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PROGRAM CCOM  
—COUPLED-CHANNELS OPTICAL MODEL CALCULATION  
WITH AUTOMATIC PARAMETER SEARCH—

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Osamu IWAMOTO

日本原子力研究所  
Japan Atomic Energy Research Institute

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Program CCOM  
— Coupled-channels Optical Model Calculation  
with Automatic Parameter Search —

Osamu IWAMOTO

Department of Nuclear Energy System  
Tokai Research Establishment  
Japan Atomic Energy Research Institute  
Tokai-mura, Naka-gun, Ibaraki-ken

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A new program of coupled-channels optical model calculation has been developed for the evaluation of actinide nuclei. The code is composed of by the modules having high independency and large flexibility. The code is written by C++ language using object oriented techniques. The program has capability of fitting of the parameters even for the several nuclei at the same time. The formulae required in the calculation, details of the numerical treatments and the input parameters are described. The examples of the input file and the output are also shown.

Keywords: Nuclear Data, Evaluation, Neutron, Proton, Computer Program, Cross Sections, Optical Model, Coupled Channel Calculation

プログラム CCOM  
— 自動パラメータ検索機能付チャンネル結合光学モデル計算 —

日本原子力研究所東海研究所エネルギーシステム研究部  
岩本 修

(2003年11月7日受理)

アクチノイド原子核の核データ評価のために、新たなチャンネル結合光学モデル計算プログラムを開発した。コード独立性の高いモジュールで構成され、大きな融通性を持っている。コードはオブジェクト指向技術を用いて、C++言語で記述されている。プログラムにはパラメータのフィッティング機能があり、複数の原子核に対しても同時に行うことが可能である。計算に必要な式及び数値的取り扱い、入力パラメータについて記述してある。また入力パラメータの例及びその出力結果を示す。

## Contents

1	Introduction .....	1
2	Model .....	1
2.1	Coupled-channels Equation .....	1
2.2	Cross Section .....	3
3	Details of Calculation .....	6
3.1	Expansion of Potential .....	6
3.2	Solution of Schrödinger Equation .....	7
4	Description of Code .....	8
4.1	List of Files, Functions and Classes .....	9
4.2	OMP Function .....	11
5	Input Parameters.....	14
5.1	Description of Input Parameters.....	14
5.2	Example of Input File.....	16
5.2.1	Example 1: Simple Calculation .....	16
5.2.2	Example 2: Fitting.....	18
	Acknowledgements .....	22
	References .....	22

## 目次

1 序論 .....	1
2 モデル .....	1
2.1 チャンネル結合方程式 .....	1
2.2 断面積 .....	3
3 計算の詳細 .....	6
3.1 ポテンシャル展開 .....	6
3.2 Schrödinger 方程式の解 .....	7
4 コードの説明 .....	8
4.1 ファイル、関数、クラスのリスト .....	9
4.2 OMP 関数 .....	11
5 入力パラメータ .....	14
5.1 入力パラメータの説明 .....	14
5.2 例 .....	16
5.2.1 例1：簡単な計算 .....	16
5.2.2 例2：フィッティング .....	18
謝辞 .....	22
参考文献 .....	22

## 1 Introduction

Optical model is one of important methods for the evaluation of the nuclear reaction data, since it is used for the estimation of total and scattering cross sections, and of transmission coefficients for statistical model calculations. In order to make the Japanese Evaluated Nuclear Data Library(JENDL), spherical optical model codes with statistical model calculation such as ELIESE[1] and CASTHY[2] were developed at the Nuclear Data Center in JAERI. For the actinide nuclei, however, coupled-channels(CC) optical model[3, 4] can describe scattering process better than spherical optical model, because such nuclei having large deformation cause strong coupling between levels in the rotational band.

A few codes[5, 6] developed in the foreign countries were used for the coupled-channels calculations for nuclear data evaluations. Such codes are usually used as a black box and it is very difficult to add new features to the codes. It is valuable to develop a new code, whose details are known, for the further improvements of the calculation. Then a new code (CCOM) of coupled-channels optical model calculation has been developed at the Nuclear Data Center in JAERI for the evaluation of actinide nuclei. The CCOM is composed of the modules having high independency and large flexibility. The CCOM is written by C++ language using object oriented technique. The CCOM also has capability of the fitting of the parameters even for the several nuclei at the same time.

This report describes the formulae required in the calculation, details of the numerical treatments and input parameters. The examples of the input file and its output are also shown.

## 2 Model

### 2.1 Coupled-channels Equation

In this model, it is assumed that the projectile has no internal structure and the interaction with the target can be described by the one body local potential. The formulations described bellow are mainly based on the Ref. [4].

The model Hamiltonian of the system is written as

$$H = H_A + H_a + T + V, \quad (1)$$

where  $H_A$  and  $H_a$  are Hamiltonians of the target and projectile, respectively.  $T$  and  $V$  are kinetic energy operator and optical potential of relative motion of target and projectile, respectively. The  $V$  depends on the relative coordinate of particles and also the internal coordinate of the target. The wave function of the target  $\Phi_n$  in the state  $n$  is the eigen

function of  $H_A$ :  $H_A \Phi_n = \epsilon_n \Phi_n$ . Expanding by angular momentum, the total wave function is written as

$$\Psi = \sum_{c,J,M} \psi_c^{JM} r^{-1} f_c^{JM}(r), \quad (2)$$

$$\psi_c^{JM} = \left[ \left[ i^l Y_l \otimes \chi_s \right]_j \otimes \Phi_{n(I)} \right]_{JM}, \quad (3)$$

where  $c$  is the channel of reaction representing  $\{l, j, n\}$ ,  $f_c(r)$  shows the radial dependence of the wave function for the channel  $c$ ,  $I$  is the spin of the target in the state  $n$ ,  $Y_l$  is spherical harmonics,  $\chi_s$  is the wave function of the projectile having the spin  $s$  and  $[\phi_{j_1} \otimes \phi_{j_2}]_{JM} = \sum_{m_1, m_2} \langle j_1 j_2 m_1 m_2 | JM \rangle \phi_{j_1} \phi_{j_2}$ . From Eq. (1) and Eq. (2) with integrating for the coordinates except  $r$  for the wave function, the coupled-channels Schrödinger equation is obtained,

$$\left( \frac{d^2}{dr^2} - \frac{l_c(l_c + 1)}{r^2} + k_c^2 \right) f_c^{JM}(r) = \sum_{c'} \langle \psi_c | V | \psi_{c'} \rangle f_{c'}^{JM}(r), \quad (4)$$

where  $k_c$  indicates wave number in the channel  $c$ .

The deformation of the potential is assumed to be described by the the potential radius which is the function of the polar angles ( $\theta'$  and  $\phi'$ ) in the body-fixed frame. The radius is written by the expansion of the spherical harmonics,

$$R(\theta', \phi') = R_o \left\{ 1 + \sum_{\lambda} \beta_{\lambda 0} Y_{\lambda 0}^*(\theta', \phi') \right\}, \quad \lambda = 2, 4, 6, \dots \quad (5)$$

