2} F_P = \left[\text{parton-branching} \right]$$

• pre-hadronic clusters

$$k \cdot \frac{\partial}{\partial \tau} F_C = \left[\begin{array}{c} \text{parton-cluster} \\ \text{cluster-hadron} \end{array} \right]$$

$$k^2 \frac{\partial}{\partial k^2} F_C \approx 0$$

• hadrons

$$k \cdot \frac{\partial}{\partial \tau} F_H = \left[\begin{array}{c} \text{cluster-hadron} \\ \text{hadron-decay} \end{array} \right]$$

$$k^2 \frac{\partial}{\partial k^2} F_H \approx 0$$

Microscopic dynamics \rightarrow macroscopic observables

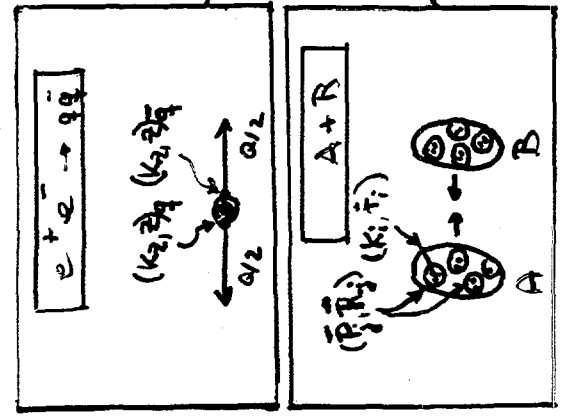
Lorentz-invariant 1-particle distribution ("Wigner function"):

$$F_a(\tau, k) d^4\tau d^4k = \begin{cases} \# \text{ of particles } a = g, q, \text{ cluster, } \pi, k, p, \dots \\ \text{at time } \tau \text{ within } d\tau, \text{ position in } d^3\vec{r} \text{ around } \vec{r}, \\ \text{energy (momentum) in } dE(d\vec{k}) \text{ around } E(\vec{k}). \end{cases}$$

$$\tau^M = (t, \vec{r}); \quad k^M = (E, \vec{k}); \quad k^2 = \begin{cases} = \frac{m^2}{m^2} & \text{on shell} \\ > \frac{m^2}{m^2} & \text{space-like} \\ < \frac{m^2}{m^2} & \text{time-like} \end{cases}$$

Initial value problem: specify initial state at time $t = t_0 = 0$:

$$F_a^{(0)}(\tau, k) \equiv F_a(t=0, \vec{r} | E, \vec{k}) = \frac{d N_a(t=0)}{d^3\tau d^3k dE}$$



$$= \mathcal{P}_a(E, \vec{k}) \bullet \mathcal{R}_a(\vec{r})$$

$$\sum_{i=q\bar{q}} \left[\delta(E_i - \frac{Q}{2}) \delta(k_{2i} \pm \frac{Q}{2}) \delta^2(k_{1i}) \right] \left[e^{-z_i^2 Q^2} \delta^2(\vec{r}_i) \right]$$

$$\sum_{j=1}^{A+B} \sum_{i=1}^{N_j} \left[f_a^j \left(\frac{k_{ij}}{P_j}, \frac{Q_j^0}{P_j} \right) g(k_{ij}) \right] \left[e^{-z_j^2 (\vec{r}_i - \vec{R}_j)^2} h^{A(0)}(\vec{r}_j, \vec{P}_j) \right]$$

e.g. \uparrow

e.g. \uparrow

For $t > t_0 = 0$, solve coupled evolutional & collisional cascade in small timesteps

$$\frac{\partial}{\partial \ln k^2} F_a(\tau, k) = E_a(F) \otimes F_a$$

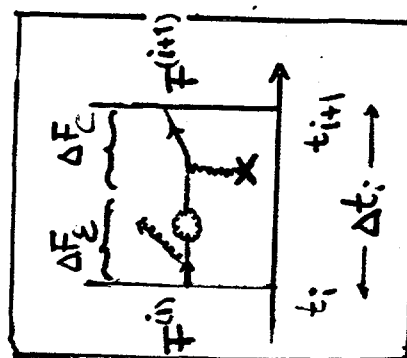
$$k \cdot \frac{\partial}{\partial \tau} F_a(\tau, k) = C_a(F) \otimes F_a$$

Iterative 2-step procedure: trace particle history $t_i = t_{i-1} + \Delta t_i$ from t_0 with $F_c^{(0)}(t_0)$:

$$F^{(i+1)} = F^{(i)} + \Delta F_E^{(i)} + \Delta F_C^{(i)}$$

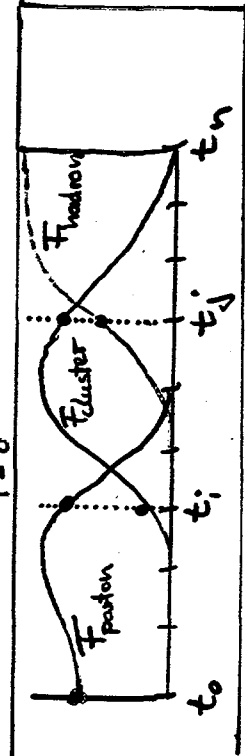
$$\Delta F_E^{(i)} = \Delta \ln k^2 E(F^{(i)}) \otimes F^{(i)} \cdot \Theta(\Delta t_i - \Delta \ln k^2)$$

$$\Delta F_C^{(i)} = \frac{\Delta t_i}{E} C(F^{(i)}) \otimes F^{(i)} \cdot \Theta(\Delta t_i - \Delta |\dot{\tau}|)$$



