An optimal method for high order mixed derivatives of bivariate functions

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The problem of optimal recovering high-order mixed derivatives of bivariate functions with finite smoothness is studied. Based on the truncation method, an algorithm for numerical differentiation is constructed, which is order-optimal both in the sense of accuracy and in terms of the amount of involved Galerkin information. Numerical examples are provided to illustrate the fact that our approach can be implemented successfully.

1. Description of the problem

Currently, many research activities on the problem of stable numerical differentiation have been taking place due to the importance of this tool in such areas of science and technology as finance, mathematical physics, image processing, analytical chemistry, viscous elastic mechanics, reliability analysis, pattern recognition, and many others. Among these investigations, we highlight [4], which is the first publication on numerical differentiation in terms of the theory of ill-posed problems. Further research [4] has been continued in numerous publications on numerical differentiation (see, for example, [21], [34], [7], [8], [1], [9], [17], [37], [38], [12], [26]), covering different classes of the functions and the types of proposed methods. Despite the abundance of works on this topic, the problem for recovery of high-order derivatives was considered only in a few publications, among which we note [5], [22], [35], [20], [17] and [26]. In particular, the results of [26] have opened perspective for further investigation of numerical methods for recovery of high-order derivatives. Namely, the main criterion of the method's efficiency has been taken as its ability to achieve the optimal order of accuracy by using a minimal amount of discrete information. Note that these aspects of numerical differentiation remain still insufficiently studied. The present paper continues the research of [26], [24], [25] and proposes a numerical method for recovering the high-order mixed derivatives of smooth bivariate functions. The method is not only stable to small perturbations of the input data, but is also optimal in terms of accuracy and quantity of involved Fourier-Legendre coefficients, and also has a simple numerical implementation. Note that the approach to numerical differentiation proposed below can be widely used in solving various practical problems. For example, the expected results can be applied to the analysis of well-known multidimensional Savitzky-Golay method [28], [10].

The article is organized as follows. In Section 1., the problem statement for optimizing numerical differentiation methods in the sense of the minimal radius of Galerkin information is given. Sections 2. and 3. describe a modification of the truncation method and establish its accuracy estimates in quadratic and uniform metrics, respectively. Section 4. is devoted to finding order estimates for the minimal radius of Galerkin information while establishing the optimality (in the power scale) of the method under consideration. In Section 5., numerical experiments will be provided to confirm the effectiveness of the proposed method.

For the further presentation of the subject matter of this paper, we need the following notation and concepts. Let $\{\varphi_k(t)\}_{k=0}^{\infty}$ be the system of Legendre polynomials orthonormal on [-1,1] as

$$\varphi_k(t) = \sqrt{k+1/2}(2^k k!)^{-1} \frac{d^k}{dt^k} [(t^2 - 1)^k], \quad k = 0, 1, 2, \dots$$

 $^{^{\}dagger}$ Key words. Numerical differentiation, Legendre polynomials, truncation method, minimal radius of Galerkin information.

By $L_2 = L_2(Q)$ we mean space of square-summable on $Q = [-1, 1]^2$ functions $f(t, \tau)$ with inner product

$$\langle f, g \rangle = \int_{-1}^{1} \int_{-1}^{1} f(t, \tau) g(t, \tau) d\tau dt$$

and norm generated by it

$$||f||_{L_2}^2 = \sum_{k,j=0}^{\infty} |\langle f, \varphi_{k,j} \rangle|^2 < \infty,$$

where

$$\langle f, \varphi_{k,j} \rangle = \int_{-1}^{1} \int_{-1}^{1} f(t, \tau) \varphi_k(t) \varphi_j(\tau) d\tau dt, \quad k, j = 0, 1, 2, \dots,$$

are Fourier-Legendre coefficients of f. Moreover, let C = C(Q) be the space of continuous bivariate functions on Q equipped with standard uniform norm and ℓ_p , $1 \le p \le \infty$, be the space of numerical sequences $\overline{x} = \{x_{k,j}\}_{k,j \in \mathbb{N}_0}$, $\mathbb{N}_0 = \{0\} \bigcup \mathbb{N}$, such that the corresponding relation

$$\|\overline{x}\|_{\ell_p} := \begin{cases} \left(\sum_{k,j \in \mathbb{N}_0} |x_{k,j}|^p\right)^{\frac{1}{p}} < \infty, & 1 \le p < \infty, \\ \sup_{k,j \in \mathbb{N}_0} |x_{k,j}| < \infty, & p = \infty, \end{cases}$$

$$(1.1)$$

is fulfilled.

We introduce a space of functions

$$L_{s,2}^{\mu}(Q) = \{ f \in L_2(Q) : \|f\|_{s,\mu}^s := \sum_{k,j=0}^{\infty} (\underline{k \cdot j})^{s\mu} |\langle f, \varphi_{k,j} \rangle|^s < \infty \},$$

where $\mu > 0, 1 \le s < \infty, \underline{k} = \max\{1, k\}, k = 0, 1, 2, \dots$

Hereinafter we will use the same notations both for the space and for a unite ball from this space, which we call a class of functions and denote as $L_{s,2}^{\mu} = L_{s,2}^{\mu}(Q) = \{f \in L_{s,2}^{\mu} \colon \|f\|_{s,\mu} \le 1\}$. What exactly is meant by $L_{s,2}^{\mu}$, space or class, will be clear depending on the context in each case. It should be noted that $L_{s,2}^{\mu}$ is a generalization of the class of bivariate functions with dominating mixed partial derivatives.

We represent a function $f(t,\tau)$ from $L_{s,2}^{\mu}$, $\mu \geq 2r + 1/2 - 1/s$, as

$$f(t,\tau) = \sum_{k,j=0}^{\infty} \langle f, \varphi_{k,j} \rangle \varphi_k(t) \varphi_j(\tau),$$

and by its mixed derivative $f^{(r,r)}$, $r=1,2,\ldots$, we mean the following series

$$f^{(r,r)}(t,\tau) = \sum_{k,j=r}^{\infty} \langle f, \varphi_{k,j} \rangle \varphi_k^{(r)}(t) \varphi_j^{(r)}(\tau).$$
(1.2)

Assume that instead of the exact values of the Fourier-Legendre coefficients $\langle f, \varphi_{k,j} \rangle$ only some their perturbations are known with the error level δ in the metrics of ℓ_p , $1 \leq p \leq \infty$. More accurately, we assume that there is a sequence of numbers $\overline{f^{\delta}} = \{\langle f^{\delta}, \varphi_{k,j} \rangle\}_{k,j \in \mathbb{N}_0}$ such that for $\overline{\xi} = \{\xi_{k,j}\}_{k,j \in \mathbb{N}_0}$, where $\xi_{k,j} = \langle f - f^{\delta}, \varphi_{k,j} \rangle$, and for some $1 \leq p \leq \infty$ the relation

$$\|\overline{\xi}\|_{\ell_p} \le \delta, \quad 0 < \delta < 1, \tag{1.3}$$

is true.

The research of this work is devoted to the optimization of methods for recovering the derivative (1.2) of functions from $L_{s,2}^{\mu}$. Further, we give a strict statement of the problem to be studied. In the coordinate plane $[r,\infty)\times[r,\infty)$ we take an arbitrary bounded domain Ω . By card(Ω) we mean the number of points that make up Ω and by the information vector

 $G(\Omega, \overline{f}^{\delta}) \in \mathbb{R}^N$, $\operatorname{card}(\Omega) = N$, we take the set of perturbed values of Fourier-Legendre coefficients

 $\{\langle f^{\delta}, \varphi_{k,j} \rangle\}_{(k,j) \in \Omega}$. Let $X = L_2(Q)$ or X = C(Q). By numerical differentiation algorithm we mean any interpolation vector $G(\Omega, \overline{f}^{\delta})$ an element mapping $\psi^{(r,r)} = \psi^{(r,r)}(\Omega)$ that corresponds to the information vector $G(\Omega, \overline{f}^{\delta})$ an element $\psi^{(r,r)}(G(\Omega,\overline{f}^{\delta})) \in X$, which is taken as an approximation to the derivative (1.2) of function f from the class $L^{\mu}_{s,2}$. We denote by $\Psi(\Omega)$ the set of all algorithms $\psi^{(r,r)}(\Omega): \mathbb{R}^N \to X$, that use the same information vector $G(\Omega, \overline{f}^{\delta})$.

We do not require, generally speaking, either linearity or even stability for algorithms from $\Psi(\Omega)$. The only condition for these algorithms is to use an input information in the form of perturbed values of the Fourier-Legendre coefficients with indices from the domain Ω of the coordinate plane. Such a general understanding of the algorithm is explained by the desire to consider the widest range of possible methods of numerical differentiation.