Here the potential is assumed to have an axial symmetric shape and not to be changed by the rotational excitation. The potential can be expanded by the spherical harmonics as

$$\begin{aligned} V(\mathbf{r}') &= \sum_{\lambda} V_{\lambda}(r') \left\{ i^{\lambda} Y_{\lambda 0}(\hat{\mathbf{r}}') \right\}^* \\ &= \sum_{\lambda} V_{\lambda}(r) \left\{ i^{\lambda} \sum_{\mu} D_{0\mu}^{\lambda}(\Omega) Y_{\lambda\mu}(\hat{\mathbf{r}}) \right\}^* \\ &= \sum_{\lambda\mu} V_{\lambda\mu}(r) \left\{ i^{\lambda} Y_{\lambda\mu}(\hat{\mathbf{r}}) \right\}^*, \end{aligned} \quad (6)$$

$$V_{\lambda}(r') = \int d\hat{\mathbf{r}}' V(\mathbf{r}') i^{\lambda} Y_{\lambda 0}(\hat{\mathbf{r}}'), \quad (7)$$

where  $r' = |\mathbf{r}'|$ ,  $\hat{\mathbf{r}}' = (\theta', \phi')$  and  $d\hat{\mathbf{r}}' = \sin\theta' d\theta' d\phi'$ . The  $D_{mm'}^l$  is the D function which is defined by the matrix element of the rotation operator which is a function of the Euler angle,  $\Omega$ . The wave function of the target nucleus  $\Phi_{IMK}$  with the spin  $I$  for the level in the rotational band is expressed as

$$\Phi_{n(I)} = \Phi_{IMK} = \frac{\hat{I}}{\sqrt{8\pi^2}} D_{KM}^I(\Omega) \phi_K(\xi), \quad (8)$$



where  $\phi_K$  and  $\xi$  are the intrinsic wave function and its coordinate, respectively, and  $\hat{I} = \sqrt{2I+1}$ . The reduced matrix element is calculated as

$$\begin{aligned} \langle I'M'|V_{\lambda\mu}|IM\rangle &= \int \int d\xi d\Omega \Phi_{I'M'K}^* V_{\lambda\mu} \Phi_{IMK} \\ &= \langle I\lambda M\mu|I'M'\rangle \langle I' || V_\lambda || I \rangle \end{aligned} \quad (9)$$

$$\langle I' || V_\lambda || I \rangle = \hat{I} \hat{I}'^{-1} V_\lambda(r) \langle I\lambda K 0 | I' K \rangle. \quad (10)$$

Finally, from Eqs. (3), (6) and (9) the coupling matrix element  $\langle \psi_c | V | \psi_{c'} \rangle$  in Eq. (4) is obtained by

$$\begin{aligned} \langle \psi_c | V | \psi_{c'} \rangle &= \frac{1}{\sqrt{4\pi}} \sum_\lambda i^{l-l'-\lambda} (-)^{J-I-s} \hat{l} \hat{l}' \hat{j} \hat{j}' \hat{I}' \langle ll' 00 | \lambda 0 \rangle \\ &\times W(jj' ll'; \lambda I) W(jj' II'; \lambda J) \langle I' || V_\lambda || I \rangle, \end{aligned} \quad (11)$$

where  $W$  shows the Racah coefficient.

## 2.2 Cross Section

In asymptotic region where the nuclear potential is negligible except for the Coulomb potential, the Schrödinger equation has the form

$$\left( \frac{d^2}{dr^2} - \frac{l(l+1)}{r^2} - \frac{2\eta k}{r} + k^2 \right) f_l(r) = 0, \quad (12)$$

where  $\eta$  is the Coulomb parameter:  $\eta = ZZ'e^2\mu/\hbar^2 k$ . The Coulomb wave functions  $F_l(\rho)$  (regular) and  $G_l(\rho)$  (irregular) are the independent solutions of Eq. (12), where  $\rho = k\tau$ . For  $\rho \rightarrow \infty$

$$F_l(\rho) \sim \sin(\rho - \eta \log 2\rho - \frac{1}{2}l\pi + \sigma_l), \quad (13)$$

$$G_l(\rho) \sim \cos(\rho - \eta \log 2\rho - \frac{1}{2}l\pi + \sigma_l), \quad (14)$$

$$\sigma_l = \arg \Gamma(l+1+i\eta). \quad (15)$$

In the case for the closed channels, the Schrödinger equation is given by

$$\left( \frac{d^2}{dr^2} - \frac{l(l+1)}{r^2} - \frac{2\eta k}{r} - k^2 \right) f_l(r) = 0. \quad (16)$$

The solution of Eq. (16) is given by Whittaker's function  $W_{-\eta, l+1/2}(2\rho)$ [7], which is decreasing exponentially. If  $\eta = 0$ , the solution is given by the irregular modified spherical Bessel function multiplied by  $\rho$ .

The cross sections are obtained from the asymptotic behaviour of the wave function, which is given by the solution of Eq. (4). In the present system, total angular momentum  $J$  and

parity  $\pi$  are conserved through the scattering process. The C-matrix element  $C_{cc'}^{J\pi}$  for  $J^\pi$  from the initial channel  $c'$  to the final channel  $c$  is defined by

$$\tilde{f}_{cc'}^{J\pi} \underset{r \rightarrow \infty}{\sim} \begin{cases} F_{l_c}(k_c r) \delta_{cc'} + \sqrt{\frac{k_{c'}}{k_c}} C_{cc'}^{J\pi} \{G_{l_c}(k_c r) + iF_{l_c}(k_c r)\} & (E > 0), \\ W_{-\eta_c, l_c + 1/2}(2k_c r) & (E \leq 0), \end{cases} \quad (17)$$

where  $\tilde{f}_{cc'}^{J\pi}$  is the solution of Eq. (4). The C-matrix is related to S-matrix by

$$C_{cc'}^{J\pi} = \frac{i}{2} (\delta_{cc'} - S_{cc'}^{J\pi}). \quad (18)$$

If  $N$  channels are coupled for  $J^\pi$ , Eq. (4) has  $N$  sets of the independent solutions having no singularity at  $r = 0$ . The  $i$ -th set of the solution is represented by a vector with the elements of  $N$  channels,

$$\mathbf{f}^i = \begin{pmatrix} f_i^1 \\ \vdots \\ f_i^N \end{pmatrix}. \quad (19)$$

The wave function having the asymptotic behaviour of Eq. (17) for the initial channel  $c'$  can be expressed by a linear combination of the  $n$  independent solutions,

$$\tilde{\mathbf{f}}^{c'} = \begin{pmatrix} \tilde{f}_{1c'} \\ \vdots \\ \tilde{f}_{Nc'} \end{pmatrix} = \sum_{i=1}^N a_i^{c'} \mathbf{f}^i, \quad (c' = 1 \dots N). \quad (20)$$

From the condition of the smooth connection of Eq. (20) and Eq. (17),  $2N$  linear equations of  $C_{cc'}$  and  $a_i^{c'}$  are obtained,

$$\begin{pmatrix} \mathbf{f}_1 & -\mathbf{H}_1 \\ \mathbf{f}_2 & -\mathbf{H}_2 \end{pmatrix} \begin{pmatrix} \mathbf{A} \\ \bar{\mathbf{C}} \end{pmatrix} = \begin{pmatrix} \mathbf{F}_1 \\ \mathbf{F}_2 \end{pmatrix}, \quad (21)$$

where

$$\mathbf{F}_i = \text{diag}(F_{l_1}(k_1 r_i), \dots, F_{l_c}(k_c r_i), \dots, F_{l_N}(k_N r_i)), \quad (22)$$

$$\mathbf{G}_i = \text{diag}(G_{l_1}(k_1 r_i), \dots, G_{l_c}(k_c r_i), \dots, G_{l_N}(k_N r_i)), \quad (23)$$

$$\mathbf{H}_i = \mathbf{G}_i + i\mathbf{F}_i, \quad (24)$$

$$\text{diag}(X_1, \dots, X_N) = \begin{pmatrix} X_1 & & 0 \\ & \ddots & \\ 0 & & X_N \end{pmatrix}, \quad (25)$$

$$\mathbf{f}_i = \left( \tilde{\mathbf{f}}^1(r_i), \dots, \tilde{\mathbf{f}}^c(r_i), \dots, \tilde{\mathbf{f}}^N(r_i) \right), \quad (26)$$

$$\mathbf{A} = \begin{pmatrix} a_1^1 & \cdots & a_N^1 \\ \vdots & \ddots & \vdots \\ a_1^N & \cdots & a_N^N \end{pmatrix}, \quad (27)$$

$$\bar{\mathbf{C}} = \begin{pmatrix} \bar{C}_{11} & \cdots & \bar{C}_{1N} \\ \vdots & \ddots & \vdots \\ \bar{C}_{N1} & \cdots & \bar{C}_{NN} \end{pmatrix}, \quad (28)$$

$$\bar{C}_{cc'} = \sqrt{\frac{k'_c}{k_c}} C_{cc'}. \quad (29)$$

Here, the matching condition is that the wave functions coincide at the two asymptotic radial points ( $r_1, r_2$ ). The matrix element of  $H_i$  will be changed to the Whittaker's function  $W(2k_c r_i)$  in the case for the closed channel. The final result of the C-matrix is obtained from Eqs. (21) and (29).