At any finite $t = t_n > t_0$ microscopic particle history is contained in

$$F_a(t, \hat{\tau}(\vec{E}, \vec{k})) = \sum_{i=0}^n F_a^{(i)}(t_i, \tau(\vec{E}, \vec{k}))$$



From statistical physics: extract from F macroscopic, spacetime-dep. quantities:

particle current: $n_a^M(t, \vec{t}) = d_a \int \frac{d^4 k}{(2\pi)^4} k^M F_a(t, \vec{t} | \epsilon, \vec{k})$

energy mom. tensor: $T_a^{MN}(t, \vec{t}) = d_a \int \frac{d^4 k}{(2\pi)^4} k^M k^N F_a(t, \vec{t} | \epsilon, \vec{k})$

entropy current: $S_a^M(t, \vec{t}) = d_a \int \frac{d^4 k}{(2\pi)^4} k^M \left\{ F_a \ln(F_a) - \theta_a (1 + \theta_a F_a) \ln(1 + \theta_a F_a) \right\}$

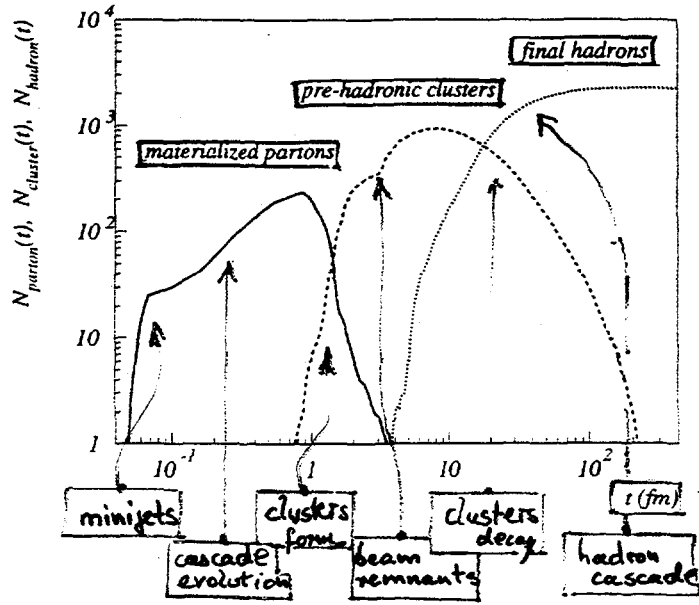
Define local flow velocity

$$u^M(t, \vec{t}) = \frac{n^M}{n}$$

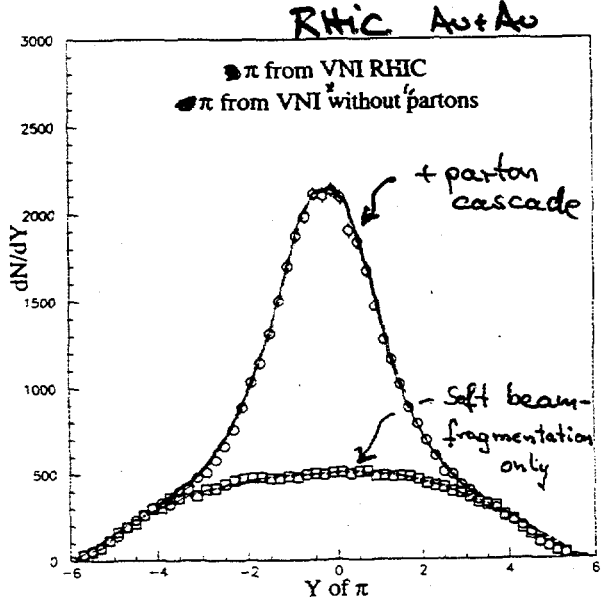
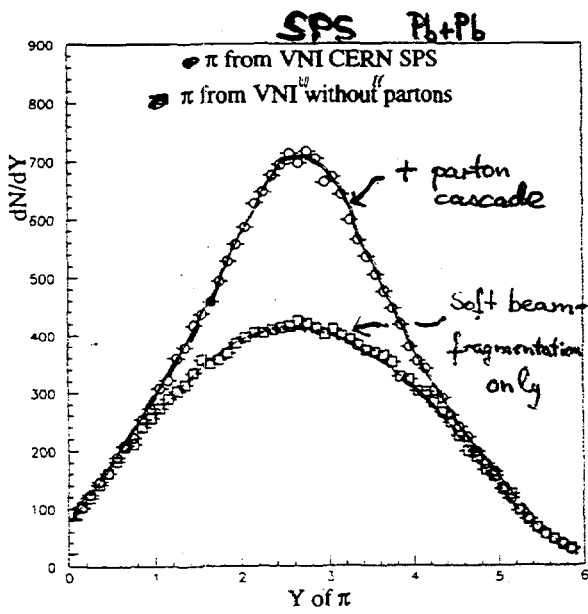
Get Lorentz-invariant, scalar density profiles