The error of the algorithm $\psi^{(r,r)}$ on the class $L_{s,2}^{\mu}$ is determined by the quantity

$$\varepsilon_{\delta}(L_{s,2}^{\mu}, \psi^{(r,r)}(\Omega), X, \ell_{p}) = \sup_{\substack{f \in L_{s,2}^{\mu}, \\ \|f\|_{s,\mu} \leq 1}} \sup_{\overline{f}^{\delta}: (1.3)} \|f^{(r,r)} - \psi^{(r,r)}(G(\Omega, \overline{f}^{\delta}))\|_{X}.$$

The minimal radius of the Galerkin information for the problem of numerical differentiation on the class $L_{s,2}^{\mu}$ is given by

$$R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu},X,\ell_p) = \inf_{\Omega: \operatorname{card}(\Omega) \leq N} \inf_{\psi^{(r,r)} \in \Psi(\Omega)} \varepsilon_{\delta}(L_{s,2}^{\mu},\psi^{(r,r)}(\Omega),X,\ell_p).$$

The quantity $R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu},X,\ell_p)$ describes the minimal possible accuracy in the metric of space X, which can be achieved by numerical differentiation of arbitrary function $f \in L_{s,2}^{\mu}$, while using not more than N values of its Fourier-Legendre coefficients that are δ -perturbed in the ℓ_p metric. Note that the minimal radius of Galerkin information in the problem of recovering the first partial derivatives was studied in [31], the mixed derivatives $f^{(2,2)}$ in [25] and for other types of ill-posed problems, similar studies were previously carried out in [16, 19]. It should be added that the minimal radius characterizes the information complexity of the considered problem and is traditionally studied within the framework of the IBC-theory (Information-Based Complexity Theory), the foundations of which are laid in monographs [33] and [32].

The goal of our research is to find order-optimal estimates (in the power scale) for $R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu},C,\ell_p)$ and $R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu},L_2,\ell_p)$. Finally, we introduce the symbolic notation for inequality and equality in order. For two

positive quantities a and b, we write $a \leq b$ if there exists a constant c > 0, such that $a \leq cb$. We will write $a \times b$ if $a \leq b$ and $b \leq a$.

Truncation method. Error estimate in the metric of L_2 2.

Many approaches to numerical differentiation have been developed so far (see e.g. [3], [2], [23], [22] and also see [27] and the references therein). However, most known methods have their drawbacks. So, in particular, some of them guarantee satisfactory accuracy only in the case of exactly given input data. Other methods are quite complex for numerical implementation.

Here we mean the representation of the problem to be solved in the form of an integral equation (for example, in the framework of the Tikhonov method and its various variations), as well as difficult-to-implement rules for determining the regularization parameters (see [22]). Thus, the problem of constructing efficient methods for numerical differentiation remains open. Recently in [26], [31] a concise numerical method called the truncation method has been proposed as a simple and optimal (in the sense of minimal radius of Galerkin information) approach to numerically differentiating bivariate functions. The essence of this method is to replace the Fourier series (1.2) with a finite Fourier sum over perturbed data $\langle f^{\delta}, \varphi_{k,j} \rangle$. In the truncation method, to ensure the stability of the approximation and achieve the required order accuracy, it is necessary to correctly choose a discretization parameter, which here serves as the regularization parameter. More strictly, the regularization process in the method under consideration consists of matching the discretization parameter with the perturbation level δ of the input data, the smoothness parameter μ , the order of the derivative r, and the values of p, s. The simplicity of implementation is the main advantage of this method. Moreover, in Section 4. it will be established that the proposed version (2.1) of the truncation method is optimal in terms of the minimal radius of Galerkin information. In other words, this method achieves the best order of accuracy while using the smallest possible amount of perturbed input data $\langle f^{\delta}, \varphi_{k,j} \rangle$.

In the case of an arbitrary bounded domain Ω of the coordinate plane $[r, \infty) \times [r, \infty)$, the truncation method for differentiating bivariate functions has the form

$$\mathcal{D}_{\Omega}^{(r,r)} f^{\delta}(t,\tau) = \sum_{(k,j) \in \Omega} \langle f^{\delta}, \varphi_{k,j} \rangle \varphi_k^{(r)}(t) \varphi_j^{(r)}(\tau).$$

To increase the efficiency of the approach under study, we take a hyperbolic cross as the domain Ω of the following form

$$\Omega = \Gamma_n := \{(k, j) : k \cdot j \le rn - 1, \quad k, j = r, \dots, n - 1\}, \quad \operatorname{card}(\Gamma_n) \times n \ln n.$$

Then the proposed version of the truncation method can be written

$$\mathcal{D}_{n}^{(r,r)} f^{\delta}(t,\tau) = \sum_{k,j \ge r, \ kj \le rn-1} \langle f^{\delta}, \varphi_{k,j} \rangle \varphi_{k}^{(r)}(t) \varphi_{j}^{(r)}(\tau). \tag{2.1}$$

We note that the idea of a hyperbolic cross for the problem of numerical differentiation was used earlier in the papers [24], [31], [26], [25] (for more details about the usage of a hyperbolic cross in solving the other ill-posed problems see [18], [6], [15], [14]).

Remark 2.1 In [26], the problem of recovering the derivatives $f^{(r,r)}$ of periodic functions was considered, when the perturbed values of the Fourier coefficients w.r.t. the trigonometric system are taken as input information. Unfortunately, it is impossible to automatically transfer the results from the periodic case to the non-periodic one. In particular, this is because the best approximation accuracy for derivatives of non-periodic functions has a worse order than the best accuracy for approximation of derivatives of periodic functions (cf. [26]). Therefore, in the non-periodic case, to construct optimal methods for numerical differentiation, a modification of the previous methodology and the development of new techniques are required. At the moment, we have already constructed optimal methods for recovering the derivatives $f^{(1,1)}$ [24], $f^{(2,2)}$ [25], $f^{(1,0)}$ [31] in the non-periodic case. The results of [24], [31], and [25] together with the results of this work, create the ground and prospects for the development of optimal methods for recovering derivatives of any order for non-periodic functions of any number of variables.

Remark 2.2 If in (2.1) we put r = 0, then the problem of numerical differentiation is transformed into the problem of numerical summation. Earlier, to solve the problem of numerical

summation of univariate functions, the truncation method was proposed in [11]. Further, this approach was extended to the case of bivariate functions in [29], [30]. It should be noted that in the last two works, the idea of a hyperbolic cross was successfully implemented, which made it possible to significantly reduce the computational resources without loss of accuracy. In the framework of the present investigation, we are going to achieve the same effect for the problem of numerical differentiation.

Let us investigate the approximation properties of the method (2.1). To this end, we write the error of the method (2.1) as

$$f^{(r,r)}(t,\tau) - \mathcal{D}_n^{(r,r)} f^{\delta}(t,\tau) = \left(f^{(r,r)}(t,\tau) - \mathcal{D}_n^{(r,r)} f(t,\tau) \right) + \left(\mathcal{D}_n^{(r,r)} f(t,\tau) - \mathcal{D}_n^{(r,r)} f^{\delta}(t,\tau) \right). \tag{2.2}$$

For the first difference on the right-hand side of (2.2), the representation

$$f^{(r,r)}(t,\tau) - \mathcal{D}_n^{(r,r)}f(t,\tau) = \Delta_1(t,\tau) + \Delta_2(t,\tau) + \Delta_3(t,\tau)$$
 (2.3)

holds, where

$$\Delta_1(t,\tau) = \sum_{k=n+1}^{\infty} \sum_{j=r}^{\infty} \langle f, \varphi_{k,j} \rangle \varphi_k^{(r)}(t) \varphi_j^{(r)}(\tau), \tag{2.4}$$

$$\Delta_2(t,\tau) = \sum_{k=r}^n \sum_{j=n+1}^\infty \langle f, \varphi_{k,j} \rangle \varphi_k^{(r)}(t) \varphi_j^{(r)}(\tau), \tag{2.5}$$

$$\triangle_3(t,\tau) = \sum_{k=r}^n \sum_{j=\frac{r_n}{k}}^n \langle f, \varphi_{k,j} \rangle \varphi_k^{(r)}(t) \varphi_j^{(r)}(\tau). \tag{2.6}$$

For our calculations, we need the following formula (see Lemma 18 [13])

$$\varphi_k'(t) = 2\sqrt{k+1/2} \sum_{l=0}^{k-1} \sqrt{l+1/2} \varphi_l(t), \quad k \in \mathbb{N},$$
(2.7)

where in aggregate $\sum_{l=0}^{k-1} \sqrt{l+1/2} \varphi_l(t)$ the summation is extended over only those terms for which k+l is odd.