The cross sections of total ( $\sigma_t$ ), scattering ( $\sigma_i$ ) to  $i$ -th level and reaction ( $\sigma_r$ ) are calculated from the C-matrix,

$$\sigma_t = \sum_{J^\pi} \frac{4\pi}{k_0^2} \frac{2J+1}{(2s+1)(2I_0+1)} \sum_{cc'(n_{c'}=0)} \text{Im} C_{cc'}^{J^\pi}, \quad (30)$$

$$\sigma_i = \sum_{J^\pi} \frac{4\pi}{k_0^2} \frac{2J+1}{(2s+1)(2I_0+1)} \sum_{cc'(n_{c'}=0, n_c=i)} |C_{cc'}^{J^\pi}|^2, \quad (31)$$

$$\sigma_r = \sum_{J^\pi} \frac{\pi}{k_0^2} \frac{2J+1}{(2s+1)(2I_0+1)} \sum_{cc'(n_{c'}=0)} \left( 1 - |S_{cc'}^{J^\pi}|^2 \right), \quad (32)$$

where  $k_0$  and  $I_0$  indicate wave number and spin of target for initial channel, respectively. The conditions of  $n_c = i$  and  $= 0$  means sum over the channels in the  $i$ -th excited or initial channel levels, respectively. The S-matrix element  $S_{cc'}^{J^\pi}$  in Eq. (32) is obtained from Eq. (18). The total cross section  $\sigma_t$  and elastic scattering cross section  $\sigma_0$  are not defined for a charged projectile. The transmission coefficient  $T_l$  for the orbital angular momentum  $l$  is obtained by

$$T_l = \frac{1}{2l+1} \sum_{J^\pi} \frac{2J+1}{(2s+1)(2I_0+1)} \sum_{cc'(n_{c'}=0, l_{c'}=l)} \left( 1 - |S_{cc'}^{J^\pi}|^2 \right). \quad (33)$$

The scattering amplitude  $A_{imM, m'M'}$  from the initial state having  $z$ -projections of the spin of the projectile ( $m'$ ) and target ( $M'$ ) to the final state ( $mM$ ) at the  $i$ -th excited level is

obtained by

$$\begin{aligned}
 A_{imM,m'M'}(\theta) = & \frac{4\pi}{k_0} \sum_{cc'(n_c=i, n_{c'}=0)} \sum_{\substack{l_c l_{c'} j_c j_{c'} J \\ m_{l_c} m_{l_{c'}} m_{j_c} m_{j_{c'}} M_J}} \\
 & \langle l_c s_c m_{l_c} m | j_c m_{j_c} \rangle \langle j_c I_c m_{j_c} M | J M_J \rangle \\
 & \times \langle l_{c'} s_{c'} m_{l_{c'}} m' | j_{c'} m_{j_{c'}} \rangle \langle j_{c'} I_{c'} m_{j_{c'}} M' | J M_J \rangle \\
 & \times e^{i(\sigma_c + \sigma_{c'})} C_{cc'} Y_{l_c m_{l_c}}(\theta, 0) Y_{l_{c'} 0}(0, 0)^*, \quad (34)
 \end{aligned}$$

where  $\theta$  is the scattering angle. The projectile is assumed to move along the  $z$ -axis at initial channel. The angular distribution of scattering cross sections to the  $i$ -th level are calculated by

$$\frac{d\sigma_i}{d\Omega}(\theta) = \sum_{mMm'M'} \frac{1}{(2s+1)(2I_0+1)} |f_C(\theta)\delta_{mm'}\delta_{MM'} + A_{imM,m'M'}(\theta)|^2, \quad (35)$$

where  $f_C$  is Coulomb scattering amplitude given by

$$f_C(\theta) = -\frac{\eta_0}{2k_0 \sin \frac{\theta}{2}} \exp \left\{ -i\eta_0 \log \left( \sin^2 \frac{\theta}{2} \right) + 2i\sigma_0 \right\}, \quad (36)$$

$$\sigma_0 = \arg \Gamma(1 + i\eta_0), \quad (37)$$

$$\eta_0 = \frac{ZZ'e^2\mu}{\hbar^2 k_0} = \frac{\mu c^2 \alpha}{\hbar c k_0}, \quad (38)$$

where  $\alpha$  is the fine-structure constant and  $c$  is the speed of light.

### 3 Details of Calculation

#### 3.1 Expansion of Potential

The optical potentials are assumed to have the shapes of Woods-Saxon and its derivative.

$$V(r) = V_i f_i(r) + W_j g_j(r) + V_s h_s(r) \boldsymbol{\sigma} \cdot \mathbf{1} + V_C(r), \quad (39)$$

where  $V_i$ ,  $W_j$  and  $V_s$  are the strength of potential of complex number and  $V_C$  shows Coulomb potential which is calculated from charge distribution of target nucleus. The  $f_i(r)$ ,  $g_j(r)$  and  $h_s(r)$  are the form factors given by

$$f_i(r) = \left( 1 + e^{\frac{r-R_i}{a_i}} \right)^{-1}, \quad (40)$$

$$g_i(r) = -4a_i \frac{d}{dr} f_i, \quad (41)$$

$$h_i(r) = \left( \frac{\hbar}{m_\pi c} \right)^2 \frac{1}{r} \frac{d}{dr} f_i(r), \quad (42)$$

where  $\left(\frac{\hbar}{m\pi c}\right)^2 = 2$ .

The charge distribution  $\rho(\mathbf{r}')$  of target is assumed to be uniform inside of the charge radius  $R(\hat{\mathbf{r}}')$ , Eq. (5), and zero outside. The Coulomb potential  $V_C$  is expanded by the spherical harmonics,

$$\begin{aligned} V_C(\mathbf{r}) &= \int d\mathbf{r}' \frac{ZZ'e^2\rho(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} \\ &= \frac{ZZ'e^2}{V_0} \int d\mathbf{r}' \Theta(R(\hat{\mathbf{r}}')-r') \sum_{lm} \frac{4\pi}{2l+1} \frac{r_{>}^{l+1}}{r_{<}^l} Y_{lm}(\hat{\mathbf{r}}) Y_{lm}(\hat{\mathbf{r}})^* \\ &= \sum_{lm} Y_{lm}(\hat{\mathbf{r}}) \int d\hat{\mathbf{r}}' Y_{lm}^*(\hat{\mathbf{r}}') \tilde{V}_{Cl}(r, \hat{\mathbf{r}}') \frac{3ZZ'e^2}{(2l+1)R_0^3 A}, \end{aligned} \quad (43)$$

for  $R(\hat{\mathbf{r}}') \leq r'$

$$\tilde{V}_{Cl}(r, \hat{\mathbf{r}}') = \frac{R(\hat{\mathbf{r}}')^{l+3}}{(l+3)r} \quad (44)$$

for  $R(\hat{\mathbf{r}}') > r'$

$$\tilde{V}_{Cl}(r, \hat{\mathbf{r}}') = \frac{r^2}{l+3} + \begin{cases} \frac{r^2}{l-2} \left(1 - \frac{r^{l-2}}{R(\hat{\mathbf{r}}')^{l-2}}\right) & (l \neq 2), \\ r^l \log \frac{R(\hat{\mathbf{r}}')}{r} & (l = 2), \end{cases} \quad (45)$$

where  $r_{<} = \min(r, r')$ ,  $r_{>} = \max(r, r')$  and  $\Theta(x)$  is the step function

$$\Theta(x) = \begin{cases} 0 & (x < 0), \\ 1 & (x > 0). \end{cases} \quad (46)$$

The  $V_0$  is volume of charge distribution and  $A$  is the ratio to the volume of the sphere of radius  $R_0$ ,

$$V_0 = \int d\mathbf{r}' \Theta(R(\hat{\mathbf{r}}')-r') \quad (47)$$

$$= R_0^3 \left[ \frac{4\pi}{3} + \sum_l \beta_l^2 + \frac{1}{6\sqrt{\pi}} \sum_{l_1 l_2 l_3} \hat{l}_1 \hat{l}_2 \hat{l}_3 \begin{pmatrix} l_1 & l_2 & l_3 \\ 0 & 0 & 0 \end{pmatrix}^2 \beta_{l_1} \beta_{l_2} \beta_{l_3} \right] \quad (48)$$

$$A = \frac{3V_0}{4\pi R_0^3} \quad (49)$$

The expanded potentials, Eq. (7) and (43), are calculated by the numerical integration of optical potentials at radial grid points to solve the Schrödinger equation, Eq. (4).