- number density: $S_a(t, \vec{t}) = n_{\mu a} n_a^M$ $\rightsquigarrow \frac{dN}{d^3 k_{\perp}}$
- energy density: $E_a(t, \vec{t}) = -\frac{1}{2} T_a^{\mu\nu} (g_{\mu\nu} - u_{\mu a} u_{\nu a})$ $\rightsquigarrow \frac{dE_{\perp}}{d^3 y}$
- pressure: $P_a(t, \vec{t}) = T_a^{\mu\nu} u_{\mu a} u_{\nu a}$ $\rightsquigarrow \frac{d\langle k_{\perp}^2 \rangle}{d^3 z}$
- entropy density: $S_a(t, \vec{t}) = S_a^M u_{\mu a}$ $\rightsquigarrow \frac{dS}{d^3 V_{\perp}}$
- eff. temperature: $T_a(t, \vec{t}) = \frac{E_{\text{eff}} P_a}{S_a}$

Time evolution of Pb+Pb @ CERN SPS



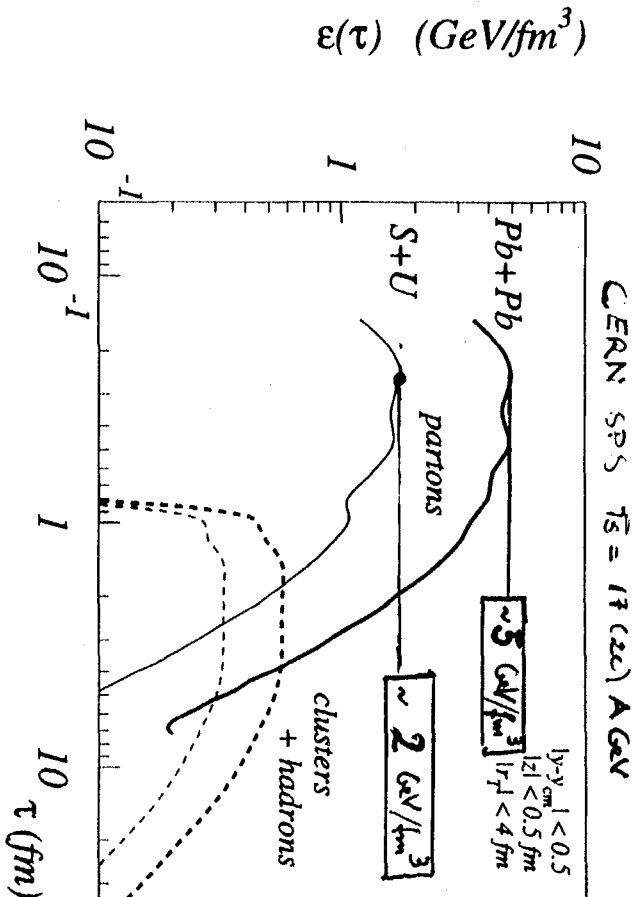
"Partonic contribution" in final hadron yield



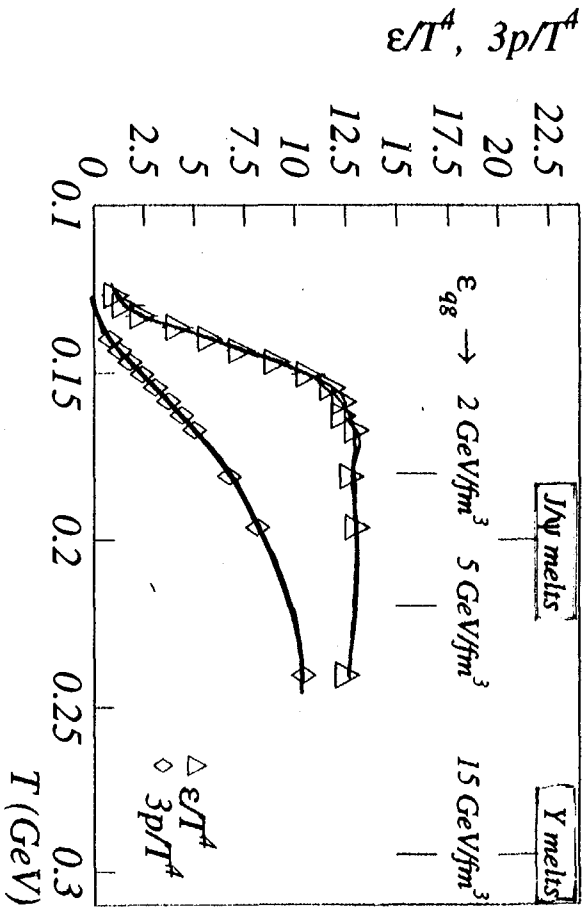
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Energy density: Parton Cascade vs. Lattice QCD

