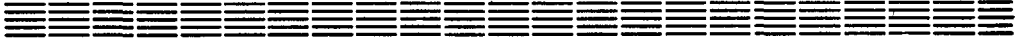




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YEREVAN PHYSICS INSTITUTE



S.M.DARBINIAN, K.A.ISPIRIAN, A.T.MARGARIAN

UNRUH RADIATION OF QUARKS AND THE SOFT PHOTON PUZZLE  
IN HADRONIC INTERACTIONS

30 - 40

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ЦНИИатоминформ  
ЕРЕВАН - 1991

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ՔՎԱՐԿՆԵՐԻ ՈՒՆՐՈՒՅԻ ՃԱՌԱԳԱՅՔՈՒՄԸ Լ ՉԱԻՐՈՆԱՅԻՆ  
ՓՈՒՍԱԶԻԵՑՈՒԹՅՈՒՆՆԵՐՈՒՄ ՓԱՓՈՒԿ ՖՈՏՈՆՆԵՐԻ

ՊՐՈՔԼԵՄԸ

Հաշվված է հաղրոն-հաղրոնային ընդհարումներում մեծ արագացումների ենթարկվող քվարկների կողմից փափուկ Ֆոտոնների առաջացման հավանականությունը, ինչպես նաև զնախառնված է համապատասխան ներդրումը  $P = 70$  ԳԷՎ-ի ժամանակ  $K_p^+$ -փոխազդեցություններում դիտվող  $P_+$ -ի փոքր արժեքներով անոմալ Ֆոտոնների թվի մեջ:

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In the transverse momentum  $P_t$ - distributions of the photon production differential cross sections measured in  $K^+p$ -interactions at 70 GeV/c [1] an excess of anomalous soft photons at  $P_t \lesssim 40$  MeV/c ( $\sim 0.5$  photon per interaction) remains after the extraction of the contribution from the decays of  $\pi^0$  and other hadrons as well as from bremsstrahlung of particles. Such an excess is observed also in  $\mu p$  [2],  $pBe$  and  $pAl$ [3] interactions. The production of these photons seems to be unexplainable since it contradicts to the Low theorem and Landau-Pomeranchuk effect (see for instance [4]). Nevertheless, some success is achieved recently in the interpretation of this phenomenon in the frames of the dense pion gas model [5] and especially in the cold quark-gluon plasma (CQGP) globs model [6].

In this work we calculate the soft photon production probability by the mechanism of the Unruh radiation of quarks and estimate the corresponding contribution into the observed number of anomalous photons [1]. Let us remember that according to the Unruh effect any body moving with acceleration  $\alpha$  in its rest frame finds itself in a bath of thermal radiation with a temperature  $T=\alpha/2\pi K$  ( $K$  is the Boltzman constant,  $\hbar=c=1$ ), which after Compton scattering on the charged body, say electron, becomes real Unruh radiation (see [7-9]). According to this conception this Unruh radiation becomes dominant at very large accelerations corresponding to  $KT \geq 10$  MeV in the case of

electrons. On the other hand the quarks produced in hadron-hadron interactions get accelerations corresponding to  $KT \sim 100$  MeV (see, for instance, [10]). Therefore the Unruh radiation of quarks is an additional source for soft photons.

With the help of the formulae of the work [8] we have calculated the dependence of the Unruh radiation intensity produced by a single quark on  $p_t$  and Feynman variable  $x$  (Fig.1 a and b, respectively) for  $KT=100$  and 800 MeV taking into account the experimental conditions [1]. It is assumed that the quark mass  $M_q = 1$  MeV and the acceleration takes place on a distance  $\sim 1/f$ . The calculation results are normalized at certain points to the experimental data [1] given with their errors. It is seen that the experimental and theoretical data have the same relative behaviour weakly depending on  $\tau$ . However as it is seen from Fig.2 the integral number of the photons  $N_\gamma$  emitted by a single quark according to the calculations depends strongly on  $\tau$ . On the right side of Fig.2 the ordinate axis shows the number of the quarks  $N_q$  with charge  $2/3$  which for the given  $\tau$  can provide a mean number  $\sim 0.5$  anomalous photons emitted per interaction act [1].