### 3.2 Solution of Schrödinger Equation

The Schrödinger equation, Eq. (4), can be written simply as

$$\frac{d^2}{dr^2} f_i(r) = \sum_j V_{ij}(r) f_j(r), \quad (50)$$

and by vector representation,

$$f''(r) = V(r)f(r). \quad (51)$$

The modified Numerov method[8] is used to integrate Eq. (51). It is obtained from the Taylor's expansion of  $f(r)$  with Eq. (51),

$$\begin{aligned} \xi(r+h) &\equiv f(r+h) - \frac{h^2}{12}f''(r+h) \\ &= 2f(r) + \frac{5}{6}f''(r) - f(r-h) + \frac{h^2}{12}f''(r-h) + O(h^6) \\ &= 2\xi(r) - \xi(r-h) + u(r) + O(h^6), \end{aligned} \quad (52)$$

$$\begin{aligned} u(r) &= h^2f''(r) \\ &= \frac{h^2}{1 - h^2V(r)/12}\xi(r) \\ &= h^2V(r) \left(1 + \frac{h^2}{12}V(r)\right)\xi(r) + O(h^6) \end{aligned} \quad (53)$$

and

$$f(r) = \xi(r) + \frac{1}{12}u(r). \quad (54)$$

The value of the  $i$ -th solution at matching radius is obtained by Eq. (52) and Eq. (54) with neglecting the terms of  $O(h^6)$  using the initial value of

$$\xi^i(0) = 0, \quad (55)$$

$$\xi^i(h) = \begin{pmatrix} 0 \\ \vdots \\ 0 \\ \xi_i^i(h) \\ 0 \\ \vdots \\ 0 \end{pmatrix}, \quad (56)$$

$$\xi_i^i(h) = \frac{(kh)^{l+1}}{(2l+1)!!}, \quad (57)$$

$$k = \sqrt{\frac{2\mu(E - V(0))}{\hbar^2}}.$$

## 4 Description of Code

The code is written by C++ language using the standard template library (STL) and GNU Scientific Library (GSL)[9]. The list of files with the brief description about functions and classes of the code is shown.

## 4.1 List of Files, Functions and Classes

- `calc.cpp`, `calc.h`: main calculation functions

*functions:*

- *cmatrix*: calculate C-matrix
- *smatrix*: calculate S-matrix from C-matrix
- *scat\_amp*: calculate scattering amplitude from C-matrix
- *cross\_section*: calculate cross sections
- *cross\_section\_da*: calculate angular distributions of scattering cross sections
- *cross\_section\_integ*: calculate total, reaction and scattering cross sections
- *calc*: main function to calculate cross sections at each incident energies

- `cblas-mat.cpp`, `cblas-mat.h`: matrix class

*classes:*

- *cblas\_matrix\_z*: matrix class of complex number used to call CBLAS functions used in the function *solve*.

- `ccom.cpp`: main program

*functions:*

- *main*: main function; makes a instance of *Input\_param* class, reads input parameters and calls *main\_fit* function
- *main\_fit*: calls the function *fit* for the fitting, and the function *calc* to get the final results

- `const.h`: define some physical constants

$\hbar c$ , fine-structure constant  $\alpha$ , atomic mass unit ...

- `fit.cpp`, `fit.h`: fitting functions

*functions:*

- *fit*: main fitting function; iterate up to the convergence
- *calc*: obtain cross sections for all experimental data points
- *expb\_f*: make a vector of difference between experiment and calculated values multiplied by the weights of inverse of experimental errors
- *expb\_df*: make Jacobian matrix, derivative of calculated values by parameters

- **func.cpp, func.h**: special functions
  - spherical harmonics, coulomb wave functions, Clebsch-Gordan coefficient, Racah coefficient, ...
- **input.cpp, input.h**: class of input parameters
  - classes*:
    - *cs\_data*: contain the experimental data
    - *Sys\_param*: parameters to define system of scattering
    - *Cal\_param*: energies and angles to calculate as final results
    - *Input\_param*: contain all input parameters
- **omp.cpp, omp.h**: calculate the optical potential parameters from the input parameters
  - functions*
    - *get\_omp*: select the function *get\_omp\_id* to calculate optical potential parameters depending on *id*
    - *get\_omp\_id*: calculate optical potential from the list of input parameters by the formula of *id*
- **potential.cpp, potential.h**: class of form factor and potential
  - classes*:
    - *FormFactor*: calculate form factors
    - *Potential*: calculate the optical potential and get multi-pole expansion of potential
- **solv.cpp, solv.h**: functions to integrate Schrödinger equation
  - functions*:
    - *gcouple*: calculate angular momentum dependent coupling coefficients
    - *set\_vcouple*: calculate coupling matrix elements
    - *set\_u*: calculate  $u$ , Eq. (53)
    - *solve*: integrate Schrödinger equation by the modified Numerov method and obtain the wave function at matching radius
- **system.cpp, system.h**: define system of scattering
  - classes*:



- *Cross\_Section*: contain calculated cross sections
- *Particle*: properties of particle, mass,  $Z$ ,  $A$ , levels, spin, parity
- *Quantum*: quantum number to define channel,  $l$ ,  $l'$ ,  $s$ ,  $s'$ ,  $j$ ,  $j'$ ,  $J^\pi$ , ...
- *Channel*: contain all channel information
- *System*: contain all parameters to define system of scattering

## 4.2 OMP Function

The optical model parameters (OMP) are calculated from the input parameters by the function *get\_omp\_id*, which can be modified to get different OMP energy dependence. For a simple example, the OMP having the linear energy and asymmetry parameter dependence are shown.

```

1 void get_omp_om0(System *sys, valarray<double> &p,
2                 valarray<double> &omp, valarray<double> &beta)
3 {
4     double en=sys->energy_lab;
5
6     int z=sys->targ.z, a=sys->targ.a, n=a-z;
7     double eta=double(n-z)/double(a);
8
9     double v0, r0, a0;
10
11    omp.resize(21);
12
13    for(int i=0, i3; i<7; i++){
14        i3=i*3;
15        v0=0;
16        r0=0;
17        a0=0;
18        switch(i){
19            case 0: // V
20                v0=-(p[0]-p[1]*eta-p[2]*en);
21                r0=p[3];
22                a0=p[4];
23                break;
24            case 1: // W
25                break;
26            case 2: // Vs
27                break;
28            case 3: // Ws
29                v0=-(p[5]-p[6]*eta+p[7]*en);
30                r0=p[8];
31                a0=p[9];
32                break;
33            case 4: // Vso
34                v0=-p[10];

```

```

35         r0=p[11];
36         a0=p[12];
37         break;
38     case 5: // Wso
39         break;
40     case 6: // Vc
41         v0=sys->get_zz();
42         r0=p[13];
43         break;
44     default:
45         break;
46     }
47     omp[i3]=v0;
48     omp[i3+1]=r0;
49     omp[i3+2]=a0;
50 }
51 beta.resize(2);
52 beta[0]=p[14];
53 beta[1]=p[15];
54 }

```

The lines 1 and 2 define the arguments of the function. The arguments `sys` and `p` are the input parameters. The `sys` contains all parameters of the present system and the `p` is the array of parameters for the OMP calculation. The function will return `omp` and `beta`, which are the OMP and deformation parameters, respectively,

$$\text{omp}[i * 3] = V_i, \quad (58)$$

$$\text{omp}[i * 3 + 1] = r_i, \quad (59)$$

$$\text{omp}[i * 3 + 2] = a_i, \quad (60)$$

$$\text{beta}[j] = \beta_{2(j+1)}, \quad (61)$$

where  $i = 0, \dots, 6$ : 0; real Woods-Saxon, 1; imaginary Woods-Saxon, 2; real derivative Woods-Saxon, 3; imaginary derivative Woods-Saxon, 4; real spin-orbit, 5; imaginary spin-orbit, 6; Coulomb. The line 4-7 define  $E$ ,  $Z_t$ ,  $A_t$ , etc. The line 9-50 calculate OMP parameters and the line 51-53 define  $\beta$ .