B. Mueller, KKs '97



VNI-32
 KKG, '97

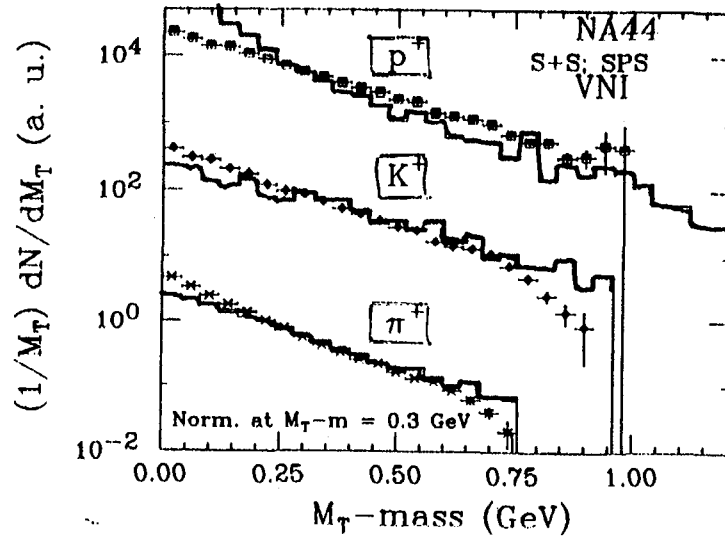
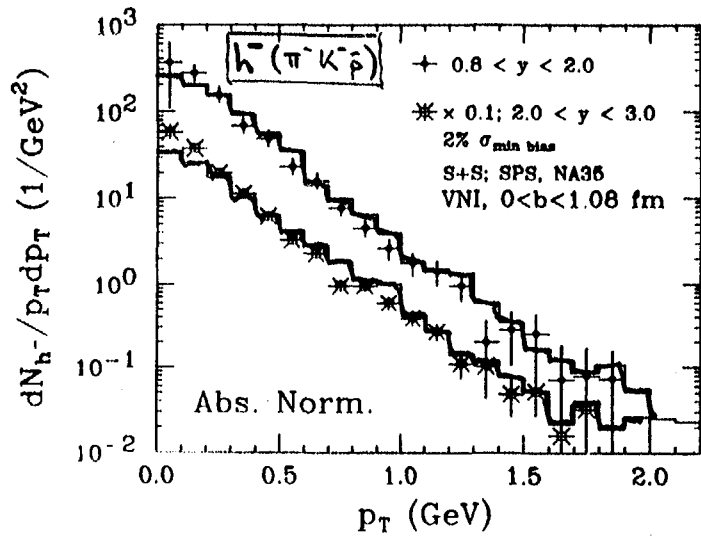
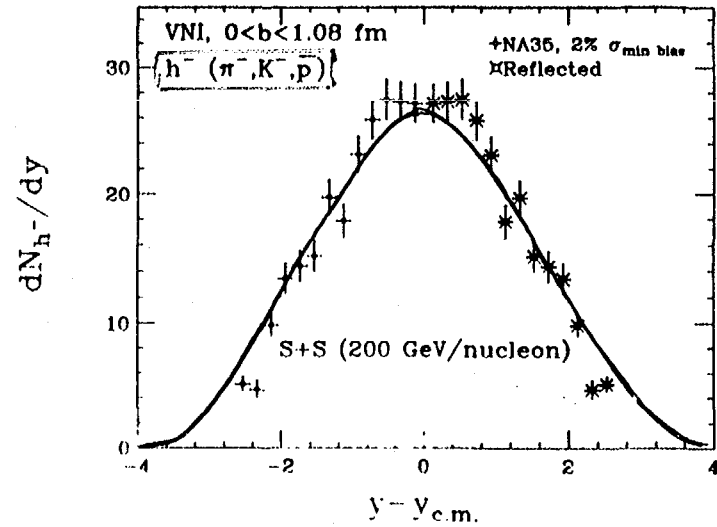
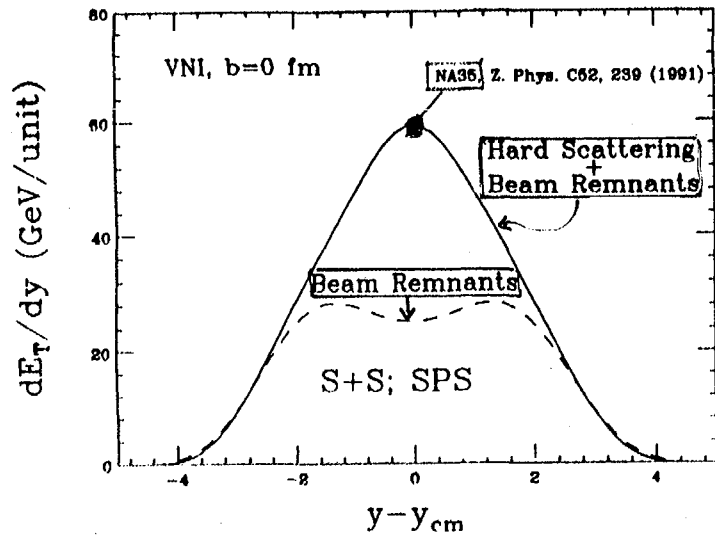