In the sequel, we adopt the convention that c denotes a generic positive coefficient, which can vary from inequality to inequality and may only depend on basic parameters such as μ, r, p, s and others which may appear below.

Let us estimate the error of the method (2.1) in the metric of L_2 . An upper bound for difference (2.3) is contained in the following statement.

Lemma 2.3 Let $f \in L_{s,2}^{\mu}$, $1 \le s < \infty$, $\mu > 2r + 1/2 - 1/s$. Then it holds

$$||f^{(r,r)} - \mathcal{D}_n^{(r,r)}f||_{L_2} \le c||f||_{s,\mu}n^{-\mu + 2r + 1/2 - 1/s}\ln^{3/2 - 1/s}n.$$

Proof. Using the formula (2.7), from (2.4) we have

$$\triangle_1(t,\tau) = 4^r \sum_{k=n+1}^{\infty} \sum_{j=r}^{\infty} \sqrt{k+1/2} \sqrt{j+1/2} \langle f, \varphi_{k,j} \rangle$$

$$\times \sum_{l_1=r-1}^{k-1} {}^{*}(l_1+1/2) \sum_{l_2=r-2}^{l_1-1} {}^{*}(l_2+1/2) \dots \sum_{l_{r-1}=1}^{l_{r-2}-1} {}^{*}(l_{r-1}+1/2) \sum_{l_r=0}^{l_{r-1}-1} {}^{*}\sqrt{l_r+1/2} \varphi_{l_r}(t)$$

$$\times \sum_{m_1=r-1}^{j-1} {}^{*}(m_1+1/2) \sum_{m_2=r-2}^{m_1-1} {}^{*}(m_2+1/2) \dots \sum_{m_{r-1}=1}^{m_{r-2}-1} {}^{*}(m_{r-1}+1/2) \sum_{m_r=0}^{m_{r-1}-1} {}^{*}\sqrt{m_r+1/2} \varphi_{m_r}(\tau).$$

We note that in the representation Δ_1 only those terms take place for which all indexes $l_1+k, l_2+l_1, ..., l_r+l_{r-1}, m_1+j, m_2+m_1, ..., m_r+m_{r-1}$ are odd. Such rule is valid also for other terms, namely Δ_2, Δ_3 and $\mathcal{D}_n^{(r,r)}f - \mathcal{D}_n^{(r,r)}f^{\delta}$, appearing in the error representation (see (2.2)-(2.6)). In the following, for simplicity, we will omit the symbol "*" when denoting such summation operations, while taking into account this rule in the calculations.

Further, we change the order of summation and get

$$\triangle_1(t,\tau) = \triangle_{11}(t,\tau) + \triangle_{12}(t,\tau),$$

where

$$\Delta_{11}(t,\tau) = 4^r \sum_{l_r=0}^{n-r+1} \sqrt{l_r + 1/2} \, \varphi_{l_r}(t) \sum_{m_r=0}^{\infty} \sqrt{m_r + 1/2} \, \varphi_{m_r}(\tau)$$

$$\times \sum_{k=n+1}^{\infty} \sum_{j=m_r+r}^{\infty} \sqrt{k + 1/2} \sqrt{j + 1/2} \, \langle f, \varphi_{k,j} \rangle B_{k,j}^r, \tag{2.8}$$

$$\triangle_{12}(t,\tau) = 4^r \sum_{l_r = n - r + 2}^{\infty} \sqrt{l_r + 1/2} \, \varphi_{l_r}(t) \sum_{m_r = 0}^{\infty} \sqrt{m_r + 1/2} \, \varphi_{m_r}(\tau)$$

$$\times \sum_{k=l_r+r}^{\infty} \sum_{j=m_r+r}^{\infty} \sqrt{k+1/2} \sqrt{j+1/2} \langle f, \varphi_{k,j} \rangle B_{k,j}^r$$
 (2.9)

and

$$B_{k,j}^{r} := \sum_{l_{1}=l_{r}+r-1}^{k-1} (l_{1}+1/2) \sum_{l_{2}=l_{r}+r-2}^{l_{1}-1} (l_{2}+1/2) \dots \sum_{l_{r-1}=l_{r}+1}^{l_{r-2}-1} (l_{r-1}+1/2) \times \sum_{m_{1}=m_{r}+r-1}^{j-1} (m_{1}+1/2) \sum_{m_{2}=m_{r}+r-2}^{m_{1}-1} (m_{2}+1/2) \dots \sum_{m_{r-1}=m_{r}+1}^{m_{r-2}-1} (m_{r-1}+1/2) \le c(kj)^{2(r-1)}.$$

$$(2.10)$$

At first, we consider the case $1 < s < \infty$. For \triangle_{11} we have

$$\|\triangle_{11}\|_{L_2}^2 \le 4^{2r} \sum_{l_r=0}^{n-r+1} (l_r+1/2) \sum_{m_r=0}^{\infty} (m_r+1/2) \left(\sum_{k=n+1}^{\infty} \sum_{j=m_r+r}^{\infty} k^{\mu} j^{\mu} \left| \langle f, \varphi_{k,j} \rangle \right| \frac{B_{k,j}^r}{(kj)^{\mu-1/2}} \right)^2.$$

Using Hölder inequality and (2.10), for $\mu > 2r + \frac{s-1}{s} - 1/2$ we get

$$\|\triangle_{11}\|_{L_{2}}^{2} \leq c \sum_{l_{r}=0}^{n-r+1} (l_{r}+1/2) \sum_{m_{r}=0}^{\infty} (m_{r}+1/2) \left(\sum_{k=n+1}^{\infty} \sum_{j=m_{r}+r}^{\infty} k^{s\mu} j^{s\mu} |\langle f, \varphi_{k,j} \rangle|^{s} \right)^{2/s}$$

$$\times \left(\sum_{k=n+1}^{\infty} \sum_{j=m_{r}+r}^{\infty} (kj)^{-(\mu-2r+3/2)s/(s-1)} \right)^{2(s-1)/s}$$

$$\leq c \|f\|_{s,\mu}^{2} n^{-2(\mu-2r+3/2) + \frac{2(s-1)}{s}} \sum_{l_{r}=0}^{n-r+1} (l_{r}+1/2) \sum_{m_{r}=0}^{\infty} (m_{r}+1/2)^{-2(\mu-2r+3/2) + \frac{2(s-1)}{s} + 1}$$

$$\leq c \|f\|_{s,\mu}^2 n^{-2(\mu-2r+3/2)+\frac{2(s-1)}{s}+2}$$

Thus, we find

$$\|\triangle_{11}\|_{L_2} \le c\|f\|_{s,\mu} n^{-\mu + 2r + \frac{s-1}{s} - 1/2}.$$

Applying the estimating technique above we can bound the norm of $\triangle_{12}(t,\tau)$:

$$\|\triangle_{12}\|_{L_{2}}^{2} \leq c\|f\|_{s,\mu}^{2} \sum_{l_{r}=n-r+2}^{\infty} (l_{r}+1/2) \sum_{m_{r}=0}^{\infty} (m_{r}+1/2) \left(\sum_{k=l_{r}+r}^{\infty} \sum_{j=m_{r}+r}^{\infty} (kj)^{-(\mu-2r+3/2)s/(s-1)} \right)^{2(s-1)/s}$$

$$\leq c\|f\|_{s,\mu}^{2} \sum_{l_{r}=n-r+2}^{\infty} (l_{r}+1/2)^{-2(\mu-2r+3/2)+\frac{2(s-1)}{s}+1} \sum_{m_{r}=0}^{\infty} (m_{r}+1/2)^{-2(\mu-2r+3/2)+\frac{2(s-1)}{s}+1}$$

$$\leq c\|f\|_{s,\mu}^{2} n^{-2(\mu-2r+3/2)+\frac{2(s-1)}{s}+2}.$$

Summing up the estimates for $\triangle_{11}(t,\tau)$ and $\triangle_{12}(t,\tau)$ we obtain

$$\|\Delta_1\|_{L_2} \le \|\Delta_{11}\|_{L_2} + \|\Delta_{12}\|_{L_2} \le c\|f\|_{s,\mu} n^{-\mu + 2r + \frac{s-1}{s} - 1/2}$$