The OMP having more complicated energy dependence is shown as the next example.

```

1 void get_omp_suk2003(System *sys, valarray<double> &p,
2     valarray<double> &omp, valarray<double> &beta)
3 {
4     double en0=sys->energy_lab, en, en2, ef;
5
6     int zp=sys->proj.z, z=sys->targ.z, a=sys->targ.a, n=a-z;
7     double c0, c1, c2, eta=double(n-z)/double(a), a3=pow(a, 1./3.);
8
9     double v0, r0, a0;
10

```

```

11  omp.resize(21);
12
13  if(zp == 1 && sys->efp != 0)
14      ef=sys->efp;
15  else if(zp == 0 && sys->efn != 0)
16      ef=sys->efn;
17  else
18      ef=p[35];
19
20  en=en0-ef;
21  en2=en*en;
22
23  c0= zp > 0 ? 1 : -1;
24  c1=1+c0*p[28]*eta/(p[0]+p[1]*(a-232)+p[4]);
25  c2=(p[5]*p[4]*exp(-p[5]*en)-p[2]-2*p[3]*en)*c1;
26
27  for(int i=0; i<7; i++){
28      v0=0;
29      r0=0;
30      a0=0;
31      switch(i){
32      case 0: // V
33          v0=(p[0]+p[1]*(a-232)+p[2]*en+p[3]*en2 +p[4]*exp(-p[5]*en))*c1
34              +p[27]*z*zp/a3*c2;
35          v0*=-1;
36          r0=p[17]*(1-p[18]*en2/(en2+p[16]*p[16]));
37          a0=p[19]+p[20]*en;
38          break;
39      case 1: // W
40          v0=-p[10]*en2/(en2+p[11]*p[11]);
41          r0=p[23];
42          a0=p[24];
43          break;
44      case 2: // Vs
45          break;
46      case 3: // Ws
47          v0=-(p[6]+p[7]*(a-232)+c0*p[29]*eta)*exp(-p[8]*en)
48              *en2/(en2+p[9]*p[9]);
49          r0=p[21];
50          a0=p[22];
51          break;
52      case 4: // Vso
53          v0=-p[12]*exp(-p[13]*en);
54          r0=p[30];
55          a0=p[31];
56          break;
57      case 5: // Wso
58          v0=-p[14]*en2/(en2+p[15]*p[15]);
59          r0=p[30];
60          a0=p[31];
61          break;
62      case 6: // Vc
63          v0=sys->get_zz();

```

```

64         r0=p[25];
65         break;
66     default:
67         break;
68     }
69
70     int i3=i*3;
71     omp[i3]=v0;
72     omp[i3+1]=r0;
73     omp[i3+2]=a0;
74 }
75
76 if(sys->beta.size() > 0){
77     beta.resize(sys->beta.size());
78     beta=sys->beta;
79 }
80 else{
81     beta.resize(3);
82     beta[0]=p[32];
83     beta[1]=p[33];
84     beta[2]=p[34];
85 }
86 }

```

The parameters used in the code is almost same as the previous example. The additional parameters are Fermi energies in lines 13-18. The lines 76-85 shows that the deformation parameter will be used if defined in the system section of the input.

## 5 Input Parameters

### 5.1 Description of Input Parameters

The input parameters are identified by tags with free format. The tags and values should be separated by white-space characters. The parenthesis is used to define variables whose number is not fixed. The input file is constructed by the repetition of “**tag** *val*<sub>1</sub> *val*<sub>2</sub> ...”.

The units used in the input and output are MeV, mb and fm for the energy, cross section and length, respectively. For the angular distributions of the cross section,  $\frac{d\sigma}{d\Omega}$ , mb/sr is used with the angle in degree.

- *comment*
  - #: comment out until end of line
  - /\* ... \*/: comment out from “/\*” to “\*/”
- *system*: define system

- **system**: add new system; system is identified by the number 0, 1, 2, ...
  - **za**  $Z_p$   $A_p$   $m_p$   $Z_t$   $A_t$   $m_t$ : proton number, mass number and mass of projectile and target
  - **lev**  $I$   $\pi$   $Ex$ : add level of spin  $I$ , parity  $\pi(+/-)$ , excitation energy  $Ex$
  - **efermi**  $\epsilon_{F_p}$   $\epsilon_{F_n}$ : Fermi energy for proton and neutron (used for some OMP parametrisation)
  - **beta** ( $\beta_2$   $\beta_4$  ...): deformation parameters
  - **mesh**  $r_m$   $N$ : matching radius and number of mesh points (should be defined before the tag **system**)
  - **jmax**  $J_m$ : maximum of total angular momentum to calculate (should be defined before the tag **system**)
  - **vlmax**  $\lambda_m$ : maximum angular momentum for multi-pole expansion of potential (should be defined before the tag **system**)
- **calculation**: make final result
    - **ene** ( $E_1$ ,  $E_2$ , ...): incident energies to be calculated (should be defined before the tag **calc**)
    - **angle**  $N_a$ ,  $\Delta_a$ : number of angles, increase of angle for angular distribution (should be defined before the tag **calc**)
    - **calc** ( $i$   $j$  ...): make calculation for the system  $i$ ,  $j$ , ...
  - **experimental data**: used for fitting
    - **data**: beginning of experimental data
    - **range**  $E_1$   $E_2$ : experimental data used in the range from  $E_1$  to  $E_2$  (should be defined before the tag **ene**)
    - **system**  $i$ : define the system number
    - **ene**  $E$ : define incident energy
    - **tot**  $\sigma_t$   $\Delta\sigma_t$ : total cross section and error
    - **reac**  $\sigma_r$   $\Delta\sigma_r$ : reaction cross section and error
    - **s0**  $S_0$   $\Delta S_0$ : s-wave neutron strength function and error in the unit of  $10^{-4}$
    - **s1**  $S_1$   $\Delta S_1$ : p-wave neutron strength function and error in the unit of  $10^{-4}$
    - **r'**  $R'$   $\Delta R'$ : scattering length and error

- `scat i j etag`  
 $\theta_0 \sigma_0 \Delta\sigma_0$   
 $\theta_1 \sigma_1 \Delta\sigma_1$   
...  
**etag:**  
angular distribution of sum of the cross sections scattering to the levels between  $i$  and  $j$ .  $\theta_i$ ,  $\sigma_i$  and  $\Delta\sigma_i$  are angle, cross section, and error, respectively. The “**etag**” defines the string indicating end of data of the scattering cross sections
- `dend:` end of experimental data
- *others:*
  - `param idomp ( p0 p1 ... ):` parameters to calculate optical potentials; can be used for fitting.  $id_{omp}$  defines the OMP parametrisation to be used.
  - `output file:` file name of output. “-” means standard output. If this tag is not used, output file name is *input.out* for input file *input.in*.
  - `fit ( i j ... ):` parameters  $p_i$ ,  $p_j$ , ... to be used for fitting
  - `end:` end of input

## 5.2 Example of Input File

### 5.2.1 Example 1: Simple Calculation

This is an example of the calculation for  $^{238}\text{U} + n$  system coupled 5 levels using the optical potential of Ref. [10].

```

test1.in
1 # ----- system -----
2 system # 0 92- U -238
3 za 0 1 1.00866 92 238 238.051
4 lev 0.00 + 0.00000
5 lev 2.00 + 0.04491
6 lev 4.00 + 0.14841
7 lev 6.00 + 0.30721
8 lev 8.00 + 0.51830
9 # ----- calc -----
10 ene ( 1 2 )
11 angle 19 10 # 0 10 20 ... 180
12 calc ( 0 )
13 # -----
14 param om1 ( 51.32134 24 0.57 0.20 1.256 0.62
15 5.04567 12 0.40 0.001 1.26 0.58
16 6.0 1.12 0.5 1.3 0.198 0.057 )
17 end