From Figs 1 and 2 one can conclude that in order to interpret the observed anomalous soft photons as result of Unruh radiation it is necessary the production of a dozen of quarks with  $KT=600-800$  MeV or from 600 to 150 quarks with  $KT=100-200$  MeV. The above considered radiation may be produced as in the initial phase of QCD parton shower as well as in the final phase when quark-gluon plasma is developed. The observed (see [4]) quadratic dependence of  $N_\gamma$  upon the number  $N_{ch}$  of the produced charged hadrons in the case of quark-gluon plasma, in particular in the CQGP model, is obtained since  $N_\gamma \sim N_{coll} \sim N_p^2 \sim N_{ch}^2$  where  $N_{coll}$  is the number of collisions and  $N_p$  is the number of

partons in plasma. In the initial phase one expects a dependence stronger than the linear one due to the increase of temperature with the growth of  $N_{ch}$  [11]. In order to make more certain conclusions it is necessary to obtain new experimental data especially on  $N_{ch}$ -dependence as well as experimental and theoretical results on low mass  $e^+e^-$ -pairs since in this case too there is an observed excess of soft  $e^+e^-$ -pairs.

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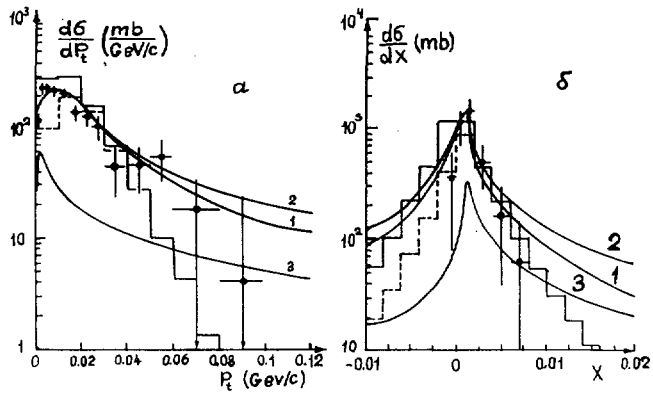


Fig. 1

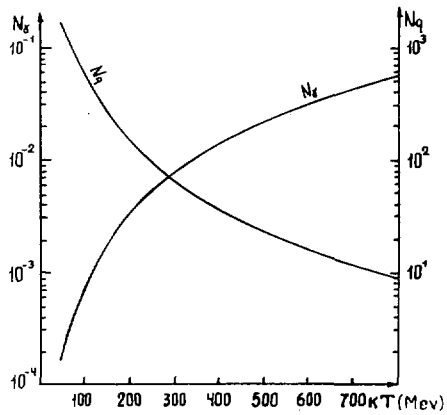


Fig. 2

#### FIGURE CAPTIONS

Fig.1. The spectral dependence of the soft photon number (a) upon  $P_t$  and (b) upon  $x$ . The points are for the experimental data [1]. The solid and dashed histograms are for the results of the CQGP model [6]. The curves 1 and 2 are for results of the present work at  $KT=100$  and  $800$  MeV. The curves 3 are for the results of the bremsstrahlung calculations [1].

Fig.2. The dependence of  $N_\gamma$  and  $N_q$  upon  $KT$ .

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ИЗЛУЧЕНИЕ УНРУ КВАРКОВ И ПРОБЛЕМА МЯГКИХ ФОТОНОВ В  
АДРОННЫХ ВЗАИМОДЕЙСТВИЯХ

(на английском языке, перевод авторов)

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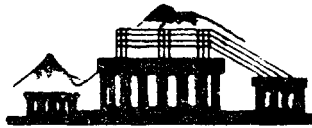
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