Karsch & Satz '96

MilC Coll.,
 Baur et al. '96

CERN SPS S+S @ 200 A GeV

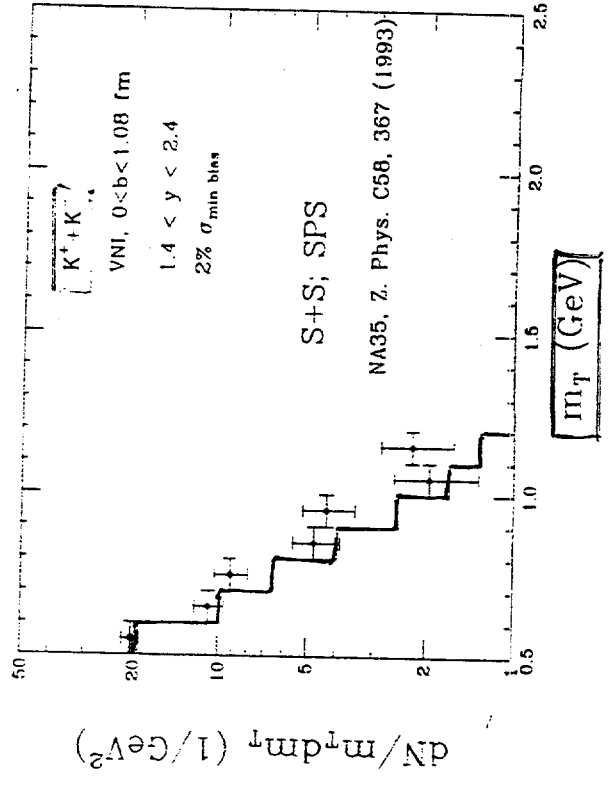
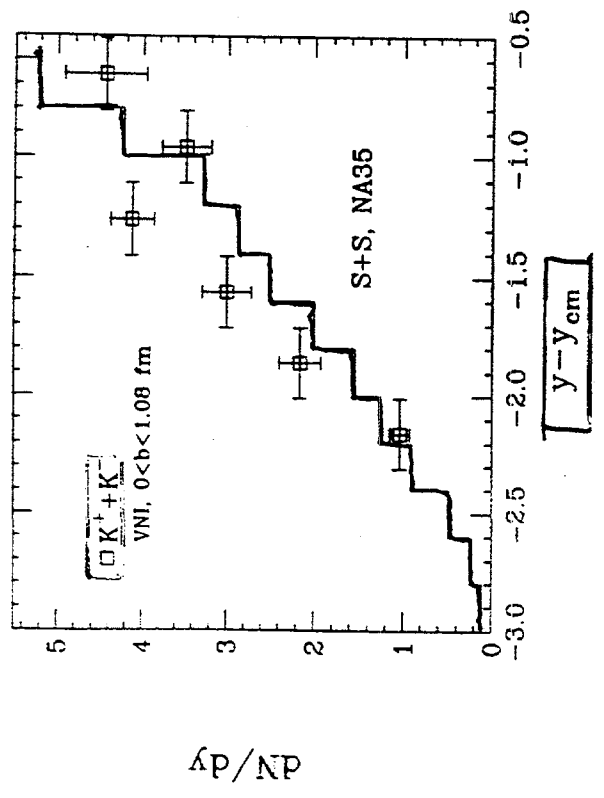
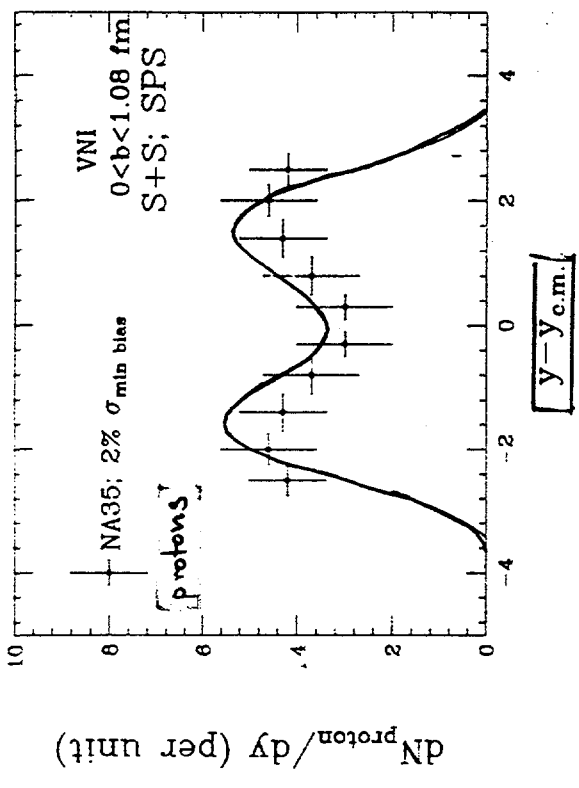
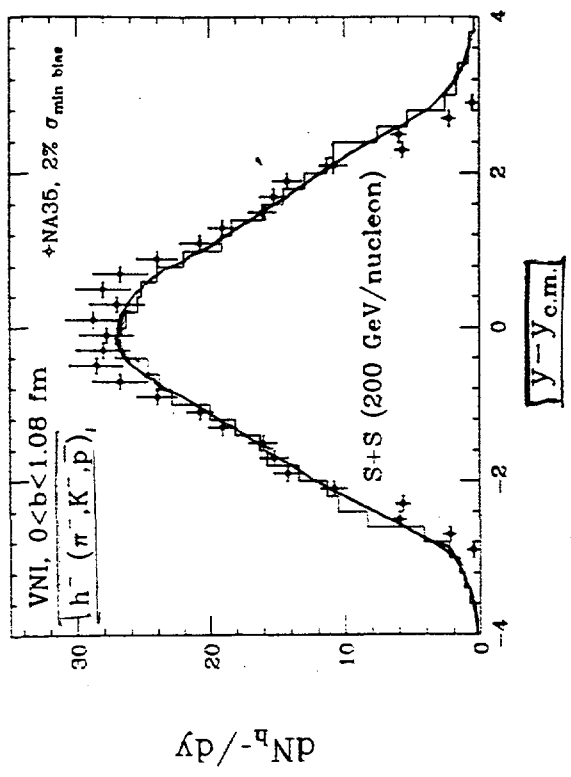
Srivastava
& KUG 97