```

The default values of  $j_{\max} = 40$ ,  $v_{\max} = 8$  and  $\text{mesh} = (20, 100)$  defined in the file `input.cpp` will be used. The following command is used to run the program with making the output file `test1.out`.

```
ccom test1.in
```

```
test1.out
```

```

1 # param om1 ( 51.3213 24 0.57 0.2 1.256 0.62 5.04567 12 0.4 0.001 ....
2 #-----
3 #   system 0 za 0 1 1.00866 92 238 238.051 nlev 5
4 #-----
5 # 0: Energy: 1
6 # cross section: 1 6929.04 2764.52 3788.55 324.37 50.6842 0.913304 0.00253666
7 # En, S0, S1, R': 1 0.655753 1.35069 5.49075
8 # En, transmission coefficient: 1 0.411152 0.625599 0.203342 0.12236 ...
9 # angular distribution
10 0 1832.45 1811.45 11.0627 9.6828 0.255607 0.000153664
11 10 1736.25 1714.76 11.6258 9.61084 0.249071 0.000154146
12 20 1476.17 1453 13.5475 9.38442 0.230463 0.000154378
13 30 1125.5 1099.09 17.2303 8.9776 0.202534 0.000151723
14 40 772.191 741.03 22.6279 8.36426 0.169133 0.000144273
15 50 486.817 450.289 28.8558 7.53802 0.134356 0.000132361
16 60 303.227 262.203 34.3925 6.52956 0.101771 0.000118423
17 70 217.042 173.816 37.7385 5.41309 0.0739624 0.00010579
18 80 197.881 155.473 38.0591 4.29596 0.0524101 9.78078e-05
19 90 206.71 167.969 35.4114 3.2915 0.0376043 9.81145e-05
20 100 210.89 177.846 30.5302 2.48474 0.0292099 0.000111437
21 110 193.006 166.61 24.4631 1.90717 0.0262382 0.000143388
22 120 152.561 132.662 18.3376 1.53377 0.0272633 0.000198327
23 130 101.484 86.877 13.272 1.30371 0.0307095 0.000276061
24 140 55.7336 44.3337 10.2117 1.15265 0.0351412 0.000369591
25 150 26.5626 15.9295 9.55475 1.03845 0.0394547 0.000465788
26 160 15.4528 3.68044 10.7807 0.948095 0.0429239 0.000548909
27 170 15.2476 1.78423 12.5303 0.887338 0.0451359 0.000604938
28 180 16.578 2.33906 13.3267 0.865736 0.0458924 0.000624663
29
30
31 # 1: Energy: 2
32 # cross section: 2 7189.9 3124.2 3495.34 443.079 111.428 15.4311 0.422913
33 # En, S0, S1, R': 2 0.551698 0.938799 5.274
34 # En, transmission coefficient: 2 0.48919 0.707341 0.382783 0.587136 ...
35 # angular distribution
36 0 3662.72 3617.48 29.0254 12.9286 3.24841 0.0326258
37 10 3298.52 3254.94 27.479 12.8424 3.22817 0.0320615
38 20 2387.83 2346.73 25.2643 12.6549 3.15435 0.0304472
39 30 1346.93 1304.44 26.9131 12.5536 2.99691 0.0280002
40 40 553.536 504.12 33.9179 12.7424 2.7311 0.0250345
41 50 154.558 96.1319 42.7727 13.2714 2.35967 0.0219209
42 60 70.1936 6.35024 47.999 13.9034 1.92191 0.019071
43 70 128.4 65.9575 46.7659 14.1764 1.48347 0.0169532
44 80 192.825 137.277 40.763 13.658 1.11104 0.0161454
45 90 208.609 161.217 34.3484 12.1807 0.8455 0.0174033

```

```

46 100 180.359 138.465 31.2792 9.9029 0.689625 0.0216317
47 110 132.553 92.1998 32.4688 7.23748 0.61741 0.0296131
48 120 85.8746 44.9148 35.6167 4.70617 0.595372 0.0415236
49 130 53.6057 13.3379 36.8889 2.72362 0.59875 0.0565461
50 140 45.7899 9.4653 34.2125 1.42488 0.61432 0.072917
51 150 68.3418 37.4305 29.4917 0.696124 0.635037 0.0884175
52 160 114.961 87.3026 26.5551 0.346579 0.655298 0.100983
53 170 162.735 135.15 26.5893 0.2159 0.669952 0.109097
54 180 183.057 154.887 27.1944 0.18808 0.675284 0.111894

```

The line 1 shows the used parameter set and the line 3 indicates  $Z$ ,  $A$ , mass and levels. The numerical values at the line 6 are the incident energy and the cross sections of total, reaction, elastic, 1st inelastic, 2nd inelastic, .... The transmission coefficients are shown at the line 8, where the values are lined in the order of incident energy,  $T_0$ ,  $T_1$ ,  $T_2$ , ...,  $T_i$ , .... The lines 10-28 show angular distributions of scattering cross sections. The columns show angle, sum of all scattering, elastic, inelastic to the 1st level, the 2nd level, .... The results for next energy begin from the 31st line.

To get the table of energy dependence of cross sections, following command may be useful.

```
awk -F : '/cross section/{print $2}' test1.out
```

### 5.2.2 Example 2: Fitting

An example of the input file for fitting of the OMP is shown. For the case of the reduced number of the coupled levels in the previous example, the energy dependence of the OMP is searched. The total cross section, the angular distributions of elastic and inelastic scattering, s- and p-wave strength functions and scattering radius are used for fitting.

`test2.in`

```

1 # ----- system -----
2 system # 0 92- U -238
3 za 0 1 1.00866 92 238 238.051
4 lev 0.00 + 0.00000
5 lev 2.00 + 0.04491
6 lev 4.00 + 0.14841
7 # ----- calc -----
8 ene ( 0.1 0.2 0.3 0.4 0.5 0.6 0.7 0.8 0.9
9      1 1.2 1.4 1.6 1.8 2 2.5 3 3.5 4 5 6 7 8 9
10     10 12 14 16 18 20 )
11 angle 0 2 # no angular distribution
12 calc ( 0 ) # energy dependence
13
14 ene ( 3.4 )
15 angle 91 2 # 0 2 4 6 ... 180
16 calc ( 0 ) # angular distribution
17 # -----

```



```

18 param om1 ( 51.32134 24 0.57 0.020 1.256 0.62
19             5.04567 12 0.40 0.001 1.26 0.58
20             6.0 1.12 0.5 1.3 0.198 0.057 )
21 fit ( 0 2 6 8 )
22 # ----- data -----
23 data
24 system 0
25     ene 0.0001 s0 1.2 0.1 s1 1.7 0.3 r' 9.6 0.1
26     ene 0.1 tot 11814.2 49.1913
27     ene 0.2 tot 10696.8 27.5866
28     ene 0.5 tot 8387.69 12.3136
29     ene 1   tot 7065.96 6.94642
30     ene 1.4 tot 6936.05 6.47929
31     ene 1.8 tot 7160.01 7.76091
32     ene 2   tot 7307.15 6.62315
33     ene 2.5 tot 7630.9 5.56785
34     ene 3   tot 7858.7 6.65072
35     ene 3.5 tot 7968.68 7.88627
36     ene 4   tot 7994.81 8.59415
37     ene 4.5 tot 7877.03 9.00727
38     ene 5   tot 7664.56 7.2575
39     ene 7   tot 6674.51 4.80028
40     ene 10  tot 5881.1 4.97937
41     ene 14  tot 5853.65 5.75409
42     ene 18  tot 6147.31 5.01827
43     ene 20  tot 6213.03 5.29372
44
45 scat 0 0 -----
46     2.0100000e+01 2.8672000e+03 1.7490000e+02
47     3.0100000e+01 1.5405000e+03 9.5500000e+01
48     4.0200000e+01 3.1000000e+02 2.3900000e+01
49     4.5200000e+01 1.4470000e+02 1.1700000e+01
50     5.0200000e+01 2.7700000e+01 3.4000000e+00
51     6.0200000e+01 3.7100000e+01 5.4000000e+00
52     6.5200000e+01 8.4600000e+01 8.2000000e+00
53     8.0200000e+01 1.1390000e+02 1.2200000e+01
54     8.5200000e+01 1.0560000e+02 9.8000000e+00
55     9.0200000e+01 8.7400000e+01 8.7000000e+00
56     1.0020000e+02 3.3400000e+01 4.2000000e+00
57     1.0520000e+02 2.7000000e+01 3.6000000e+00
58     1.1020000e+02 1.0700000e+01 1.9000000e+00
59     1.2020000e+02 2.5000000e+00 7.0000000e-01
60     1.3020000e+02 5.8000000e+00 1.5000000e+00
61     1.4020000e+02 1.4500000e+01 2.7000000e+00
62     1.5010000e+02 2.4600000e+01 3.7000000e+00
63     1.6010000e+02 3.3900000e+01 4.4000000e+00
64     -----
65
66 scat 1 1 -----
67     4.0200000e+01 8.2200000e+01 9.7000000e+00
68     4.5200000e+01 6.1900000e+01 6.6000000e+00
69     5.0200000e+01 4.5900000e+01 6.2000000e+00
70     6.0200000e+01 3.6500000e+01 5.9000000e+00