Now using the formula (2.7), from (2.5) we have

$$\Delta_2(t,\tau) = 4^r \sum_{k=r}^n \sum_{j=n+1}^\infty \sqrt{k+1/2} \sqrt{j+1/2} \left\langle f, \varphi_{k,j} \right\rangle$$

$$\times \sum_{l_1=r-1}^{k-1} (l_1+1/2) \sum_{l_2=r-2}^{l_1-1} (l_2+1/2) \dots \sum_{l_{r-1}=1}^{l_{r-2}-1} (l_{r-1}+1/2) \sum_{l_r=0}^{l_{r-1}-1} \sqrt{l_r+1/2} \varphi_{l_r}(t)$$

$$\times \sum_{m_1=r-1}^{j-1} (m_1+1/2) \sum_{m_2=r-2}^{m_1-1} (m_2+1/2) \dots \sum_{m_r=r-1}^{m_{r-2}-1} (m_{r-1}+1/2) \sum_{m_r=0}^{m_{r-1}-1} \sqrt{m_r+1/2} \varphi_{m_r}(\tau).$$

Further, we change the order of summation and get

$$\Delta_2(t,\tau) = \Delta_{21}(t,\tau) + \Delta_{22}(t,\tau),$$

where

$$\Delta_{21}(t,\tau) = 4^r \sum_{l_r=0}^{n-r} \sqrt{l_r + 1/2} \, \varphi_{l_r}(t) \sum_{m_r=0}^{n-r+1} \sqrt{m_r + 1/2} \, \varphi_{m_r}(\tau)$$

$$\times \sum_{k=l_r+r}^{n} \sum_{j=n+1}^{\infty} \sqrt{k+1/2} \sqrt{j+1/2} \, \langle f, \varphi_{k,j} \rangle \, B_{k,j}^r,$$

$$\Delta_{22}(t,\tau) = 4^r \sum_{l_r=0}^{n-r} \sqrt{l_r + 1/2} \, \varphi_{l_r}(t) \sum_{m_r=n-r+2}^{\infty} \sqrt{m_r + 1/2} \, \varphi_{m_r}(\tau)$$

$$\times \sum_{k=l_r+r}^{n} \sum_{j=m_r+r}^{\infty} \sqrt{k+1/2} \sqrt{j+1/2} \, \langle f, \varphi_{k,j} \rangle B_{k,j}^r.$$

Further, we estimate

$$\|\triangle_{21}\|_{L_2}^2 \le 4^{2r} \sum_{l_r=0}^{n-r} (l_r + 1/2) \sum_{m_r=0}^{n-r+1} (m_r + 1/2) \left(\sum_{k=l_r+r}^n \sum_{j=n+1}^{\infty} k^{\mu} j^{\mu} |\langle f, \varphi_{k,j} \rangle| \frac{B_{k,j}^r}{(kj)^{\mu-1/2}} \right)^2.$$

Again, applying the Hölder inequality, we get for $\mu > 2r + \frac{s-1}{s} - 1/2$:

$$\|\triangle_{21}\|_{L_2}^2 \le c\|f\|_{s,\mu}^2 n^{-2(\mu-2r+3/2)+\frac{2(s-1)}{s}}$$

$$\times \sum_{l_{n}=0}^{n-r} (l_{r}+1/2)^{-2(\mu-2r+3/2)+\frac{2(s-1)}{s}+1} \sum_{m_{r}=0}^{n-r+1} (m_{r}+1/2) \le c \|f\|_{s,\mu}^{2} n^{-2(\mu-2r+3/2)+\frac{2(s-1)}{s}+2}.$$

Further, we bound the norm of \triangle_{22} :

$$\|\triangle_{22}\|_{L_2}^2 \le c\|f\|_{s,\mu}^2 \sum_{l_r=0}^{n-r} (l_r+1/2) \sum_{m_r=n-r+2}^{\infty} (m_r+1/2) \left(\sum_{k=l_r+r}^n \sum_{j=m_r+r}^{\infty} (kj)^{-(\mu-2r+3/2)s/(s-1)} \right)^{2(s-1)/s}$$

$$\leq c \|f\|_{s,\mu}^2 n^{-2(\mu-2r+3/2)+\frac{2(s-1)}{s}+2}$$

Summing up the estimates for \triangle_{21} and \triangle_{22} we obtain

$$\|\Delta_2\|_{L_2} \le \|\Delta_{21}\|_{L_2} + \|\Delta_{22}\|_{L_2} \le c \|f\|_{s,\mu} n^{-\mu + 2r + \frac{s-1}{s} - 1/2}.$$

Using the formula (2.7), from (2.6) we have

$$\Delta_3(t,\tau) = 4^r \sum_{k=r}^n \sum_{j=\frac{rn}{k}}^n \sqrt{k+1/2} \sqrt{j+1/2} \left\langle f, \varphi_{k,j} \right\rangle$$

$$\times \sum_{l_1=r-1}^{k-1} (l_1+1/2) \sum_{l_2=r-2}^{l_1-1} (l_2+1/2) \dots \sum_{l_{r-1}=1}^{l_{r-2}-1} (l_{r-1}+1/2) \sum_{l_r=0}^{l_{r-1}-1} \sqrt{l_r+1/2} \varphi_{l_r}(t)$$

$$\times \sum_{m_1=r-1}^{j-1} (m_1+1/2) \sum_{m_2=r-2}^{m_1-1} (m_2+1/2) \dots \sum_{m_{r-1}=1}^{m_{r-2}-1} (m_{r-1}+1/2) \sum_{m_r=0}^{m_{r-1}-1} \sqrt{m_r+1/2} \varphi_{m_r}(\tau).$$

Further, using arguments similar to those above, we get

$$\|\Delta_3\|_{L_2} \le c \|f\|_{s,\mu} n^{-\mu + 2r + \frac{s-1}{s} - 1/2} \ln^{3/2 - 1/s} n.$$

The combination of (2.3) and bounds for the norms of \triangle_1 , \triangle_2 , \triangle_3 makes it possible to establish the desired inequality.

In the case of s = 1, Lemma is proved similarly.

The following statement contains an estimate for the second difference from the right-hand side of (2.2) in the metric of L_2 .

Lemma 2.4 Let condition (1.3) be satisfied. Then for an arbitrary function $f \in L_2(Q)$ it holds

$$\|\mathcal{D}_{n}^{(r,r)}f - \mathcal{D}_{n}^{(r,r)}f^{\delta}\|_{L_{2}} \leq c\delta n^{2r-1/p+1/2}\ln^{3/2-1/p}n$$

Proof. Let us write down the representation

$$\mathcal{D}_{n}^{(r,r)}f(t,\tau) - \mathcal{D}_{n}^{(r,r)}f^{\delta}(t,\tau) = \sum_{k,j \geq r, \ kj \leq rn-1} \langle f - f^{\delta}, \varphi_{k,j} \rangle \varphi_{k}^{(r)}(t) \varphi_{j}^{(r)}(\tau).$$

Using the formula (2.7), we get

$$\mathcal{D}_{n}^{(r,r)}f(t,\tau) - \mathcal{D}_{n}^{(r,r)}f^{\delta}(t,\tau) = 4^{r} \sum_{k=r}^{n-1} \sum_{j=r}^{\frac{r-1}{k}} \sqrt{k+1/2} \sqrt{j+1/2} \left\langle f - f^{\delta}, \varphi_{k,j} \right\rangle$$

$$\times \sum_{l_1=r-1}^{k-1} (l_1+1/2) \sum_{l_2=r-2}^{l_1-1} (l_2+1/2) \dots \sum_{l_{r-1}=1}^{l_{r-2}-1} (l_{r-1}+1/2) \sum_{l_r=0}^{l_{r-1}-1} \sqrt{l_r+1/2} \varphi_{l_r}(t)$$

$$\times \sum_{m_1=r-1}^{j-1} (m_1+1/2) \sum_{m_2=r-2}^{m_1-1} (m_2+1/2) \dots \sum_{m_{r-1}=1}^{m_{r-2}-1} (m_{r-1}+1/2) \sum_{m_r=0}^{m_{r-1}-1} \sqrt{m_r+1/2} \varphi_{m_r}(\tau).$$

Further, we change the order of summation and get

$$\mathcal{D}_{n}^{(r,r)} f(t,\tau) - \mathcal{D}_{n}^{(r,r)} f^{\delta}(t,\tau) = 4^{r} \sum_{l_{r}=0}^{n-r-1} \sqrt{l_{r}+1/2} \, \varphi_{l_{r}}(t) \sum_{m_{r}=0}^{\frac{r_{n}-1}{l_{r}+r}-r} \sqrt{m_{r}+1/2} \, \varphi_{m_{r}}(\tau)$$

$$\times \sum_{k=l_{r}+r}^{\frac{r_{n}-1}{m_{r}+r}} \sum_{i=m_{r}+r}^{\frac{r_{n}-1}{k}} \sqrt{k+1/2} \sqrt{j+1/2} \, \langle f - f^{\delta}, \varphi_{k,j} \, \rangle B_{k,j}^{r}.$$