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CERN SPS S+S @ 200 AGeV

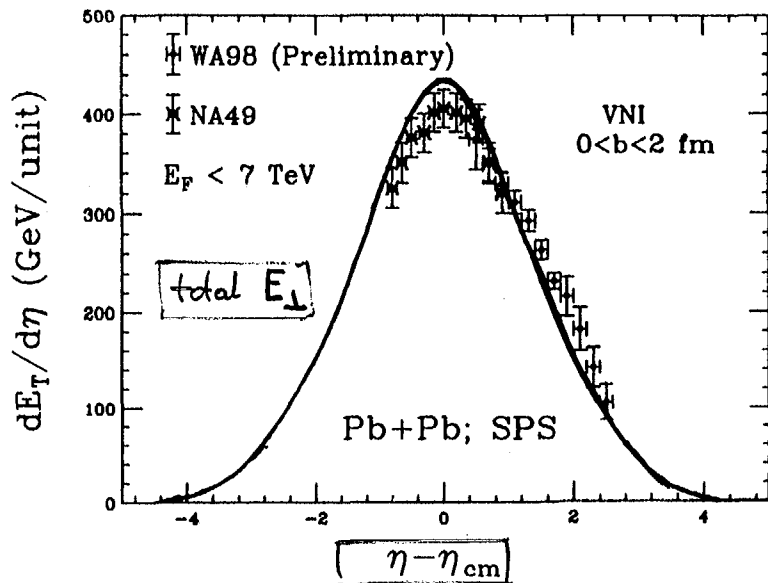
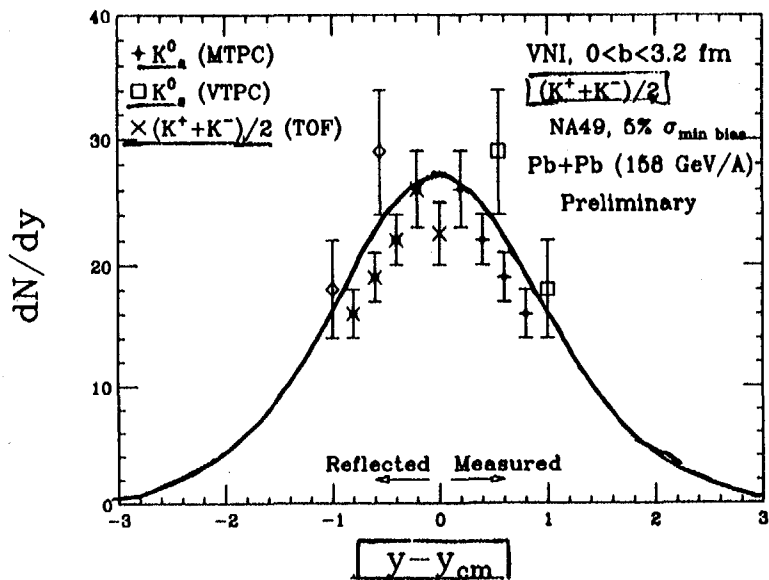
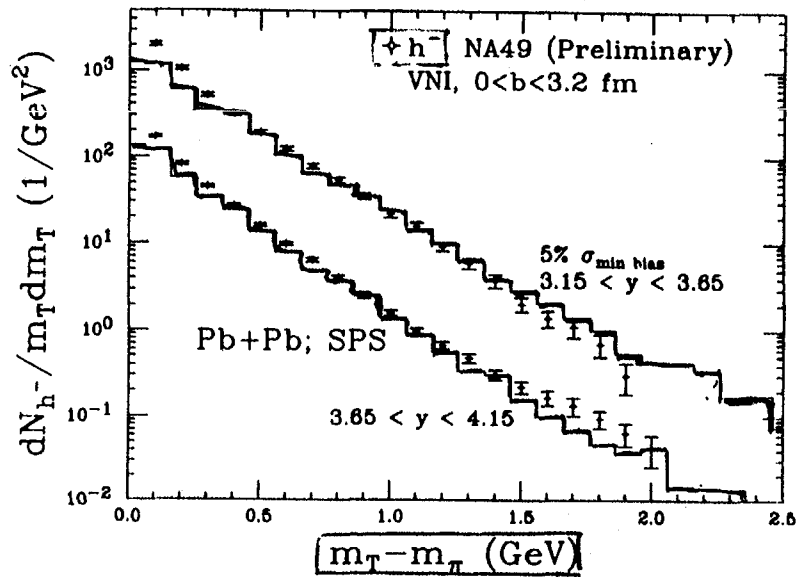
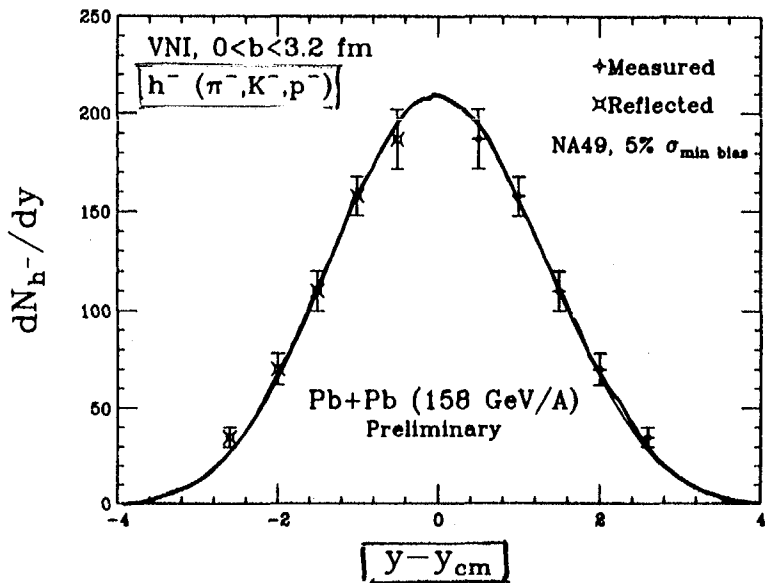
Srivastava '97
2 KKG