```

```

71 6.520000e+01 4.390000e+01 5.100000e+00
72 8.020000e+01 5.130000e+01 7.400000e+00
73 8.520000e+01 5.030000e+01 5.800000e+00
74 9.020000e+01 5.290000e+01 5.800000e+00
75 1.002000e+02 3.070000e+01 4.100000e+00
76 1.052000e+02 2.040000e+01 3.100000e+00
77 1.102000e+02 1.720000e+01 2.800000e+00
78 1.202000e+02 1.040000e+01 2.100000e+00
79 1.302000e+02 3.460000e+01 4.600000e+00
80 1.402000e+02 4.300000e+01 5.200000e+00
81 1.501000e+02 2.840000e+01 3.900000e+00
82 1.601000e+02 1.990000e+01 3.100000e+00
83 -----
84
85 scat 2 2 -----
86 4.020000e+01 2.190000e+01 4.500000e+00
87 4.520000e+01 2.220000e+01 3.500000e+00
88 5.020000e+01 1.580000e+01 3.500000e+00
89 6.020000e+01 1.220000e+01 3.300000e+00
90 6.520000e+01 9.900000e+00 2.100000e+00
91 8.020000e+01 9.000000e+00 2.800000e+00
92 8.520000e+01 9.900000e+00 2.200000e+00
93 9.020000e+01 1.010000e+01 2.600000e+00
94 1.052000e+02 9.700000e+00 2.200000e+00
95 1.102000e+02 8.400000e+00 2.000000e+00
96 1.202000e+02 6.800000e+00 1.800000e+00
97 1.302000e+02 6.600000e+00 1.700000e+00
98 1.402000e+02 4.800000e+00 1.500000e+00
99 1.501000e+02 3.900000e+00 1.300000e+00
100 1.601000e+02 3.100000e+00 1.100000e+00
101 -----
102 dend
103 end

```

The parameters to be fitted is shown at the line 21. In the *om1* parametrisation,  $p[0]$  and  $p[2]$  are constant and linear terms of energy dependence of real potential, respectively, and  $p[6]$  and  $p[8]$  for imaginary potential. The fitting data are shown from the line 23. The line 24 shows the system to be calculated; it is defined at lines 2-6. At the line 25 s- and p-wave strength functions and scattering radius are written. The lines 26-43 show the total cross sections. The data for angular distribution for the elastic scattering are written at the lines 45-64. The line 66-83 and line 85-101 are for the inelastic scattering to 1st and 2nd excited levels, respectively.

`test2.out`

```

1 # param om1 ( 51.3213 24 0.57 0.02 1.256 0.62 5.04567 12 0.4 0.001 1.26 0.58 ...
2 # fit ( 0 2 6 8 )
3 # fitting started
4 # iter: 0 x = 51.3213400 0.5700000 5.0456700 0.4000000 |f(x)| = 142.68
5 # iter: 1 x = 52.2212321 0.5886197 5.1215875 0.5132671 |f(x)| = 72.1206
6 # iter: 2 x = 52.0315561 0.5575361 5.1235479 0.4939006 |f(x)| = 66.3312
7 # iter: 3 x = 52.0803911 0.5601253 4.9158329 0.5312031 |f(x)| = 66.0273
8 # iter: 4 x = 52.0769483 0.5571478 4.8741173 0.5409622 |f(x)| = 65.9834

```

```

9 # iter: 5 x = 52.0802561 0.5565554 4.8467365 0.5468730 |f(x)| = 65.9735
10 # iter: 6 x = 52.0810143 0.5561466 4.8358709 0.5492956 |f(x)| = 65.9715
11 # iter: 7 x = 52.0814727 0.5559865 4.8307129 0.5504341 |f(x)| = 65.9708
12 # iter: 8 x = 52.0816644 0.5559102 4.8284067 0.5509443 |f(x)| = 65.9706
13 # iter: 9 x = 52.0817541 0.5558761 4.8273536 0.5511770 |f(x)| = 65.9706
14 # iter: 10 x = 52.0817945 0.5558606 4.8268760 0.5512826 |f(x)| = 65.9705
15 # p(0) = 52.08179 +/- 0.00892 | 1.000 0.409 -0.709 0.515
16 # p(1) = 0.55586 +/- 0.00243 | 0.409 1.000 -0.049 -0.029
17 # p(2) = 4.82688 +/- 0.03573 | -0.709 -0.049 1.000 -0.836
18 # p(3) = 0.55128 +/- 0.00749 | 0.515 -0.029 -0.836 1.000
19 # status = success
20 # param om1 ( 52.0818 24 0.555861 0.02 1.256 0.62 4.82688 12 0.551283 0.001 ...
21 # -----
22 # system 0 za 0 1 1.00866 92 238 238.051 nlev 3
23 # -----
24 # 0: Energy: 0.1
25 # cross section: 0.1 11975.1 3167.02 8796.25 11.865 0
26 # En, S0, S1, R': 0.1 0.639172 2.63677 8.3665
27 # En, transmission coefficient: 0.1 0.12673 0.11523 0.00191502 5.78636e-05 ...
28
29
30 # 1: Energy: 0.2
31 # cross section: 0.2 11018.1 3408.16 7563.34 46.4354 0.155721
32 # En, S0, S1, R': 0.2 0.632752 2.65966 7.75804
33 # En, transmission coefficient: 0.2 0.177423 0.269376 0.009617 0.000656405 ...
34

```

The lines 4-14 show results for the each iterations. The last column shows  $\sqrt{\sum_i (y_i - f_i)^2 / \sigma_i^2}$ , where  $y_i$  and  $\sigma_i$  are the experimental data and its error for fitting, respectively, and  $f_i$  is the calculated value. At the line 15-18, the final results with the estimation of errors and correlations of the fitted parameters are shown. The line 20 shows the new parameter set obtained by fitting. From the line 21, the final calculated cross sections are shown. The result of fitting is shown in Fig. 1.

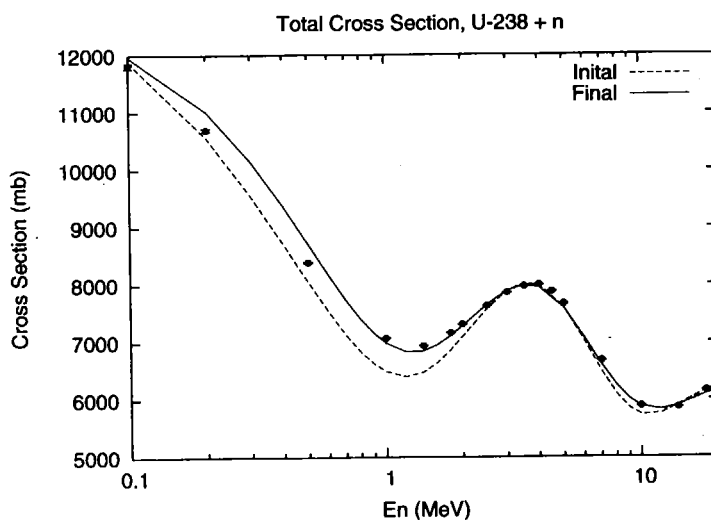


Fig. 1: Total cross sections calculated by initial and final parameters in the fitting.