Let 1 first. Then, using the Hölder inequality and estimate (2.10), we find

$$\|\mathcal{D}_{n}^{(r,r)}f - \mathcal{D}_{n}^{(r,r)}f^{\delta}\|_{L_{2}}^{2} \leq c\delta^{2} \sum_{l_{r}=0}^{n-r-1} (l_{r}+1/2) \sum_{m_{r}=0}^{\frac{rn-1}{l_{r}+r}-r} (m_{r}+1/2) \left(n^{\frac{(4r-3)p}{2(p-1)}+1} \sum_{k=l_{r}+r}^{\frac{rn-1}{m_{r}+r}} \frac{1}{k}\right)^{2(p-1)/p}$$

$$\leq c\delta^{2}n^{4r-3+\frac{2(p-1)}{p}}\ln^{\frac{2(p-1)}{p}}n\sum_{l=0}^{n-r-1}(l_{r}+1/2)\sum_{m_{r}=0}^{\frac{rn-1}{l_{r}+r}-r}(m_{r}+1/2) \times \delta^{2}n^{4r-1+2(p-1)/p}\ln^{3-2/p}n,$$

which was required to prove.

In the cases of p=1 and $p=\infty$, Lemma is proved similarly.

The combination of Lemmas 2.3 and 2.4 gives

Theorem 2.5 Let $f \in L^{\mu}_{s,2}$, $1 \le s < \infty$, $\mu > 2r - 1/s + 1/2$, condition (1.3) be satisfied. Then for $n \asymp \left(\delta^{-1} \ln^{1/p - 1/s} \frac{1}{\delta}\right)^{\frac{1}{\mu - 1/p + 1/s}}$ the following bound is valid

$$||f^{(r,r)} - \mathcal{D}_n^{(r,r)} f^{\delta}||_{L_2} \le c \left(\delta \ln^{1/s - 1/p} \frac{1}{\delta}\right)^{\frac{\mu - 2r + 1/s - 1/2}{\mu - 1/p + 1/s}} \ln^{3/2 - 1/s} \frac{1}{\delta}.$$

Proof. Taking into account Lemmas 2.3, 2.4, from (2.2) we get

$$||f^{(r,r)} - \mathcal{D}_n^{(r,r)} f^{\delta}||_{L_2} \le ||f^{(r,r)} - \mathcal{D}_n^{(r,r)} f||_{L_2} + ||\mathcal{D}_n^{(r,r)} f - \mathcal{D}_n^{(r,r)} f^{\delta}||_{L_2}$$

$$\le c||f||_{s,\mu} n^{-\mu + 2r + 1/2 - 1/s} \ln^{3/2 - 1/s} n + c\delta n^{2r - 1/p + 1/2} \ln^{3/2 - 1/p} n.$$

Substituting the rule $n \approx \left(\delta^{-1} \ln^{1/p-1/s} \frac{1}{\delta}\right)^{\frac{1}{\mu-1/p+1/s}}$ into the relation above completely proves Theorem. \square

Corollary 2.6 In the considered problem, the truncation method $\mathcal{D}_n^{(r,r)}$ (2.1) achieves (in the L_2 -metric) the accuracy

$$O\left(\left(\delta \ln^{1/s - 1/p} \frac{1}{\delta}\right)^{\frac{\mu - 2r + 1/s - 1/2}{\mu - 1/p + 1/s}} \ln^{3/2 - 1/s} \frac{1}{\delta}\right)$$

on the class $L_{s,2}^{\mu}$, $\mu > 2r - 1/s + 1/2$, and requires

$$\operatorname{card}(\Gamma_n) \simeq n \ln n \simeq \left(\delta^{-1} \ln^{\mu} \frac{1}{\delta}\right)^{\frac{1}{\mu - 1/p + 1/s}}$$

perturbed Fourier-Legendre coefficients.

Remark 2.7 Let us consider the standard version of the truncation method with $\Omega = \Box_n := [r, n] \times [r, n]$. Similar to [24] one can check that such an approach guarantees (in the L_2 -metric) the accuracy

$$O\left(\delta^{\frac{\mu-2r+1/s-1/2}{\mu+2r-2/p+1/s+1/2}}\right)$$

on the class $L_{s,2}^{\mu}$, $\mu > 2r - 1/s + 1/2$, and requires

$$\operatorname{card}(\square_n) \asymp n^2 \asymp \delta^{-\frac{2}{\mu + 2r - 2/p + 1/s + 1/2}}$$

perturbed Fourier-Legendre coefficients. The difference between these two versions of the truncation method (standard and (2.1)) is that within the method (2.1) we use only the Fourier-Legendre coefficients $\langle f^{\delta}, \varphi_{k,j} \rangle$ with indexes (k,j) from $\Gamma_n \subset \square_n$. The proposed approach (2.1) allows not only to reduce the amount of input information from $O(n^2)$ (for the standard version) to $O(n \ln n)$ (for (2.1)), but also to achieve higher accuracy.

3. Truncation method. Error estimate in the metric of C

Now we have to bound the error of (2.1) in the metric of C. An upper estimate for the norm of the difference (2.3) is contained in the following statement.

Lemma 3.1 Let $f \in L_{s,2}^{\mu}$, $1 \le s < \infty$, $\mu > 2r - 1/s + 3/2$. Then

$$||f^{(r,r)} - \mathcal{D}_n^{(r,r)}f||_C \le c||f||_{s,\mu}n^{-\mu+2r-1/s+3/2}\ln^{2-1/s}n.$$

Proof. Let us start with the case $1 < s < \infty$. Using (2.8) and (2.10), we get

$$\|\triangle_{11}\|_{C} \le c\|f\|_{s,\mu} \sum_{l_{r}=0}^{n-r+1} (l_{r}+1/2) \sum_{m_{r}=0}^{\infty} (m_{r}+1/2) \left(\sum_{k=n+1}^{\infty} \sum_{j=m_{r}+r}^{\infty} (kj)^{\frac{(-\mu+2r-3/2)s}{s-1}} \right)^{(s-1)/s}$$

$$\leq c \|f\|_{s,\mu} n^{-\mu+2r+3/2-1/s}.$$

Moreover, from (2.9) it follows

$$\|\triangle_{12}\|_{C} \le c\|f\|_{s,\mu} \sum_{l_{r}=n-r+2}^{\infty} (l_{r}+1/2) \sum_{m_{r}=0}^{\infty} (m_{r}+1/2) \left(\sum_{k=l_{r}+r}^{\infty} \sum_{j=m_{r}+r}^{\infty} (kj)^{\frac{(-\mu+2r-3/2)s}{s-1}} \right)^{(s-1)/s}$$

$$\leq c \|f\|_{s,\mu} n^{-\mu+2r+3/2-1/s}.$$

Thus, we get

$$\|\Delta_1\|_C \le c\|f\|_{s,\mu} n^{-\mu+2r+3/2-1/s}$$
.

Similarly, we find

$$\|\Delta_2\|_C \le c\|f\|_{s,\mu} n^{-\mu + 2r + 3/2 - 1/s},$$

$$\|\Delta_3\|_C \le c \|f\|_{s,\mu} n^{-\mu + 2r + 3/2 - 1/s} \ln^{2-1/s} n.$$

Substituting estimates for the norms of \triangle_1 , \triangle_2 , \triangle_3 into the relation (2.3) allows to establish the desired inequality.

In the case of s = 1, Lemma is proved similarly.

Remark 3.2 By the conditions of Lemma 3.1 we have that $\mu - 2r > 3/2 - 1/s \ge 1/2$. Thus, $f^{(r,r)}$ is continuous and the problem of estimating $f^{(r,r)}$ in the C-metric is well-defined.

The following statement contains an estimate for the second difference from the right-hand side of (2.2) in the metric of C.