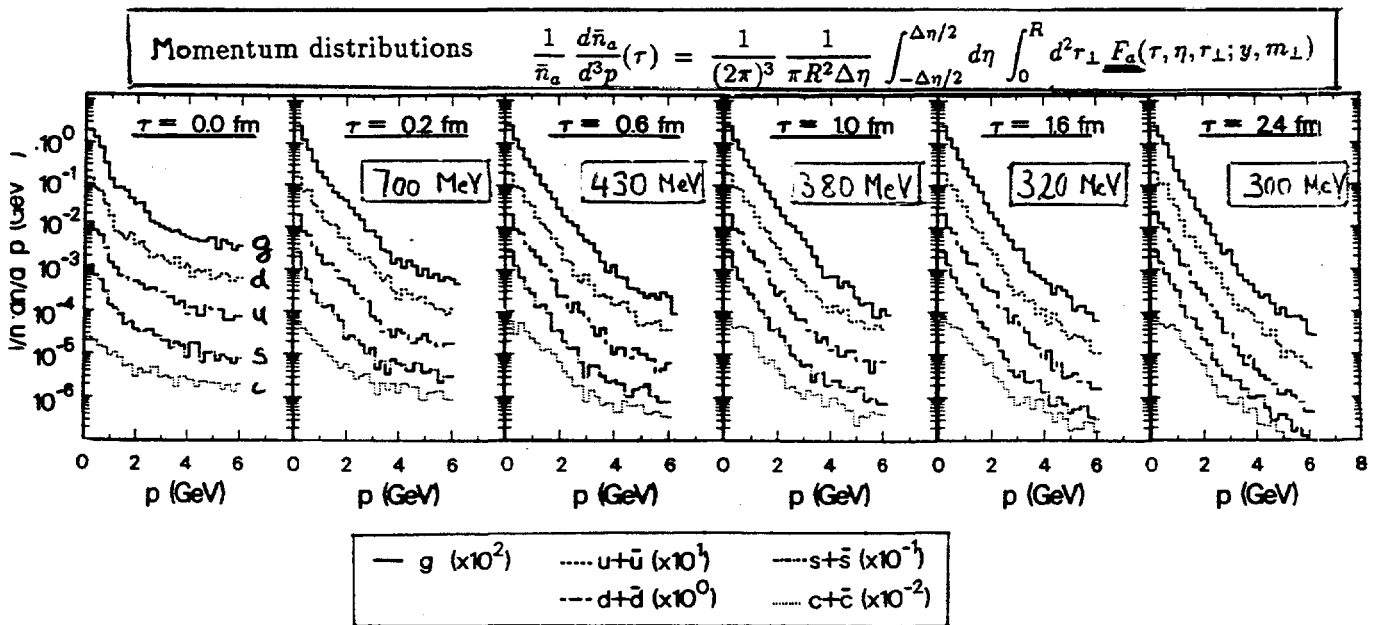


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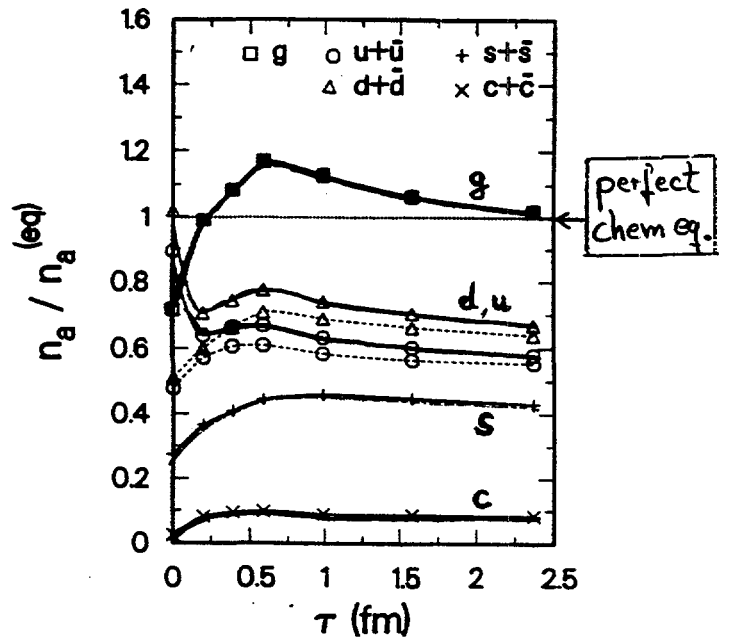
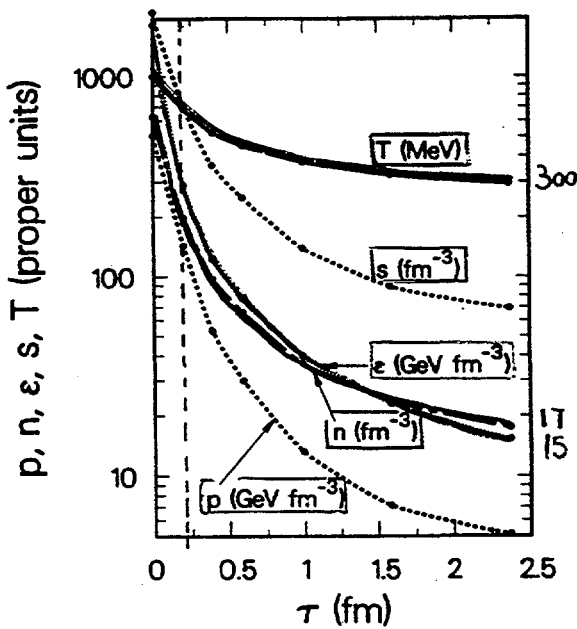
CERN SPS Pb+Pb @ 158 A GeV