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# 国際単位系 (SI) と換算表

表1 SI基本単位および補助単位

量	名称	記号
長さ	メートル	m
質量	キログラム	kg
時間	秒	s
電流	アンペア	A
熱力学温度	ケルビン	K
物質質量	モル	mol
光度	カンデラ	cd
平面角	ラジアン	rad
立体角	ステラジアン	sr

表2 SIと併用される単位

名称	記号
分, 時, 日	min, h, d
度, 分, 秒	°, ', "
リットル	l, L
トン	t
電子ボルト	eV
原子質量単位	u

$1 \text{ eV} = 1.60218 \times 10^{-19} \text{ J}$   
 $1 \text{ u} = 1.66054 \times 10^{-27} \text{ kg}$

表5 SI接頭語

倍数	接頭語	記号
$10^{18}$	エクサ	E
$10^{15}$	ペタ	P
$10^{12}$	テラ	T
$10^9$	ギガ	G
$10^6$	メガ	M
$10^3$	キロ	k
$10^2$	ヘクト	h
$10^1$	デカ	da
$10^{-1}$	デシ	d
$10^{-2}$	センチ	c
$10^{-3}$	ミリ	m
$10^{-6}$	マイクロ	$\mu$
$10^{-9}$	ナノ	n
$10^{-12}$	ピコ	p
$10^{-15}$	フェムト	f
$10^{-18}$	アト	a

表3 固有の名称をもつSI組立単位

量	名称	記号	他のSI単位による表現
周波数	ヘルツ	Hz	$\text{s}^{-1}$
力	ニュートン	N	$\text{m}\cdot\text{kg}/\text{s}^2$
圧力, 応力	パスカル	Pa	$\text{N}/\text{m}^2$
エネルギー, 仕事, 熱量	ジュール	J	$\text{N}\cdot\text{m}$
工率, 放射束	ワット	W	$\text{J}/\text{s}$
電気量, 電荷	クーロン	C	$\text{A}\cdot\text{s}$
電位, 電圧, 起電力	ボルト	V	$\text{W}/\text{A}$
静電容量	ファラド	F	$\text{C}/\text{V}$
電気抵抗	オーム	$\Omega$	$\text{V}/\text{A}$
コンダクタンス	ジーメン	S	$\text{A}/\text{V}$
磁束	ウェーバ	Wb	$\text{V}\cdot\text{s}$
磁束密度	テスラ	T	$\text{Wb}/\text{m}^2$
インダクタンス	ヘンリー	H	$\text{Wb}/\text{A}$
セルシウス温度	セルシウス度	$^{\circ}\text{C}$	
光束	ルーメン	lm	$\text{cd}\cdot\text{sr}$
照射度	ルクス	lx	$\text{lm}/\text{m}^2$
放射能	ベクレル	Bq	$\text{s}^{-1}$
吸収線量	グレイ	Gy	$\text{J}/\text{kg}$
線量当量	シーベルト	Sv	$\text{J}/\text{kg}$

表4 SIと共に暫定的に維持される単位

名称	記号
オングストローム	$\text{\AA}$
バーン	b
バル	bar
ガリ	Gal
キュリー	Ci
レントゲン	R
ラド	rad
レム	rem

$1 \text{ \AA} = 0.1 \text{ nm} = 10^{-10} \text{ m}$   
 $1 \text{ b} = 100 \text{ fm} = 10^{-28} \text{ m}^2$   
 $1 \text{ bar} = 0.1 \text{ MPa} = 10^5 \text{ Pa}$   
 $1 \text{ Gal} = 1 \text{ cm}/\text{s}^2 = 10^{-2} \text{ m}/\text{s}^2$   
 $1 \text{ Ci} = 3.7 \times 10^{10} \text{ Bq}$   
 $1 \text{ R} = 2.58 \times 10^{-4} \text{ C}/\text{kg}$   
 $1 \text{ rad} = 1 \text{ cGy} = 10^{-2} \text{ Gy}$   
 $1 \text{ rem} = 1 \text{ cSv} = 10^{-2} \text{ Sv}$

(注)

- 表1-5は「国際単位系」第5版、国際度量衡局 1985年刊行による。ただし、1 eV および 1 uの値はCODATAの1986年推奨値によった。
- 表4には海里、ノット、アール、ヘクトールも含まれているが日常の単位なのでここでは省略した。
- barは、JISでは流体の圧力を表わす場合に限り表2のカテゴリーに分類されている。
- EC閣僚理事会指令ではbar, barnおよび「血圧の単位」mmHgを表2のカテゴリーに入れている。

## 換算表

力	N (=10 <sup>5</sup> dyn)	kgf	lbf
	1	0.101972	0.224809
	9.80665	1	2.20462
	4.44822	0.453592	1

粘 度  $1 \text{ Pa}\cdot\text{s} (\text{N}\cdot\text{s}/\text{m}^2) = 10 \text{ P} (\text{ポアズ}) (\text{g}/(\text{cm}\cdot\text{s}))$

動粘度  $1 \text{ m}^2/\text{s} = 10^4 \text{ St} (\text{ストークス}) (\text{cm}^2/\text{s})$

圧	MPa (=10 bar)	kgf/cm <sup>2</sup>	atm	mmHg (Torr)	lbf/in <sup>2</sup> (psi)
	1	10.1972	9.86923	$7.50062 \times 10^3$	145.038
力	0.0980665	1	0.967841	735.559	14.2233
	0.101325	1.03323	1	760	14.6959
	$1.33322 \times 10^{-4}$	$1.35951 \times 10^{-3}$	$1.31579 \times 10^{-3}$	1	$1.93368 \times 10^{-2}$
	$6.89476 \times 10^{-3}$	$7.03070 \times 10^{-2}$	$6.80460 \times 10^{-2}$	51.7149	1

エネルギー・仕事・熱量	J (=10 <sup>7</sup> erg)	kgf·m	kW·h	cal (計量法)	Btu	ft·lbf	eV
	1	0.101972	$2.77778 \times 10^{-7}$	0.238889	$9.47813 \times 10^{-4}$	0.737562	$6.24150 \times 10^{18}$
	9.80665	1	$2.72407 \times 10^{-6}$	2.34270	$9.29487 \times 10^{-3}$	7.23301	$6.12082 \times 10^{19}$
	$3.6 \times 10^6$	$3.67098 \times 10^5$	1	$8.59999 \times 10^5$	3412.13	$2.65522 \times 10^6$	$2.24694 \times 10^{25}$
	4.18605	0.426858	$1.16279 \times 10^{-6}$	1	$3.96759 \times 10^{-3}$	3.08747	$2.61272 \times 10^{19}$
	1055.06	107.586	$2.93072 \times 10^{-4}$	252.042	1	778.172	$6.58515 \times 10^{21}$
	1.35582	0.138255	$3.76616 \times 10^{-7}$	0.323890	$1.28506 \times 10^{-3}$	1	$8.46233 \times 10^{18}$
	$1.60218 \times 10^{-19}$	$1.63377 \times 10^{-20}$	$4.45050 \times 10^{-26}$	$3.82743 \times 10^{-20}$	$1.51857 \times 10^{-22}$	$1.18171 \times 10^{-19}$	1

$1 \text{ cal} = 4.18605 \text{ J} (\text{計量法})$   
 $= 4.184 \text{ J} (\text{熱化学})$   
 $= 4.1855 \text{ J} (15^{\circ}\text{C})$   
 $= 4.1868 \text{ J} (\text{国際蒸気表})$   
 仕事率  $1 \text{ PS} (\text{仏馬力})$   
 $= 75 \text{ kgf}\cdot\text{m}/\text{s}$   
 $= 735.499 \text{ W}$

放射能	Bq	Ci
	1	$2.70270 \times 10^{-11}$
	$3.7 \times 10^{10}$	1

吸収線量	Gy	rad
	1	100
	0.01	1

照射線量	C/kg	R
	1	3876
	$2.58 \times 10^{-4}$	1

線量当量	Sv	rem
	1	100
	0.01	1

Program CCOM - Coupled-channels Optical Model Calculation with Automatic Parameter Search -

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