Lemma 3.3 Assume that condition (1.3) is satisfied. Then for an arbitrary function $f \in C$ the following bound is valid

$$\|\mathcal{D}_{n}^{(r,r)}f - \mathcal{D}_{n}^{(r,r)}f^{\delta}\|_{C} \le c\delta n^{2r+3/2-1/p}\ln^{2-1/p}n.$$

Proof. Let 1 first. Then, using the Hölder inequality and the estimate (2.10), we find

$$\|\mathcal{D}_{n}^{(r,r)}f - \mathcal{D}_{n}^{(r,r)}f^{\delta}\|_{C} \le c\delta n^{2r-3/2 + \frac{p-1}{p}} \ln^{\frac{p-1}{p}} n \sum_{l_{r}=0}^{n-r-1} (l_{r} + 1/2) \sum_{m_{r}=0}^{\frac{rn-1}{l_{r}+r}-r} (m_{r} + 1/2)$$

$$\leq c\delta n^{2r+1/2+\frac{p-1}{p}}\ln^{\frac{p-1}{p}}n\sum_{l_n=0}^{n-r-1}\frac{1}{l_r+1/2} \asymp \delta n^{2r+3/2-1/p}\ln^{2-1/p}n,$$

which was required to prove.

In the cases of p = 1 and $p = \infty$, Lemma is proved similarly.

The combination of Lemmas 3.1 and 3.3 gives

Theorem 3.4 Let $f \in L_{s,2}^{\mu}$, $1 \le s < \infty$, $\mu > 2r - 1/s + 3/2$, and condition (1.3) be satisfied. Then for $n \asymp \left(\delta^{-1} \ln^{1/p - 1/s} \frac{1}{\delta}\right)^{\frac{1}{\mu - 1/p + 1/s}}$ the following bound is valid

$$||f^{(r,r)} - \mathcal{D}_n^{(r,r)} f^{\delta}||_C \le c \left(\delta \ln^{1/s - 1/p} \frac{1}{\delta}\right)^{\frac{\mu - 2r + 1/s - 3/2}{\mu - 1/p + 1/s}} \ln^{2-1/s} \frac{1}{\delta}.$$

Corollary 3.5 In the considered problem, the truncation method $\mathcal{D}_n^{(r,r)}$ (2.1) achieves (in the *C*-metric) the accuracy

$$O\left(\left(\delta \ln^{1/s - 1/p} \frac{1}{\delta}\right)^{\frac{\mu - 2r + 1/s - 3/2}{\mu - 1/p + 1/s}} \ln^{2 - 1/s} \frac{1}{\delta}\right)$$

on the class $L_{s,2}^{\mu}$, $\mu > 2r - 1/s + 3/2$, and requires

$$\operatorname{card}(\Gamma_n) \approx n \ln n \approx \left(\delta^{-1} \ln^{\mu} \frac{1}{\delta}\right)^{\frac{1}{\mu - 1/p + 1/s}}$$

perturbed Fourier-Legendre coefficients.

Remark 3.6 Consider the standard variant of the truncation method with $\Omega = \square_n$. Similar to [24] one can verify that this approach guarantees (in the *C*-metric) the accuracy

$$O\left(\delta^{\frac{\mu-2r+1/s-3/2}{\mu+2r-2/p+1/s+3/2}}\right)$$

on the class $L_{s,2}^{\mu}$, $\mu > 2r - 1/s + 3/2$, and requires

$$\operatorname{card}(\square_n) \asymp n^2 \asymp \delta^{-\frac{2}{\mu + 2r - 2/p + 1/s + 3/2}}$$

perturbed Fourier-Legendre coefficients. Comparison of the estimates found above with the corresponding estimates for the method $\mathcal{D}_n^{(r,r)}$ (2.1) (see Corollary 3.5) demonstrates that (2.1) is more efficient both in terms of accuracy and the amount of discrete information used.

Remark 3.7 Earlier, the method $\mathcal{D}_{n}^{(r,r)}$ (2.1) was studied for the problem of numerical differentiation of functions from $L_{s,2}^{\mu}$ in the case of r=1 & p=s=2 (see [24]) and in the case of r=2 & s=2 (see [25]). Thus, the results of Theorems 2.5 and 3.4 generalize studies of [24], [25] for the case of arbitrary r, p, s.

4. Minimal radius of Galerkin information

Now, we are in the position to find sharp estimates (in the power scale) for the minimal radius. First, we establish a lower estimate for the quantity $R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu},C,\ell_p)$. We fix an arbitrarily chosen domain $\hat{\Omega}$, $\operatorname{card}(\hat{\Omega}) \leq N$, of the coordinate plane $[r,\infty) \times [r,\infty)$ and construct an auxiliary function

$$f_1(t,\tau) = \widetilde{c} \left(\varphi_0(t) \varphi_0(\tau) + N^{-\mu - 1/s} r^{-\mu} \varphi_r(\tau) \sum_{k=N+r}^{3N+r} \varphi_k(t) \right),$$

where the sum $\sum_{k=N+r}^{3N+r}$ is taken over any N pairwise distinct functions $\varphi_k(t)$ such that $N+r \leq$

 $k \leq 3N + r$ and $(k,r) \notin \hat{\Omega}$. There is at least one set of such functions.

Now we estimate the norm of f_1 in the space metric $L_{s,2}^{\mu}$:

$$||f_1||_{s,\mu}^s = \widetilde{c}^s \left(1 + N^{-s\mu - 1} \sum_{k=N+r}^{3N+r} {}' k^{s\mu} \right) \le \widetilde{c}^s \left(1 + 4^{s\mu} \right).$$

Whence it follows that to satisfy the condition $||f_1||_{s,\mu} \leq 1$ it suffices to take

$$\widetilde{c} = \left(1 + 4^{s\mu}\right)^{-1/s}.\tag{4.1}$$

Next, we take another function from the class $L_{s,2}^{\mu}$:

$$f_2(t,\tau) = \widetilde{c}\,\varphi_0(t)\varphi_0(\tau).$$

Let us find a lower bound for the quantity $||f_1^{(r,r)} - f_2^{(r,r)}||_C$. For this we need formulas

$$\varphi_r^{(r)}(t) = \frac{\sqrt{r+1/2}}{2^{r-1/2}} \frac{(2r)!}{r!} \varphi_0(t), \quad f_2^{(r,r)}(t,\tau) \equiv 0,$$

$$f_1^{(r,r)}(t,\tau) = \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \sum_{k=N+r}^{3N+r} \varphi_k^{(r)}(t)$$

$$= 2^r \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \sum_{k=N+r}^{3N+r} \sqrt{k+1/2} \sum_{l_1=r-1}^{k-1} (l_1+1/2) \sum_{l_2=r-2}^{l_1-1} (l_2+1/2)$$

$$\dots \sum_{l_{r-1}=1}^{l_{r-2}-1} (l_{r-1}+1/2) \sum_{l_r=0}^{l_{r-1}-1} \sqrt{l_r+1/2} \varphi_{l_r}(t). \tag{4.2}$$

We note that in the right-hand side of (4.2) only terms with odd indexes $l_1 + k, l_2 + l_1, ..., l_r + l_{r-1}$ take part.

It is easy to see that

$$||f_1^{(r,r)} - f_2^{(r,r)}||_C \ge |f_1^{(r,r)}(1,1)| \ge \overline{c} N^{-\mu + 2r - 1/s + 3/2},$$

where

$$\overline{c} = \frac{\widetilde{c}\sqrt{r+1/2}}{4^r r^{\mu}} \frac{(2r)!}{(r!)^2}.$$
(4.3)

Since for any $1 \le p \le \infty$ it holds true

$$\|\overline{f}_1 - \overline{f}_2\|_{\ell_p} = \frac{\widetilde{c}}{r^{\mu}} N^{-\mu - 1/s + 1/p},$$

then in the case of $N^{-\mu-1/s+1/p} \leq r^{\mu}\delta/\widetilde{c}$, the functions

$$f_1^{\delta}(t,\tau) = f_2(t,\tau), \qquad f_2^{\delta}(t,\tau) = f_1(t,\tau)$$

can be considered as δ -perturbations of f_1 and f_2 , respectively.

Let us find an upper bound for $||f_1^{(r,r)} - f_2^{(r,r)}||_C$. Taking into account the relation $G(\hat{\Omega}, \overline{f_1}^{\delta}) = G(\hat{\Omega}, \overline{f_2}^{\delta})$, for any $\psi^{(r,r)}(\hat{\Omega}) \in \Psi(\hat{\Omega})$ we find

$$||f_1^{(r,r)} - f_2^{(r,r)}||_C \le ||f_1^{(r,r)} - \psi^{(r,r)}(G(\hat{\Omega}, \overline{f}_1^{\delta}))||_C + ||f_2^{(r,r)} - \psi^{(r,r)}(G(\hat{\Omega}, \overline{f}_2^{\delta}))||_C \le 2 \sup_{\substack{f \in L_{s,2}^{\mu}, \\ ||f||_{s,\mu} \le 1}} \sup_{\overline{f^{\delta}}: (1.3)} ||f^{(r,r)} - \psi^{(r,r)}(G(\hat{\Omega}, \overline{f}^{\delta}))||_C =: 2 \varepsilon_{\delta}(L_{s,2}^{\mu}, \psi^{(r,r)}(\hat{\Omega}), C, \ell_p).$$

That is

$$\varepsilon_{\delta}(L_{s,2}^{\mu}, \psi^{(r,r)}(\hat{\Omega}), C, \ell_p) \ge \frac{\overline{c}}{2} N^{-\mu + 2r - 1/s + 3/2}$$

From the fact that the domain $\hat{\Omega}$ and the algorithm $\psi^{(r,r)}(\hat{\Omega}) \in \Psi(\hat{\Omega})$ are arbitrary, follows that

$$R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu},C,\ell_p) \ge \frac{\overline{c}}{2} N^{-\mu+2r-1/s+3/2}.$$

Thus, the following assertion is proved.