Srivastava '97
 & KKG





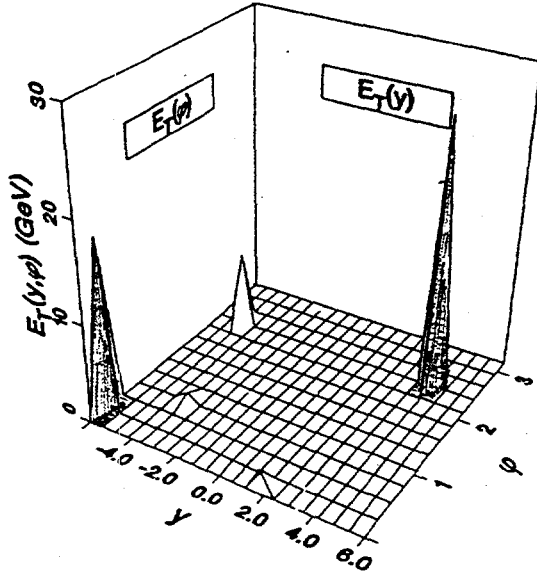
Central Au+Au at $s^{1/2} = 200 \text{ A GeV}$



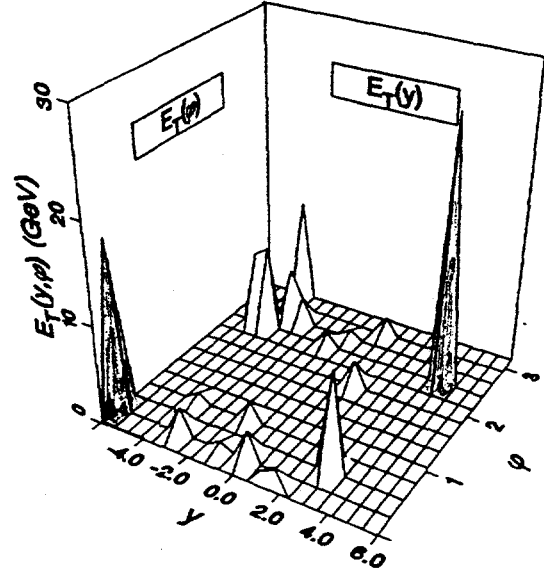
CENTRAL REGION:
 $-1 \leq \eta \leq 1$ $0 \leq r_{\perp} \leq 5 \text{ fm}$

p+p (b=0) at $s^{1/2}=200$ GeV

$\tau = 0.1$ fm

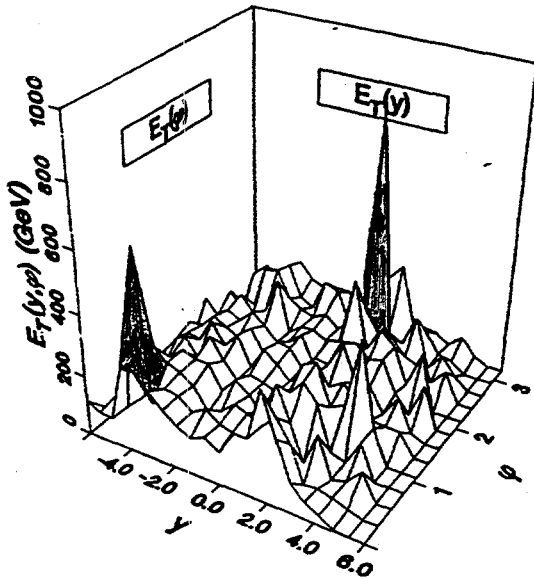


$\tau = 0.4$ fm

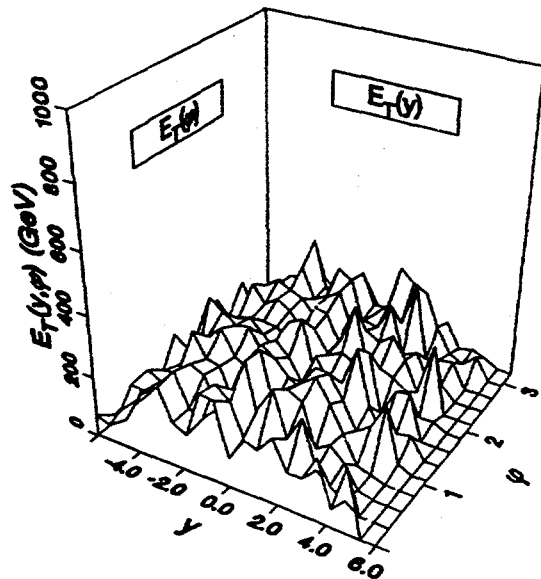


Au+Au (b=0) at $s^{1/2}=200$ A GeV

$\tau = 0.2$ fm



$\tau = 2.4$ fm



"JET QUENCHING" (Wang; Gulassy)