Theorem 4.1 Let $1 \le s < \infty$, $\mu > 2r - 1/s + 3/2$, $1 \le p \le \infty$, $N \ge \left(r^{\mu} \delta/\tilde{c}\right)^{-1/(\mu + 1/s - 1/p)}$. Then

$$R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu},C,\ell_p) \ge \frac{\overline{c}}{2} N^{-\mu+2r-1/s+3/2},$$

where the constants \tilde{c} and \bar{c} are defined by (4.1) and (4.3) respectively.

The following assertion contains sharp estimates (in the power scale) for the minimal radius in the uniform metric.

Theorem 4.2 Let $1 \le s < \infty$, $\mu > 2r - 1/s + 3/2$, $1 \le p \le \infty$. Then for $N \asymp (\delta^{-1} \ln^{\mu} \frac{1}{\delta})^{\frac{1}{\mu - 1/p + 1/s}}$ it holds

$$N^{-\mu+2r-1/s+3/2} \leq R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu},C,\ell_p) \leq N^{-\mu+2r-1/s+3/2} \ln^{\mu-2r+1/2} N$$

or

$$\left(\delta \ln^{-\mu} \frac{1}{\delta}\right)^{\frac{\mu - 2r + 1/s - 3/2}{\mu - 1/p + 1/s}} \preceq R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu}, C, \ell_p) \preceq \left(\delta \ln^{1/s - 1/p} \frac{1}{\delta}\right)^{\frac{\mu - 2r + 1/s - 3/2}{\mu - 1/p + 1/s}} \ln^{2 - 1/s} \frac{1}{\delta}.$$

The upper bound is realized by (2.1) for $n \simeq \left(\delta^{-1} \ln^{1/p-1/s} \frac{1}{\delta}\right)^{\frac{1}{\mu-1/p+1/s}}$.

Proof. The upper bound for $R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu},C,\ell_p)$ follows from Theorem 3.4. The lower bound is found in Theorem 4.1.

Let us turn to estimate the minimal radius in the integral metric.

Theorem 4.3 Let $1 \le s < \infty$, $\mu > 2r - 1/s + 1/2$, $1 \le p \le \infty$. Then for any $N \ge \left(r^{\mu}\delta/\widetilde{c}\right)^{-1/(\mu+1/s-1/p)}$ it holds

$$R_{N\delta}^{(r,r)}(L_{s,2}^{\mu}, L_2, \ell_p) \ge \overline{\overline{c}} N^{-\mu + 2r - 1/s + 1/2},$$

where $\overline{c} = \frac{\widetilde{c}\sqrt{r+1/2}}{2^{4r-3/2}r^{\mu-1}} \frac{(2r)!}{(r!)^2}$ and \widetilde{c} is defined by (4.1).

Proof. The proof of Theorem 4.3 almost completely coincides with the proof of Theorem 4.1, including the form of the auxiliary functions f_1 , f_1^{δ} , f_2 , f_2^{δ} . The only difference is in the lower estimate of the norm of the difference $f_1^{(r,r)} - f_2^{(r,r)}$. Changing the order of summation in (4.2) yields to the representation

$$f_1^{(r,r)}(t,\tau) = 2^r \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \sum_{k=N+r}^{3N+r} \sqrt{k+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \sum_{k=N+r}^{3N+r} \sqrt{k+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \sum_{k=N+r}^{3N+r} \sqrt{k+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \sum_{k=N+r}^{3N+r} \sqrt{k+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \sum_{k=N+r}^{3N+r} \sqrt{k+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \sum_{k=N+r}^{3N+r} \sqrt{k+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \sum_{k=N+r}^{3N+r} \sqrt{k+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \sum_{k=N+r}^{3N+r} \sqrt{k+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \sum_{k=N+r}^{3N+r} \sqrt{k+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu-1/s} \varphi_r^{(r)}(\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu} N^{-\mu} (\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu} (\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu} (\tau) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu} (\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu} (\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}} N^{-\mu} (\tau) \Big(\sum_{l_r=0}^N \sqrt{l_r+1/2} \varphi_{l_r}(t) \Big) + \frac{\widetilde{c}}{r^{\mu}}$$

$$+\sum_{l_r=N+1}^{3N}\sqrt{l_r+1/2}\,\varphi_{l_r}(t)\sum_{k=l_r+r}^{3N+r}\sqrt{k+1/2}\Big)\sum_{l_1=l_r+r-1}^{k-1}(l_1+1/2)\sum_{l_2=l_r+r-2}^{l_1-1}(l_2+1/2)\dots\sum_{l_{r-1}=l_r+1}^{l_{r-2}-1}(l_{r-1}+1/2).$$

It is easy to verify that

$$||f_1^{(r,r)} - f_2^{(r,r)}||_{L_2}^2 \ge \frac{c'}{4^{4r-2}((r-1)!)^2} N^{-2\mu+4r-2/s+1},$$

where

$$c' = \frac{2\tilde{c}^2}{r^{2\mu}}(r+1/2)\frac{((2r)!)^2}{(r!)^2}.$$

Whence we obtain the relation

$$\varepsilon_{\delta}(L_{s,2}^{\mu},\psi^{(r,r)}(\hat{\Omega}),L_2,\ell_p) \geq \overline{\overline{c}} N^{-\mu+2r-1/s+1/2}$$

is true for any $N \ge \left(r^{\mu}\delta/\widetilde{c}\right)^{-1/(\mu+1/s-1/p)}$. From the fact that the domain $\hat{\Omega}$ and the algorithm $\psi^{(r,r)}(\hat{\Omega}) \in \Psi(\hat{\Omega})$ are arbitrary, it follows that

$$R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu}, L_2, \ell_p) \ge \overline{\overline{c}} N^{-\mu + 2r - 1/s + 1/2}.$$

Thus, the proof of Theorem 4.3 has been complied.

The following statement contains sharp estimates (in the power scale) for the minimal radius in the integral metric.

Theorem 4.4 Let $1 \le s < \infty$, $\mu > 2r - 1/s + 1/2$, $1 \le p \le \infty$. Then for $N \asymp \left(\delta^{-1} \ln^{\mu} \frac{1}{\delta}\right)^{1/(\mu - 1/p + 1/s)}$ it holds

$$N^{-\mu+2r-1/s+1/2} \leq R_{N\delta}^{(r,r)}(L_{s,2}^{\mu}, L_2, \ell_p) \leq N^{-\mu+2r-1/s+1/2} \ln^{\mu-2r+1} N$$

or

$$\left(\delta \ln^{-\mu} \frac{1}{\delta}\right)^{\frac{\mu - 2r + 1/s - 1/2}{\mu - 1/p + 1/s}} \preceq R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu}, L_2, \ell_p) \preceq \left(\delta \ln^{1/s - 1/p} \frac{1}{\delta}\right)^{\frac{\mu - 2r + 1/s - 1/2}{\mu - 1/p + 1/s}} \ln^{3/2 - 1/s} \frac{1}{\delta}.$$

The upper bound is realized by (2.1) for $n \simeq \left(\delta^{-1} \ln^{1/p-1/s} \frac{1}{\delta}\right)^{\frac{1}{\mu-1/p+1/s}}$.

Proof. The upper bound for $R_{N,\delta}^{(r,r)}(L_{s,2}^{\mu},L_2,\ell_p)$ follows from Theorem 2.5. The lower bound is found in Theorem 4.3.

5. Computational experiments

To demonstrate the effectiveness of the proposed method for recovering high-order derivatives some numerical experiments were carried out. The calculations were performed on a computer with a 4-core Intel Core i5 processor and 16 GB memory in the mathematical modeling environment MATLAB 2022a.

5.1. Example 1

We consider the function $F_1(t,\tau) = f(t)f(\tau)/C$, where C = 754 and

$$f(t) = \begin{cases} -1/8t^2 + 1/12t^4 - 1/20t^5 + 1/42t^7 - 3/224t^8, & -1 \le t < 0, \\ -1/8t^2 + 1/12t^4 - 1/20t^5 + 1/45t^7 - 3/240t^8, & 0 \le t \le 1. \end{cases}$$

Taking into account the definition of space $L_{s,2}^{\mu}$ it is easy to see that $||F_1||_{2,\mu} \approx 1$ for $\mu = 5, 5$ and $||F_1^{(2,2)}||_{L_2} \approx 10^{-4}$.

The simulation of the noise in the input data was done in two different ways:

- a random noise adds to the values of the Fourier-Legendre coefficients. The noise is generated by the $\mathbf{randn}(\mathbf{size}(\mathcal{F}))\delta$ command, where \mathbf{randn} and \mathbf{size} are standard functions of the MATLAB system, and \mathcal{F} is a matrix for exact values of the Fourier-Legendre coefficients;
- the values of the Fourier-Legendre coefficients recovered by the quadrature trapezoid formula on a uniform grid with a step h so that condition (1.3) is satisfied for a given δ .

Numerical experiments were carried out for the following error levels: $\delta = 10^{-6}, 10^{-7}, 10^{-8}$. Tables 1 and 2 show the results of numerical calculations for the approximation of $F_1^{(2,2)}$ by the truncation method (2.1) for two different types of noise (random data noise and trapezoid formula errors). "Error in L_2 " and "Error in C" columns contain the recovery accuracy in the L_2 — and C— metrics respectively, n and $\operatorname{card}(\Gamma_n)$ columns indicate the highest degree of

Legendre polynomial and the number of Fourier-Legendre coefficients involved, resp. Also in Table 2, h means the step size in the quadrature formula.

The graphs 1 and 2 show the exact derivative $F_1(2,2)$ and its approximations constructed on data with random noise and with a noise, generated by the trapezoid formula, respectively.

Tabl. 1: The results of recovering derivative $F_1^{(2,2)}$ for random noise

δ	10^{-6}	10^{-7}	10^{-8}
Error in L_2	$1, 1 \cdot 10^{-4}$	$2,73 \cdot 10^{-5}$	$6,7 \cdot 10^{-6}$
Error in C	$1, 2 \cdot 10^{-3}$	$3,4\cdot 10^{-4}$	$5 \cdot 10^{-5}$
n	19	24	31
$\operatorname{card}(\Gamma_n)$	52	80	106

Tabl. 2: The results of recovering derivative $F_1^{(2,2)}$ for noise from quadrature formula

δ	10^{-6}	10^{-7}	10^{-8}		
Error in L_2	$4.8 \cdot 10^{-5}$	$3, 2 \cdot 10^{-5}$	$6, 6 \cdot 10^{-6}$		
Error in C	$7,53 \cdot 10^{-4}$	$4,9 \cdot 10^{-4}$	$2,53 \cdot 10^{-5}$		
n	19	24	31		
h	$1,16\cdot 10^{-4}$	$8 \cdot 10^{-5}$	$4 \cdot 10^{-5}$		

As can be seen from the graphs and tables above, for both types of noise, the truncation method gives the same order of accuracy for recovering the derivative $F_1^{(2,2)}$. At the same time, applying the quadrature formula expands the area of using the proposed method in computational problems, especially in the situation when the input data are given in the form of a set of function values at the grid nodes.

5.2. Example 2

Let us test the method (2.1) on an analytic function. Following [36], we take the function $F_2(t,\tau)=(2-(2t-1)^2)^2\cos(4\tau)/C$. Let us put $\mu=6$. It is easy to check that $\|F_2\|_{2,6}\approx 1$ and $\|F_2^{(2,2)}\|_{L_2}\approx 10^{-4}$, if C=43940129. The Fourier-Legendre coefficients of the considered function are calculated using the quadrature trapezoid formula for $h=4\cdot 10^{-4}, 10^{-4}, 4\cdot 10^{-5}$, which in turn according to formula (1.3) matches $\delta\approx 10^{-6}, 10^{-7}, 10^{-8}$, respectively.

Tabl. 3: The results of recovering derivative $F_2^{(2,2)}$ for noise from quadrature formula

δ	10^{-6}	10^{-7}	10^{-8}
Error in L_2	$3.8 \cdot 10^{-5}$	$1 \cdot 10^{-6}$	$1.53 \cdot 10^{-7}$
Error in C	$1,85 \cdot 10^{-4}$	$6.37 \cdot 10^{-6}$	$8.17 \cdot 10^{-7}$
n	11	18	25
h	$4 \cdot 10^{-4}$	10^{-4}	$4 \cdot 10^{-5}$

The results of the numerical experiment are shown in Table 3 and Graph 3.

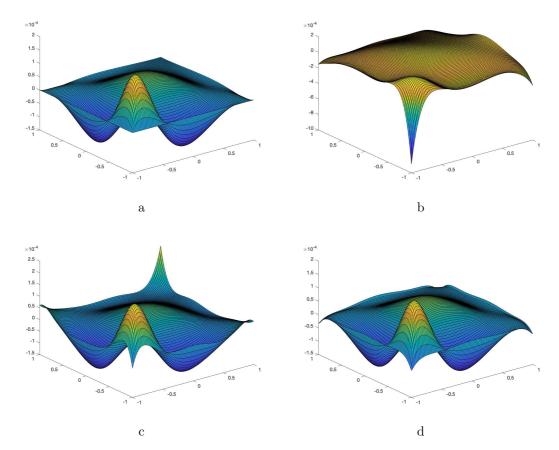


Fig. 1: Recovery of the derivative $F_1^{(2,2)}$ with random noise in the input data . The exact derivative $F_1^{(2,2)}$ (Fig. a); approximation to $F_1^{(2,2)}$ for $\delta=10^{-6}$ (Fig. b); for $\delta=10^{-7}$ (Fig. c) and $\delta=10^{-8}$ (Fig. d),

The numerical results given above show that the proposed method (2.1) works efficiently and coincides well with the theoretical results.

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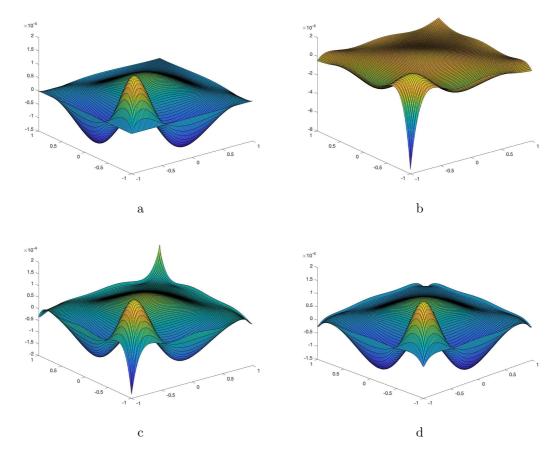


Fig. 2: Recovery of the derivative $F_1^{(2,2)}$ with noise in the input data arising from the quadrature formula. The exact derivative $F_1^{(2,2)}$ (Fig. a); approximation to $F_1^{(2,2)}$ for $\delta=10^{-6}$ (Fig. b); for $\delta=10^{-7}$ (Fig. c) and $\delta=10^{-8}$ (Fig. d).

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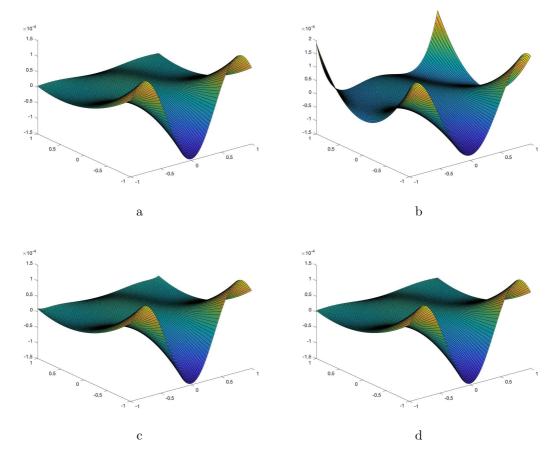


Fig. 3: Recovery of the derivative $F_2^{(2,2)}$ with noise in the input data arising from the quadrature formula. The exact derivative $F_2^{(2,2)}$ (Fig. a); approximation to $F_2^{(2,2)}$ for $\delta=10^{-6}$ (Fig. b); for $\delta=10^{-7}$ (Fig. c) and $\delta=10^{-8}$ (Fig. d).

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